

# Approximate Mortar Conditions for the CR Finite Element on Non-matching Grids

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## The Crouzeix-Raviart Mortar FE

We consider the problem: Find  $u^* \in H^1_0(\Omega)$  such that

$$a(u^*,v)=f(v),\quad \forall v\in H^1_0(\Omega),$$

where  $\overline{\Omega} = \bigcup_i \overline{\Omega_i}$  (Non-overlapping),  $a(u, v) = \sum_{i=1}^N \int_{\Omega_i} \nabla u \cdot \nabla v \, dx$ and  $f(v) = \sum_{i=1}^N \int_{\Omega_i} fv \, dx$ .



 $X_h(\Omega_i)$ : The CR (P1 Non-conforming) FE space on  $\Omega_i$ , vanishing at the edge-mid nodes on the boundary  $\partial \Omega$ .

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Mortar condition: For  $u_i \in X_h(\Omega_i)$  and  $u_j \in X_h(\Omega_j)$ ,

$$Q_m u_i = Q_m u_j$$
 on  $\delta_m$ .

 $Q_m : L^2(\Gamma_{ij}) \to M(\delta_m)$  is the  $L^2$  projection operator.  $M(\delta_m)$  is the set of piecewise constant functions on the  $\delta_m$ -discretization.



CR mortar FE space: [cf. Marcinkowski]

 $V_h = \{ u \in \prod_i X_h(\Omega_i) : u \text{ satisfies the Mortar condition on } \delta_m \}$ 

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#### **Approximate Mortar Condition**

Basis functions of  $V_h$ , associated with the subdomain interior, may have non-zero support on the non-mortar side (Not desirable, specially in 3D).



Aim is to avoid their direct use in applying the mortar condition. Alternative: Approximate mortar condition  $(I_m : X_h(\gamma_m) \to L^2(\gamma_m))$ :

$$Q_m I_m u_i = Q_m u_j \quad \text{ on } \delta_m$$

Approximate mortar condition in Wavelet context, cf. [Bertoluzza]

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### Choice of $I_m$

### **First choice**

Being inspired by the fact that the basis functions associated with the subdomain interior, have zero integrals on the subdomain boundary, an easy choice would be to define  $I_m$  as

$$I_m u_{i|_e} = \frac{1}{|e|} \int_e u_i \, dx \quad \text{ on } e \subset \gamma_m$$

The operator  $I_m$  preserves only constants, and is therefore not enough for getting optimal error estimate. We need locally at least P1-preserving.

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The interpolation is done basically by first joining values at the neighboring edge-mid nodes by straight lines, and then simply extending the two end straight lines towards the end of the mortar  $\gamma_m$ .

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### The discrete problem

Since  $V_h$  does not belong to  $H_0^1(\Omega)$ , we use the broken bilinear form  $a_h(\cdot, \cdot)$  defined as  $a_h(u, v) = \sum_{i=1}^N \sum_{\tau \in \mathcal{T}_h(\Omega_i)} \int_{\tau} \nabla u \cdot \nabla v \, dx$ . The discrete problem takes the following form: Find  $u_h^* \in V_h$  such that

$$a_h(u_h^*, v_h) = f(v_h), \quad \forall v_h \in V_h.$$

 $V_h$  is a Hilbert space with an inner product defined by  $a_h(u_h, v_h)$ . An error estimate for the new mortar technique for CR FE (optimal):

$$|u^* - u^*_h|^2_{H^1_h(\Omega)} \le c \sum_{i=1}^N h^2_i |u^*|^2_{H^2(\Omega_i)}$$

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## An additive Schwarz method

### The splitting

$$V_h = \sum_{\gamma_m} V_\gamma + V_0 + \sum_i^N V_i$$

 $V_i$  and  $V_{\gamma}$  are the standard local subspaces of  $V_h$  on subdomain  $\Omega_i$  and mortar edge  $\gamma_m$ .  $V_0 \subset V_h = span\{\Phi_i\}_{i=1,\dots,N}$  is a coarse space similar to the one in [Dryja-Bjørstad-Rahman].

The function  $\Phi_i$  is defined by its values at the CR nodes:

$$\Phi_i(x) = \begin{cases} 1, & x \in \Omega_i^{CR}, \\ \frac{\rho_i}{(\rho_i + \rho_j)}, & x \in \gamma_m^{CR}, \gamma_m \subset \partial \Omega_i \cap \partial \Omega_j, \\ 0, & \text{elsewhere.} \end{cases}$$

Values on the non-mortar sides are defined using the mortar condition.

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### **Projection-like**

For 
$$u \in V_h$$
, and  $\alpha \in \{\{\gamma_m\}, 0, 1, \cdots, N\}, \quad T_\alpha : V_h \to V_\alpha$ 
$$a_h(T_\alpha u, v) = a_h(u, v), \quad v \in V_\alpha$$

#### An equivalent discrete problem

Set 
$$T = \sum_{\gamma_m} T_{\gamma_m} + T_0 + \sum_{i=1}^N T_i.$$
  
$$a_h(Tu, v) = a_h(u, v), \quad v \in V_\alpha$$

### Convergence

$$\kappa_2(T) \le c \frac{H}{h}$$

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## **Numerical Results**

Unit square domain  $\Omega$ , initially divided into 9 square subdomains, each consisting of  $2m_i^2$  triangles corresponding to the mesh size  $h_i \approx 1/m_i$ . For each neighboring subdomain pair  $\{\Omega_i, \Omega_j\}, m_i \neq m_j$  in order to get non-matching grids. The function f is so chosen that the exact solution is equal to  $sin(\pi x)sin(\pi y)$ .

$\{h_i, h_j\}$	Standard Mortar		Approximate Mortar	
	$L^2 - error$	$H^1 - error$	$L^2 - error$	H1 - error
$\{1/6, 1/5\}$	0.00202040	0.065292	0.00248430	0.078409
$\{1/12, 1/10\}$	0.00049658	0.032843	0.00066682	0.038768
$\{1/24, 1/20\}$	0.00012307	0.016479	0.00017493	0.019321

 $L^2$ -norm and  $H^1$ -seminorm of the error of the two mortar methods, showing that the two methods are quite close.

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$\{h_i,h_j\}$	Standard Mortar	Approximate Mortar
$\{1/6, 1/5\}$	28.85 (25)	30.11 (23)
$\{1/12, 1/10\}$	63.44 (35)	60.90 (31)
$\{1/24, 1/20\}$	134.18 (49)	122.55 (45)

Condition number estimates (PCG-iteration counts in parentheses), showing similar performance for the two mortar methods.



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