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Hydro-Elastic Contributions to Fatigue Damage on a Large Monopile

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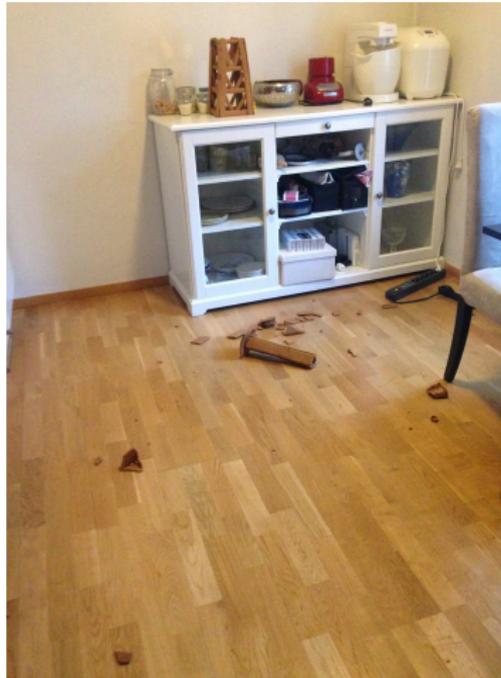
^cStatkraft

January 21, 2016

Motivation



Motivation



Motivation



Overview

Introduction

Structural model

Hydrodynamic models

Challenges

Results

Evaluation by time-frequency analysis

Conclusions



Hydrodynamic Modeling

Motivation:

- Need accurate and realistic load models to evaluate control strategies
- Upscaling of monopiles will give new response characteristics
- Load theory validation and relative impact on lifetime estimation



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Findings:

- Significant higher order contributions to fatigue damage in sea-states with $H_S > D/2$
- Necessary to include diffraction effects on second order inertia forces
- Higher order loads more prominent at low damping levels



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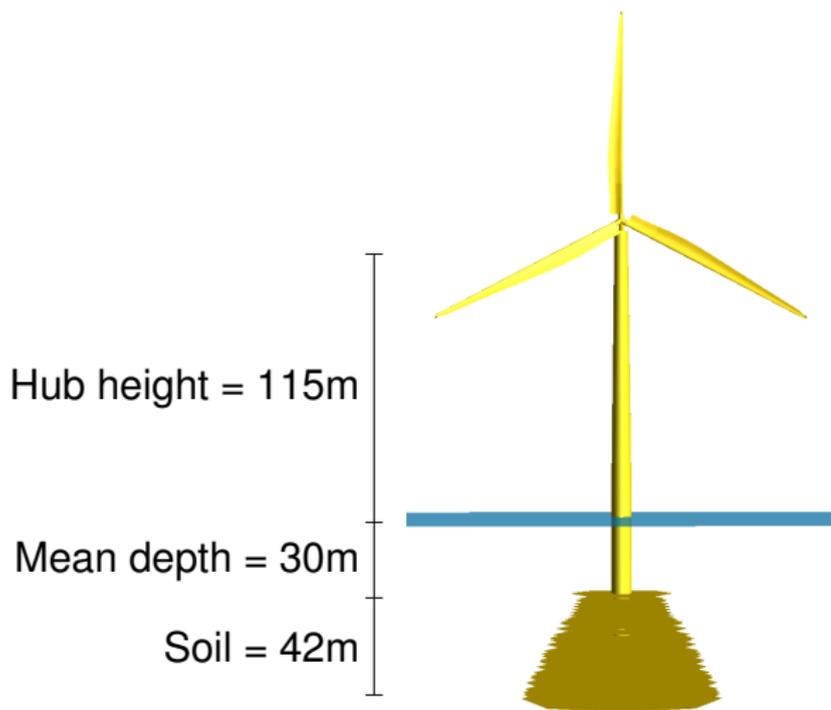
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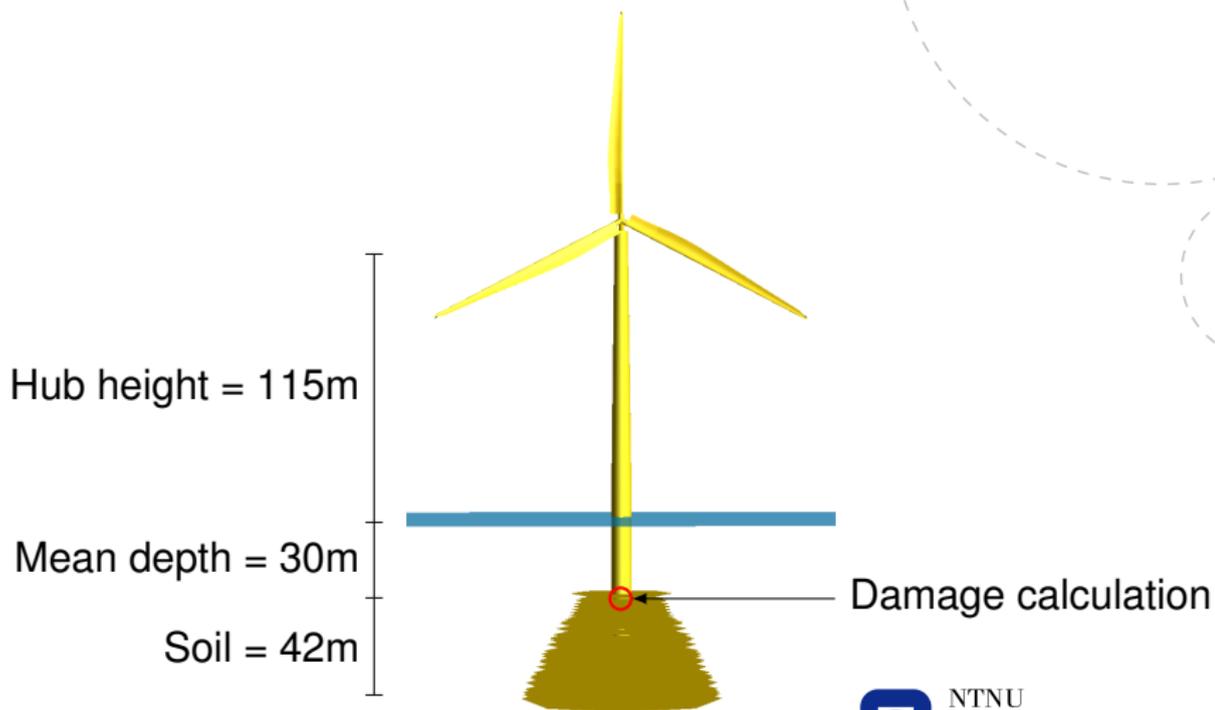
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Model



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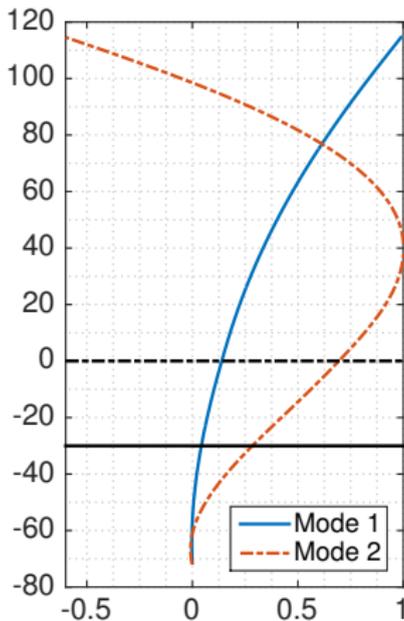
Parameters:

| | |
|---------------------|--|
| Diameter | 9m |
| Depth | 30m |
| Structural damping | 3% of critical damping (Rayleigh) |
| Aerodynamic damping | Constant Rayleigh included in structural |
| Soil | Non-linear springs for sand and clay |
| Natural periods | Mode 1: 4.2s, Mode 2: 1.0s |
| Sea-states | FLS |



Modal analysis

Mode 1: 4.1 seconds
Mode 2: 1.0 seconds



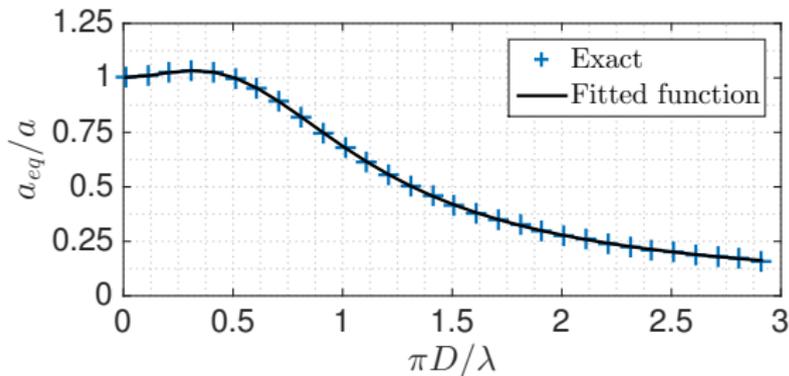
Wave load models

| | |
|------|--|
| O1 | Linear waves |
| O1D | Linear waves with diffraction (MacCamy and Fuchs) |
| O2 | Second order contribution from kinematics stretching |
| O3 | Third+fourth order contribution from kinematics stretching |
| FNV3 | Third order FNV - direct implementation |
| O1P | First order distributed pressure from panel code |
| O2P | Second order total force from panel code |



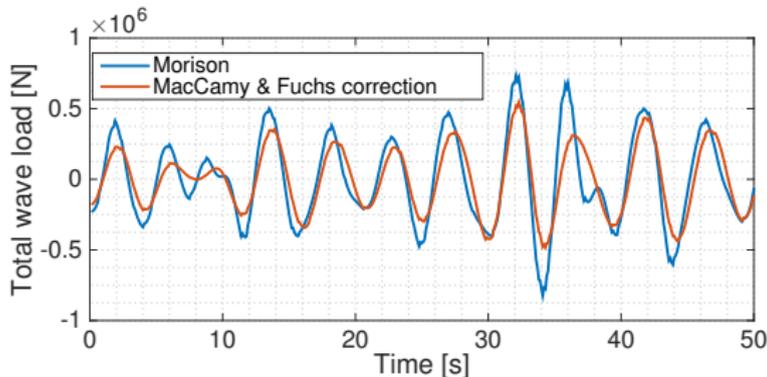
Diffraction - MacCamy and Fuchs

Correction of wave load due to interaction with large-volume structure. a_{eq} = equivalent water particle acceleration.

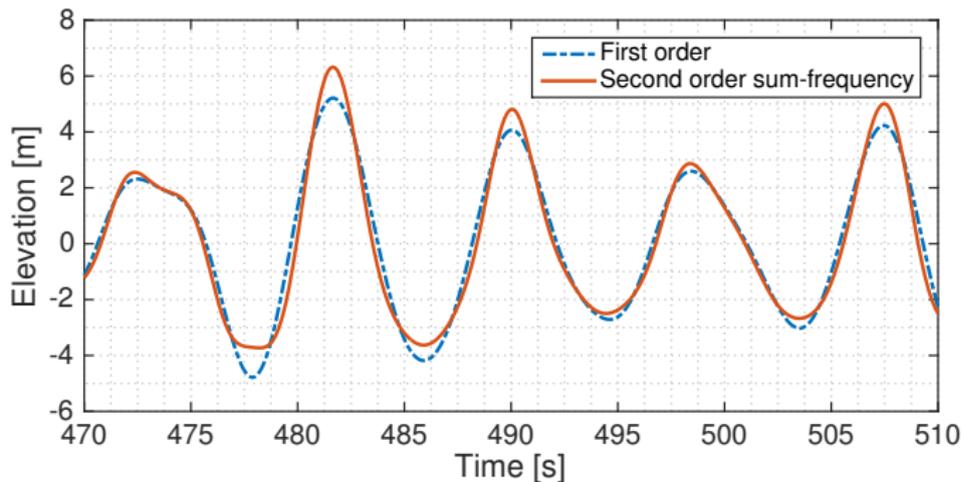


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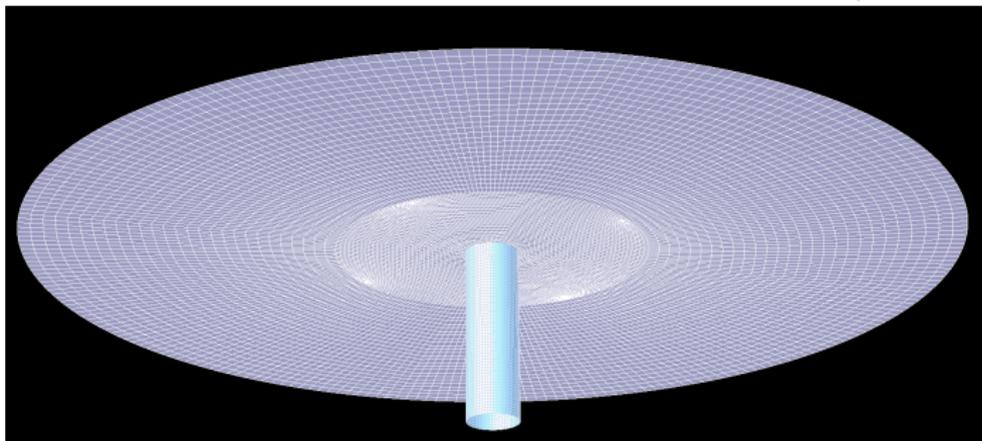


Second order wave elevation



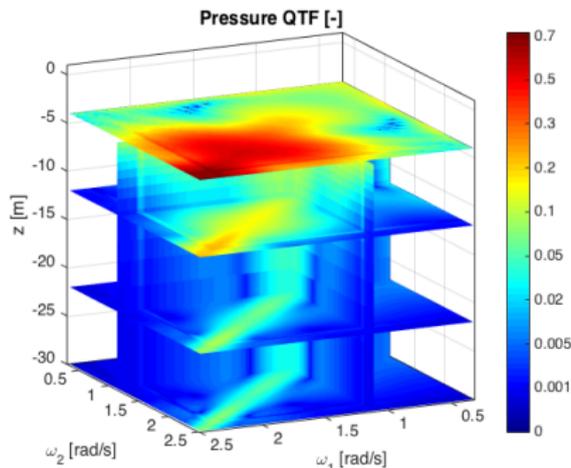
Second order sum-frequency from panel code

HydroD - Wadam



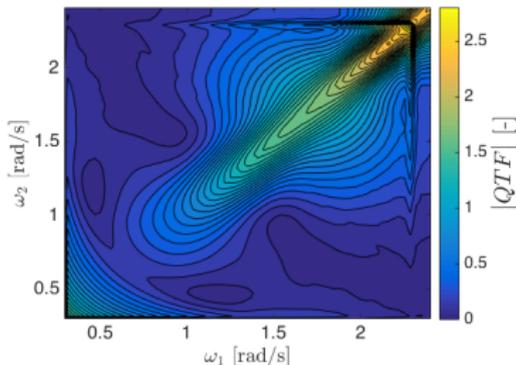
Second order sum-frequency from panel code

Non-dimensional
resulting pressure
in x-direction over
the column as a
function of ω_1 , ω_2
and z .



Second order sum-frequency from panel code

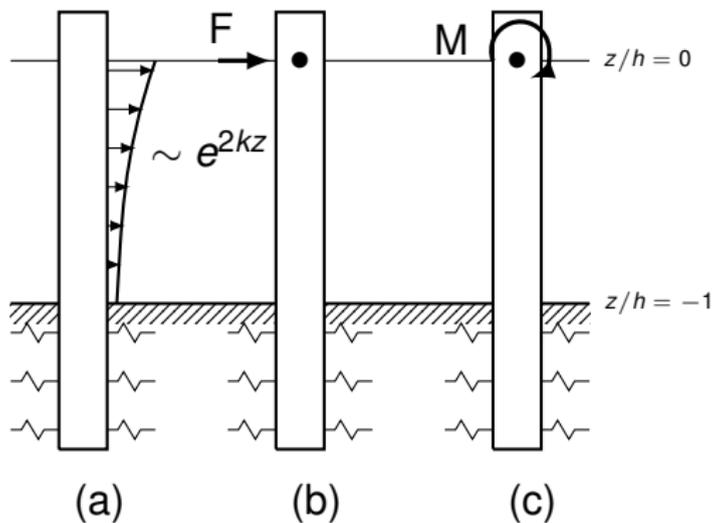
The second-order pressures are lumped to $z = 0$ and act as a point force.



$$\bar{F} = \Re \left\{ \sum_{n=1}^N \sum_{m=1}^M \zeta_{a,n} \cdot \zeta_{a,m} \cdot QTF \cdot e^{-i(\phi_m + \phi_n)} \right\}$$

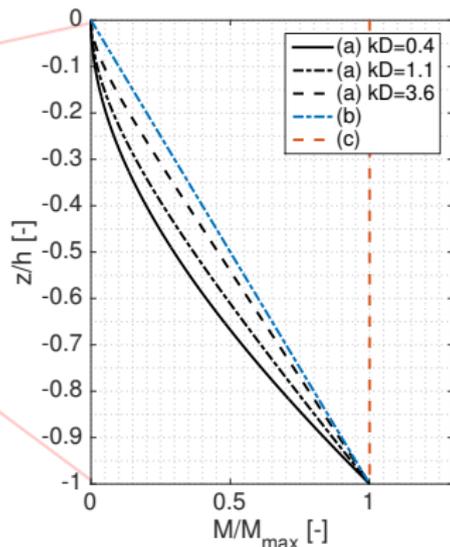
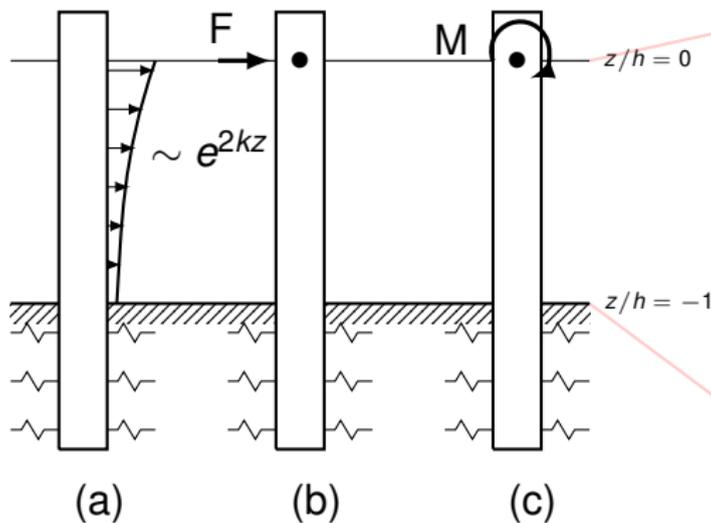
Load application

Distributed, point force or moment?



Load application

Distributed, point force or moment?



Third Order FNV

Third order horizontal force from linear elevation and diffraction potential assuming deep water:

$$F_x^{FNV(3)} = \rho\pi r^2 \left[\zeta_1 \left(\zeta_1 u_{tz} + 2ww_x + uu_x - \frac{2}{g} u_t w_t \right) - \left(\frac{u_t}{g} \right) (u^2 + w^2) + \frac{\beta}{g} u^2 u_t \right]_{z=0}$$



Kinematics models

Assuming

$$k\zeta_a = O(\epsilon)$$

$$kD = O(\delta)$$

where $\epsilon \ll 1$ and $\epsilon \approx \delta$

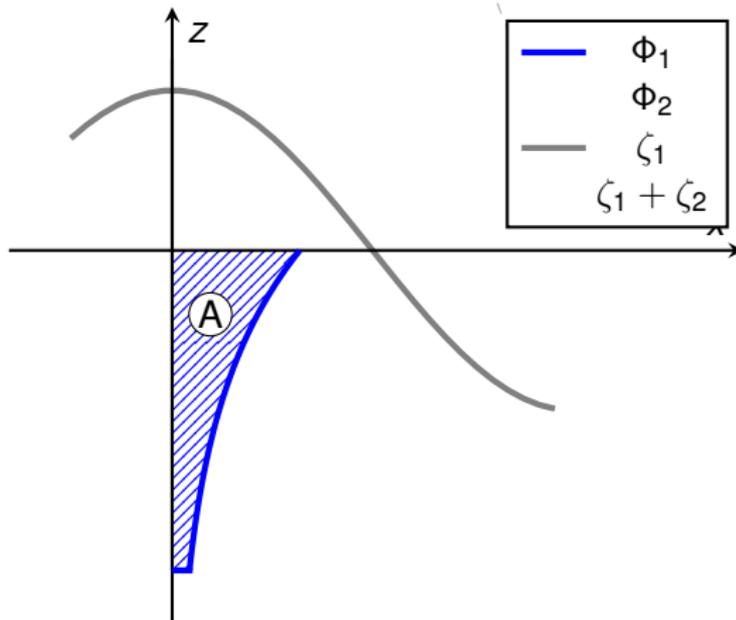


Kinematics models

Assuming
 $k\zeta_a = O(\epsilon)$
 $kD = O(\delta)$

Order of horizontal
inertia forces:

— A: $\epsilon^1 \delta^2$

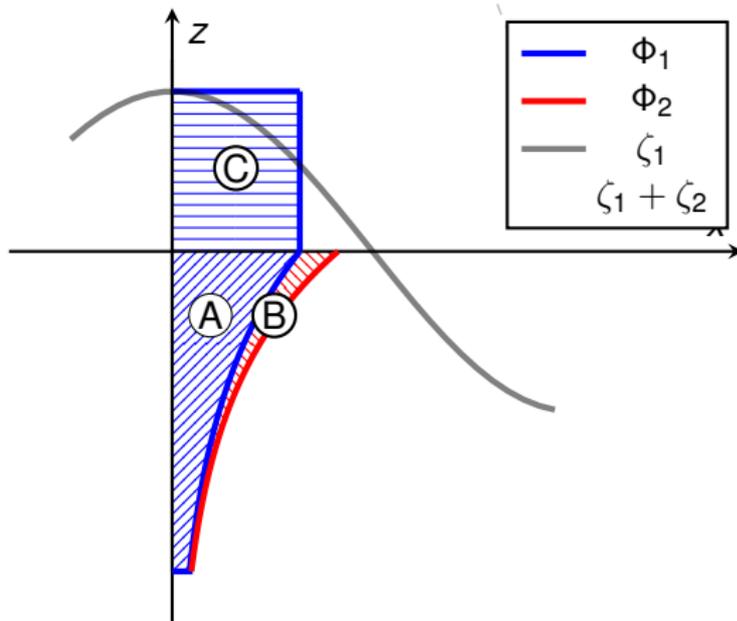


Kinematics models

Assuming
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Order of horizontal
inertia forces:

- A: $\epsilon^1 \delta^2$
- B+C: $\epsilon^2 \delta^2$

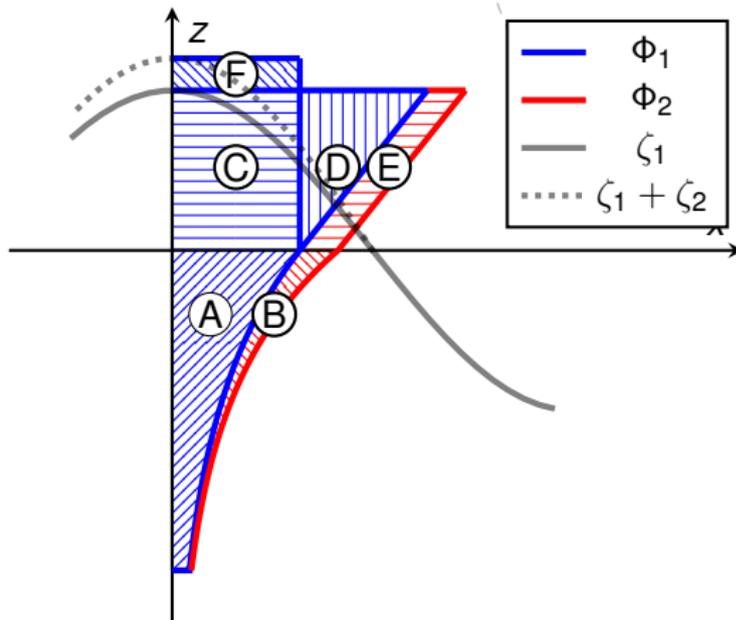


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Order of horizontal
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- A: $\epsilon^1 \delta^2$
- B+C: $\epsilon^2 \delta^2$
- D+E+F: $\epsilon^3 \delta^2$

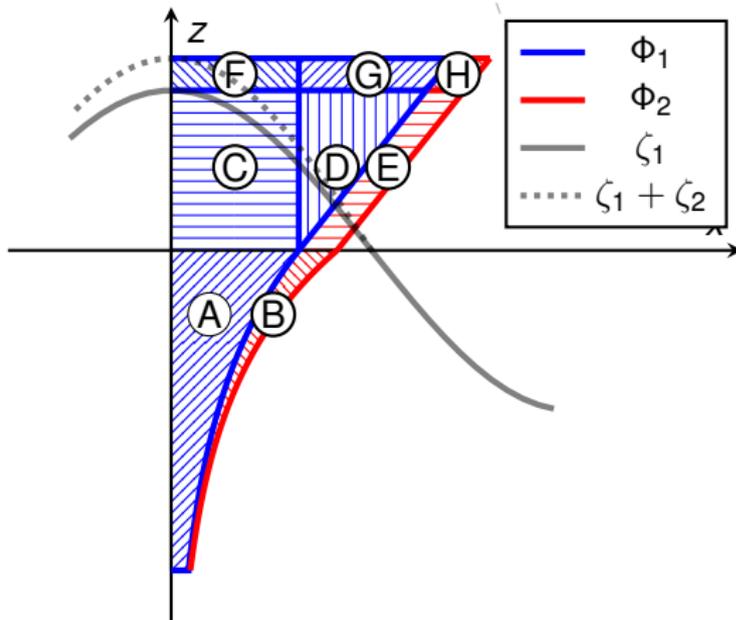


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Assuming
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Order of horizontal
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- A: $\epsilon^1 \delta^2$
- B+C: $\epsilon^2 \delta^2$
- D+E+F: $\epsilon^3 \delta^2$
- G+H: $\epsilon^4 \delta^2$

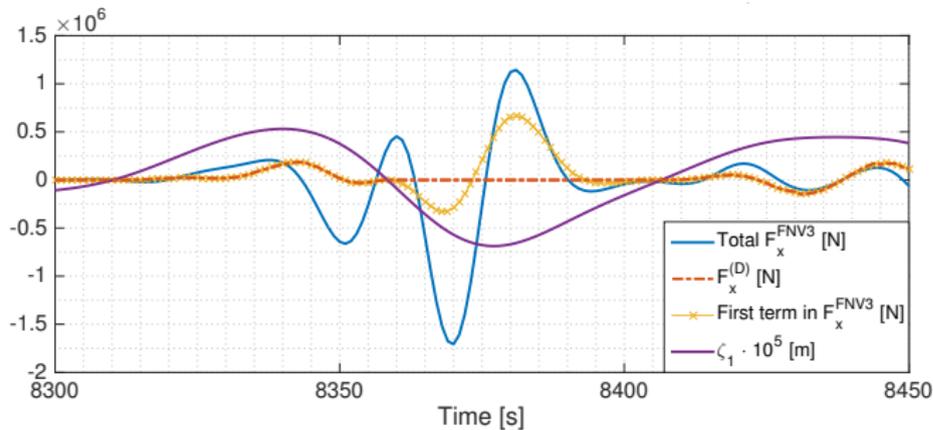


Kinematics models

| | $F_x / (0.5\pi\rho D^2)$ | $O(F_x)$ | $F_x \propto$ | Frequency |
|---|--|----------------------|---------------|---------------------|
| A | $\int_{-h}^0 u_{1,t}(z) dz$ | $\epsilon\delta^2$ | ζ_a | 1ω |
| B | $\int_{-h}^0 u_{2,t}(z) dz$ | $\epsilon^2\delta^2$ | ζ_a^2 | 2ω |
| C | $\int_0^{\max(0,\zeta_1)} u_{1,t}(0) dz$ | $\epsilon^2\delta^2$ | ζ_a^2 | 2ω |
| D | $\int_0^{\max(0,\zeta_1)} zu_{1,tz} dz$ | $\epsilon^3\delta^2$ | ζ_a^3 | $1\omega + 3\omega$ |
| E | $\int_0^{\max(0,\zeta_1)} u_{2,t}(0) dz$ | $\epsilon^3\delta^2$ | ζ_a^3 | $1\omega + 3\omega$ |
| F | $\int_{\max(0,\zeta_1)}^{\max(0,\zeta_1+\zeta_2)} u_{1,t}(0) dz$ | $\epsilon^3\delta^2$ | ζ_a^3 | $1\omega + 3\omega$ |
| G | $\int_{\max(0,\zeta_1)}^{\max(0,\zeta_1+\zeta_2)} zu_{1,tz}(0) dz$ | $\epsilon^4\delta^2$ | ζ_a^4 | 4ω |
| H | $\int_{\max(0,\zeta_1)}^{\max(0,\zeta_1+\zeta_2)} u_{2,t}(0) dz$ | $\epsilon^4\delta^2$ | ζ_a^4 | 4ω |



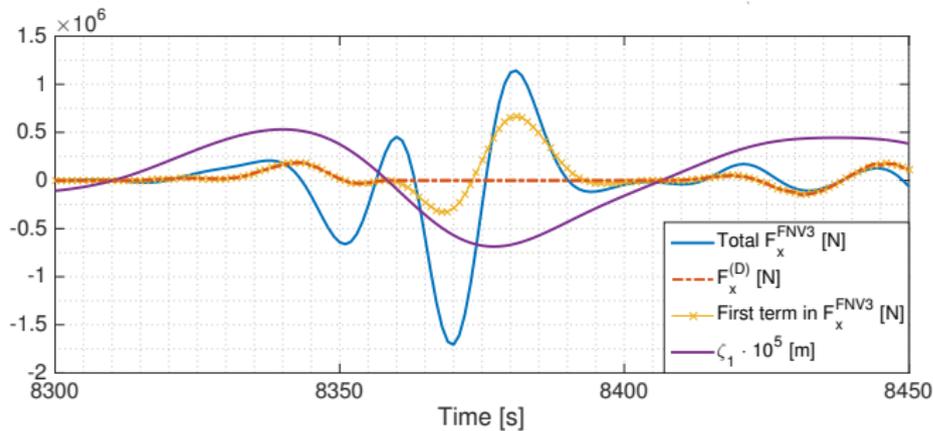
Third order forces



$$F_x^{FNV(3)} = \rho \pi r^2 \left[\zeta_1^2 u_{tz} + 2\zeta_1 w w_x + \zeta_1 u u_x - \frac{2}{g} \zeta_1 u_t w_t - \left(\frac{u_t}{g} \right) (u^2 + w^2) + \frac{\beta}{g} u^2 u_t \right]_{z=0}$$



Third order forces



$$F_x^{FNV(3)} = \rho \pi r^2 \left[\zeta_1^2 u_t z + 2\zeta_1 w w_x + \zeta_1 u u_x - \frac{2}{g} \zeta_1 u_t w_t - \left(\frac{u_t}{g} \right) (u^2 + w^2) + \frac{\beta}{g} u^2 u_t \right]_{z=0}$$



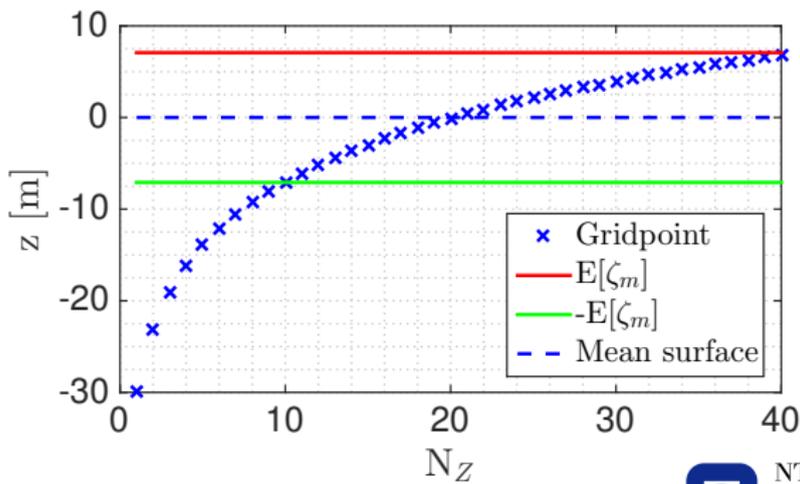
Kinematics models

| Notation | Fields | Description |
|----------|-----------|--|
| O1 | A | First order incident wave potential |
| O1D | A | First order incident wave potential w/diffraction |
| O2 | B+C | Second order incident wave potential and stretched first order potential |
| O3 | D+E+F+G+H | Third and fourth order force from stretched first and second order potential |
| O1P | A | First order diffraction pressure from panel code modeled as acceleration |
| O2P | B+C | Total second order diffraction force from panel code |
| FNV3 | N/A | Third order FNV ringing force based on first order incident potential |



Wave kinematics grid

- Logarithmically distributed in z -direction to increase accuracy in wave-zone
- 4.3 million points for 30 minute simulation with $dt = 0.1$ and $N_z = 40$, \rightarrow large files!



Sea-states

Chosen sea-states for Dogger Bank conditions. JONSWAP spectrum with peak parameter 3.3 is used.

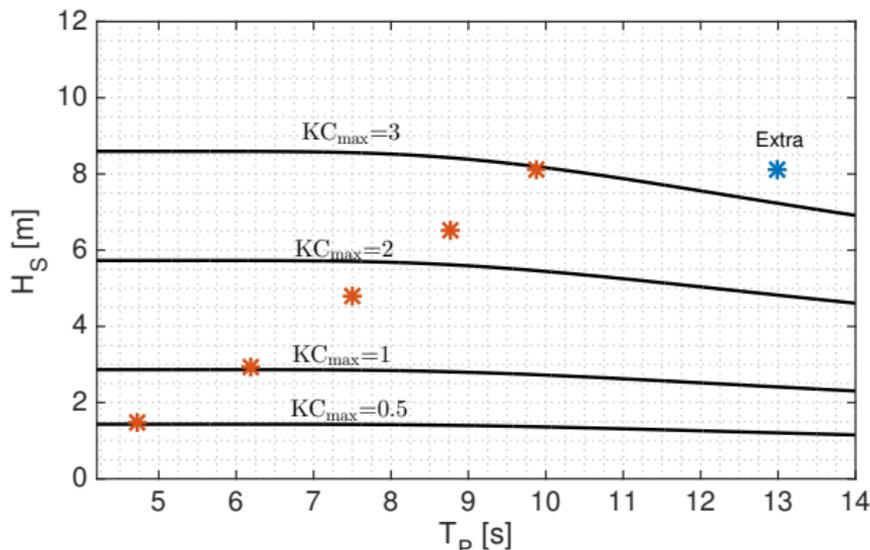


Figure: Sea-states with finite depth KC number for $h=30$ m.



Sea-states

Chosen sea-states for Dogger Bank conditions. JONSWAP spectrum with peak parameter 3.3 is used.

| No. | H_S [m] | T_P [s] | f_{H_S, T_P} [-] | KC_{max} [-] | $\pi D/\lambda$ [-] |
|-----|-----------|-----------|--------------------|----------------|---------------------|
| 1 | 1.46 | 4.72 | 0.1002 | 0.5 | 1.28 |
| 2 | 2.95 | 6.18 | 0.0314 | 1.0 | 0.75 |
| 3 | 4.79 | 7.50 | 0.0092 | 1.7 | 0.50 |
| 4 | 6.54 | 8.76 | 0.0016 | 2.3 | 0.37 |
| 5 | 8.13 | 9.88 | 0.0002 | 3.0 | 0.29 |
| 6* | 8.13 | 13.00 | 0.0000 | 3.5 | 0.17 |



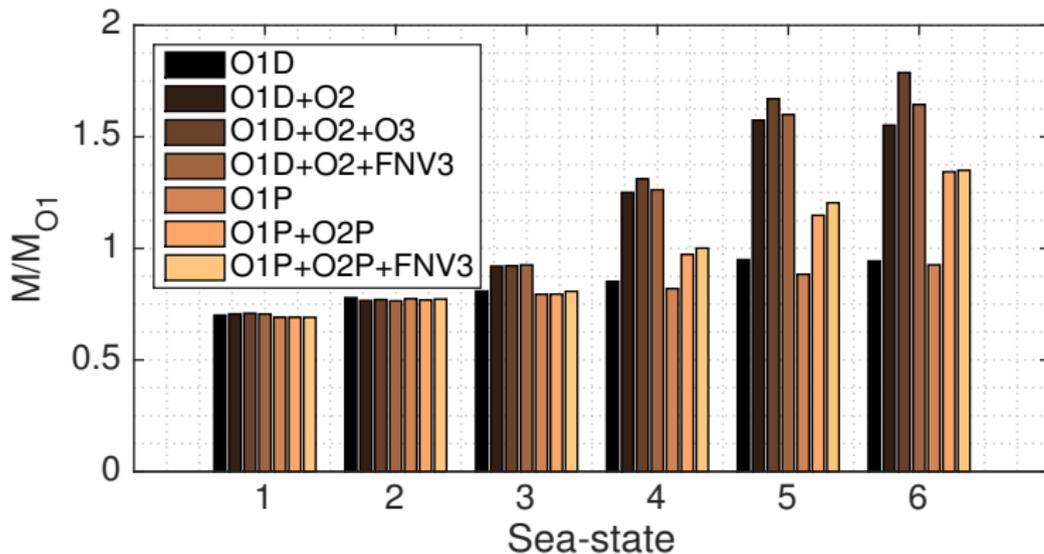
Results

- For each sea-state and hydrodynamic model, 3×30 minute simulations have been run without wind
- Average findings presented
- Small variances between the seeds



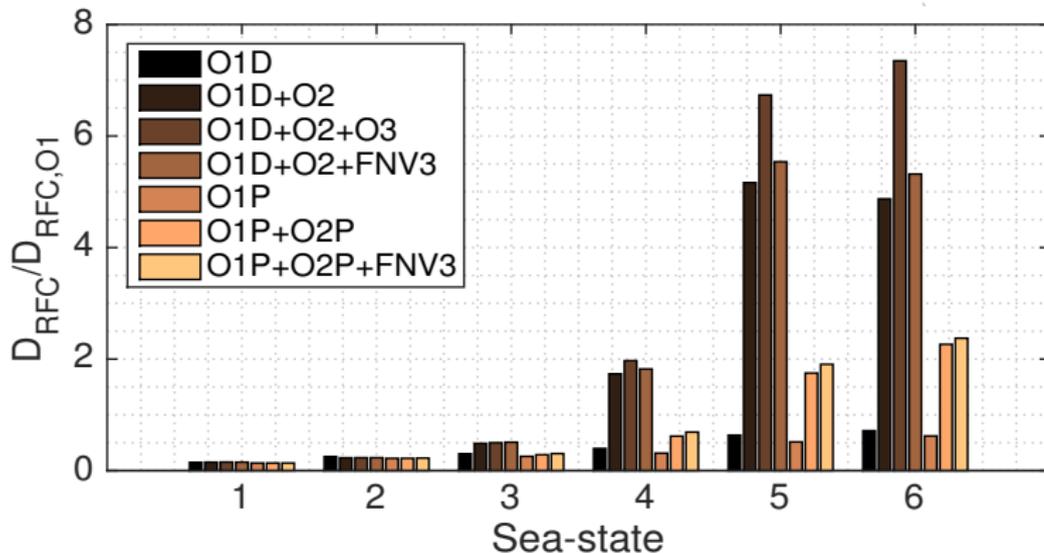
Results

Baseline maximum moment



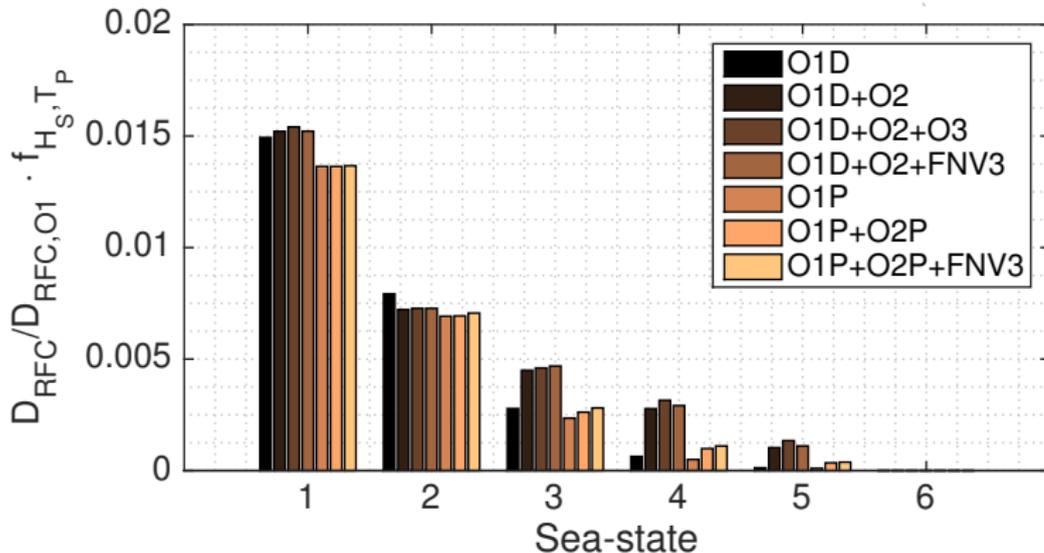
Results

Fatigue damage relative to first order incident wave



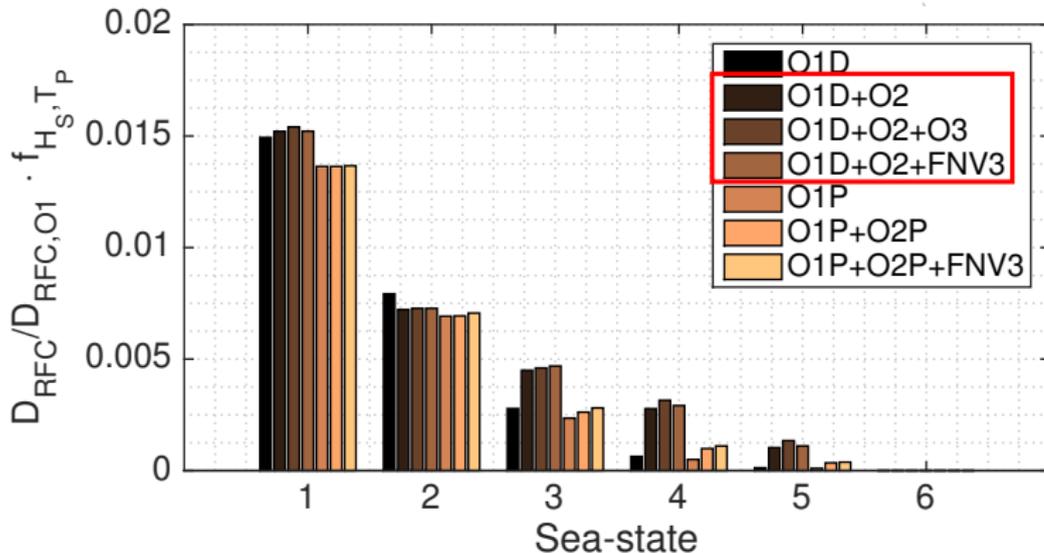
Results

Relative fatigue damage accounting for probability of occurrence



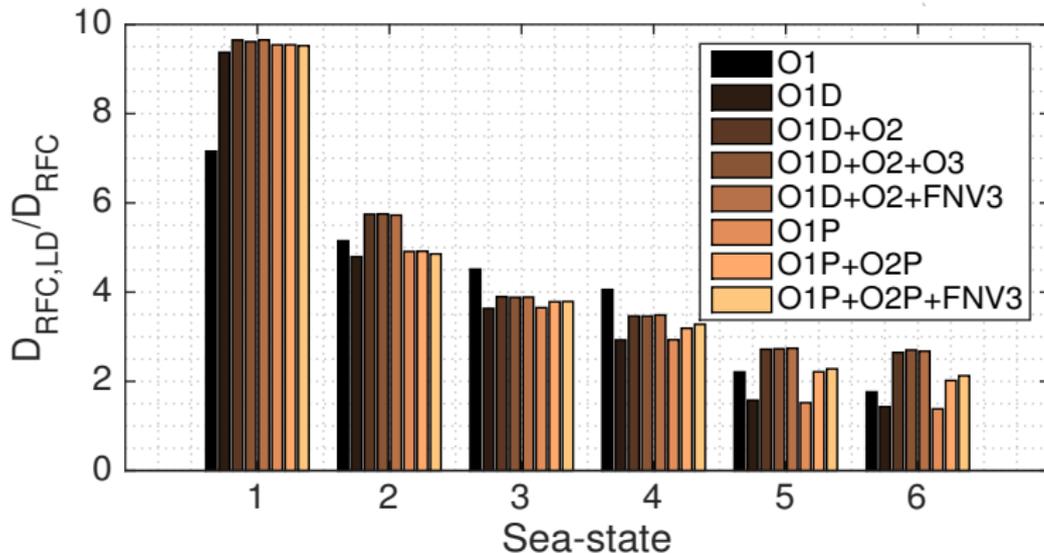
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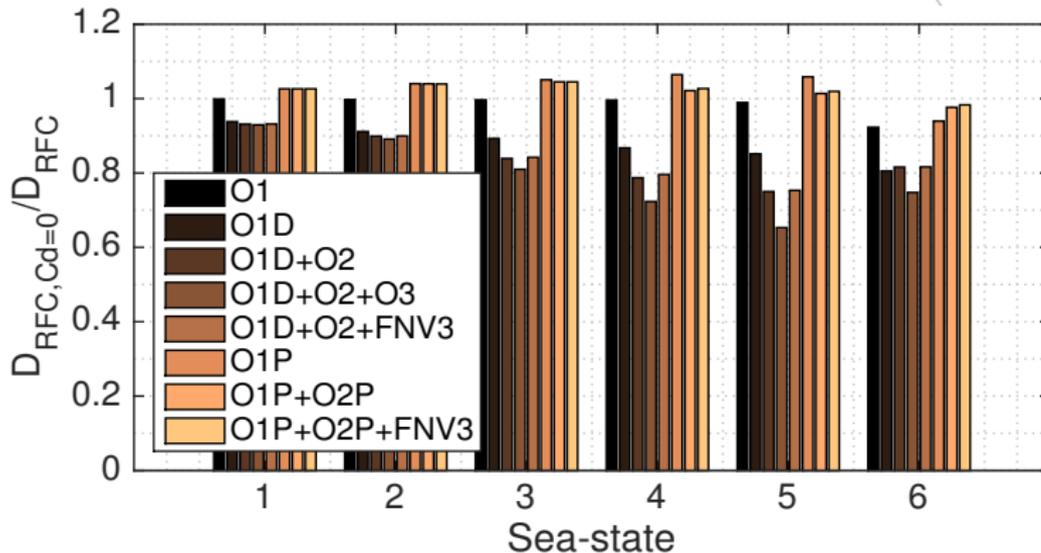
Results

Increased fatigue damage for lightly damped system: 3% → 1%



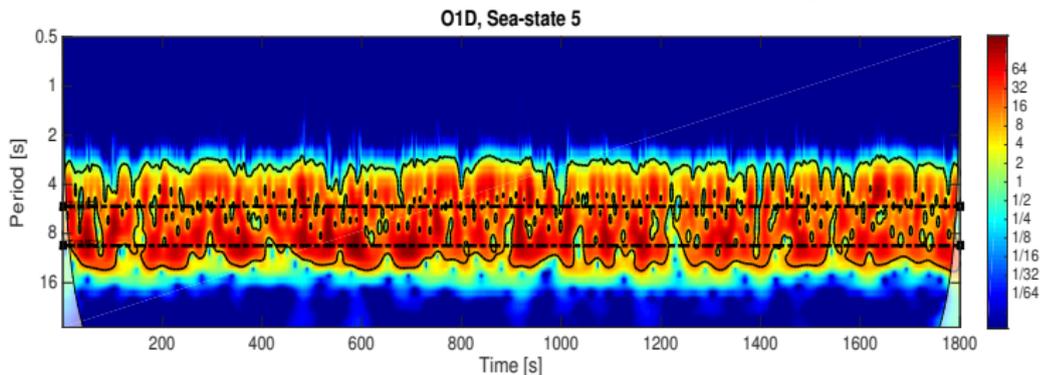
Results

Drag force contribution due to wave elevation and increasing KC-number



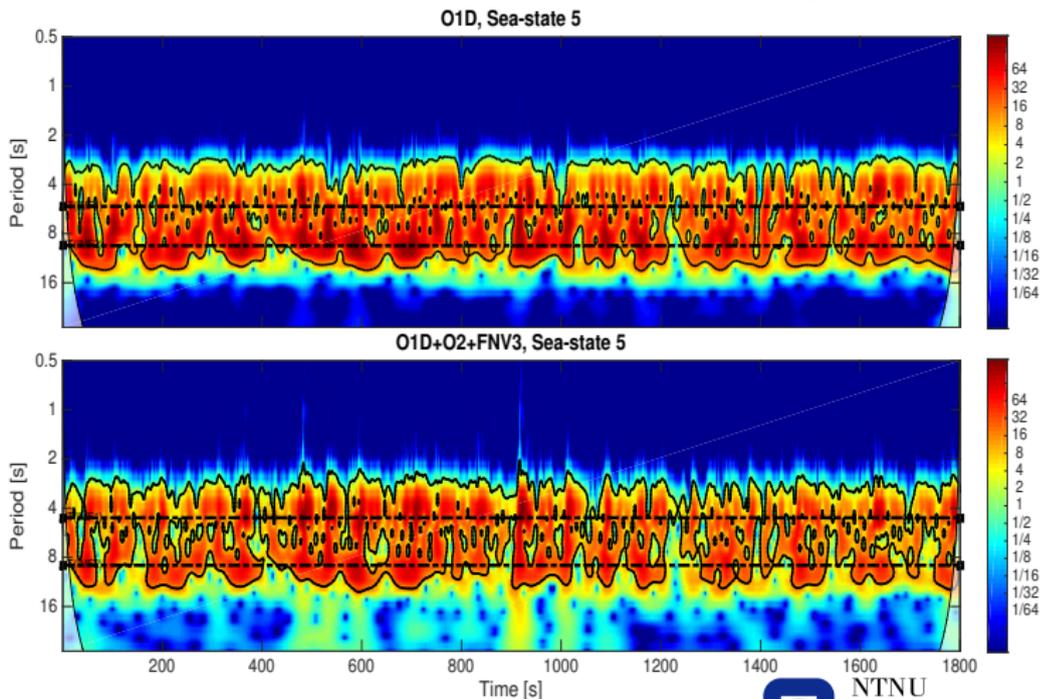
Time-frequency analysis

Wavelet analysis revealing most dominating oscillation periods.



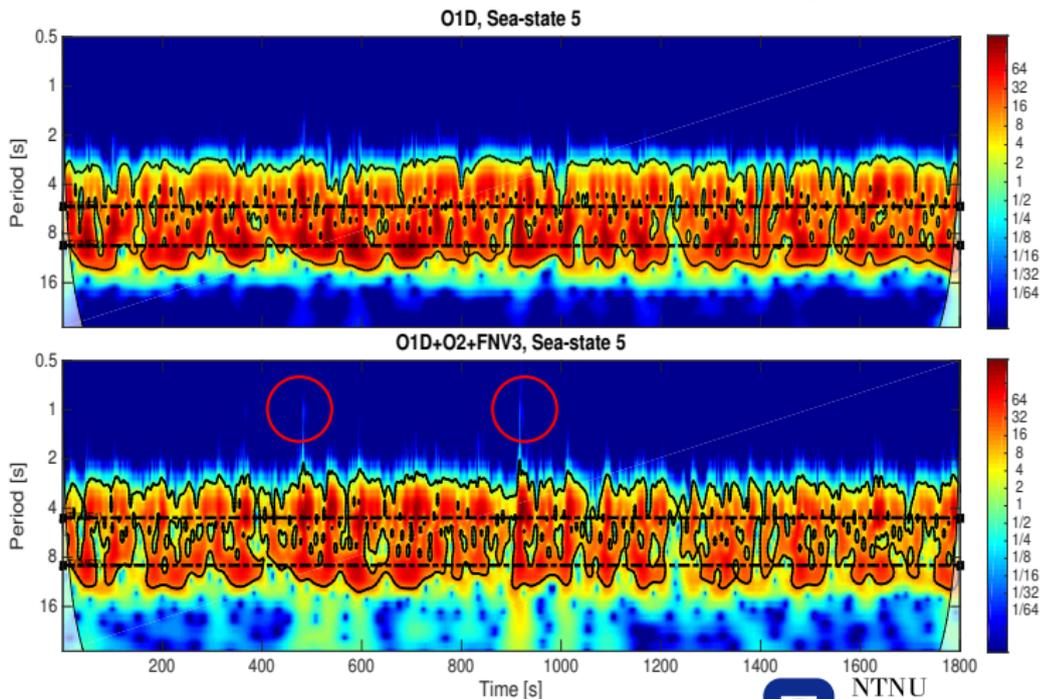
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Conclusions

- When $H_S > D/2$, significant contributions to fatigue damage from higher order loads are observed
- Higher order effects not important for smaller sea-states - overall small contributions when frequency of occurrence is accounted for
- Lower damping level results in more prominent contributions from higher order forces
- Drag forces still important when wave elevation is accounted for - need sensitivity study of C_D
- A Morison type loading for second order load seems to be predicting very large responses and fatigue damage for large sea-states - elevation important
- Important to include diffraction effects - both first and second order



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Thank you for your attention.
- Questions?



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