

Interface Tracking and Multi-Fluid Simulations of Bubbly Flows in Bubble Columns

Akio Tomiyama and Kosuke Hayashi

Kobe University

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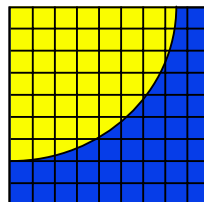
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Numerical Methods for Bubbly Flow Simulation

02/45

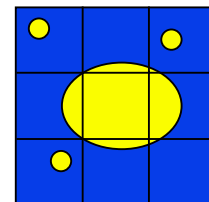
Interface Tracking Method (ITM)

$$d > 10\Delta x$$



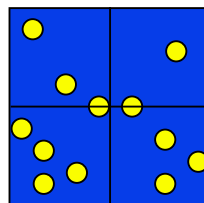
Bubble Tracking Method (BTM)

$$d \cong \Delta x$$



Two-Fluid Model (TFM)

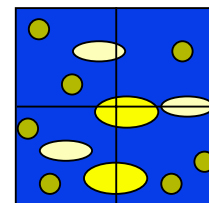
$$d < \Delta x$$



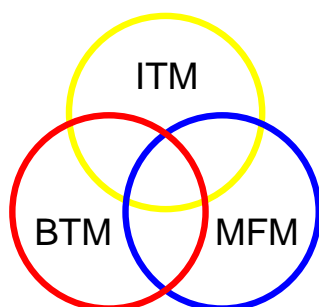
Multi-Fluid Model (MFM)

$$d_m < \Delta x$$

$(m = 1, \dots, N)$



Hybrid CMFD (Computational Multi-Fluid Dynamics)



HCMFD

BTM & MFM

ITM & MFM

ITM & BTM

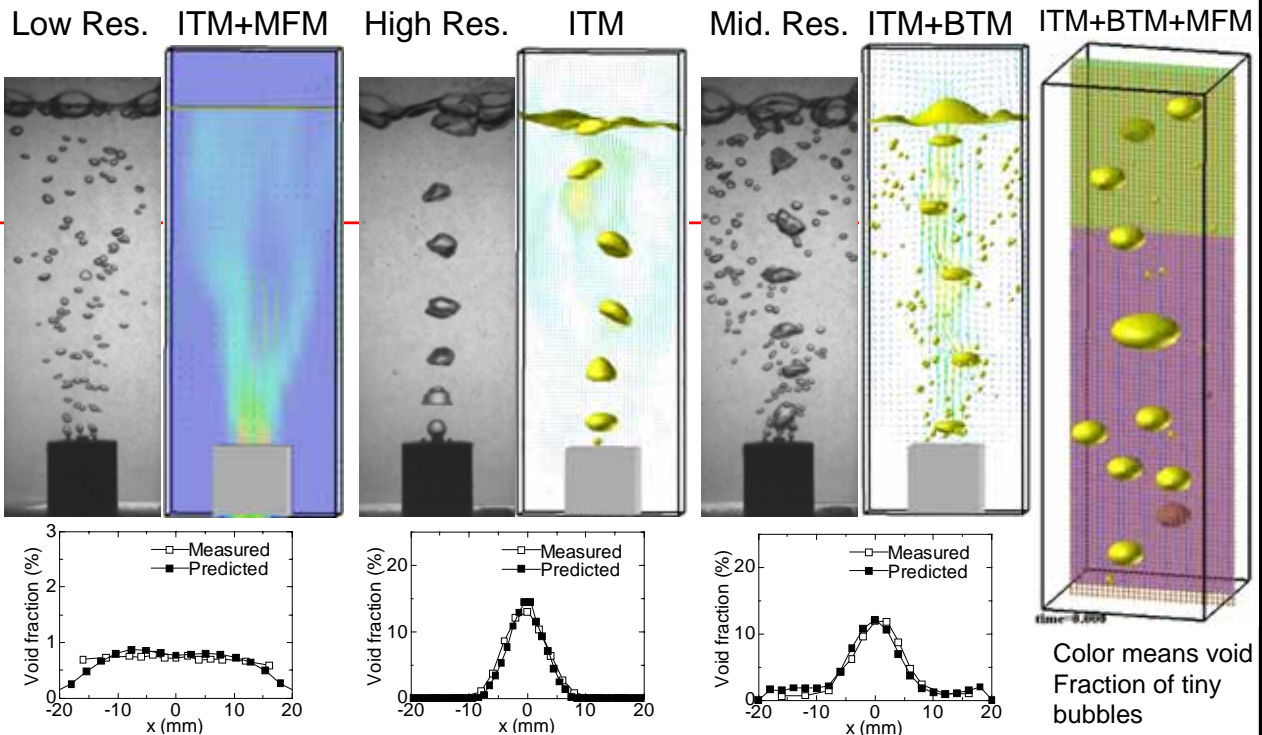
BTM

MFM

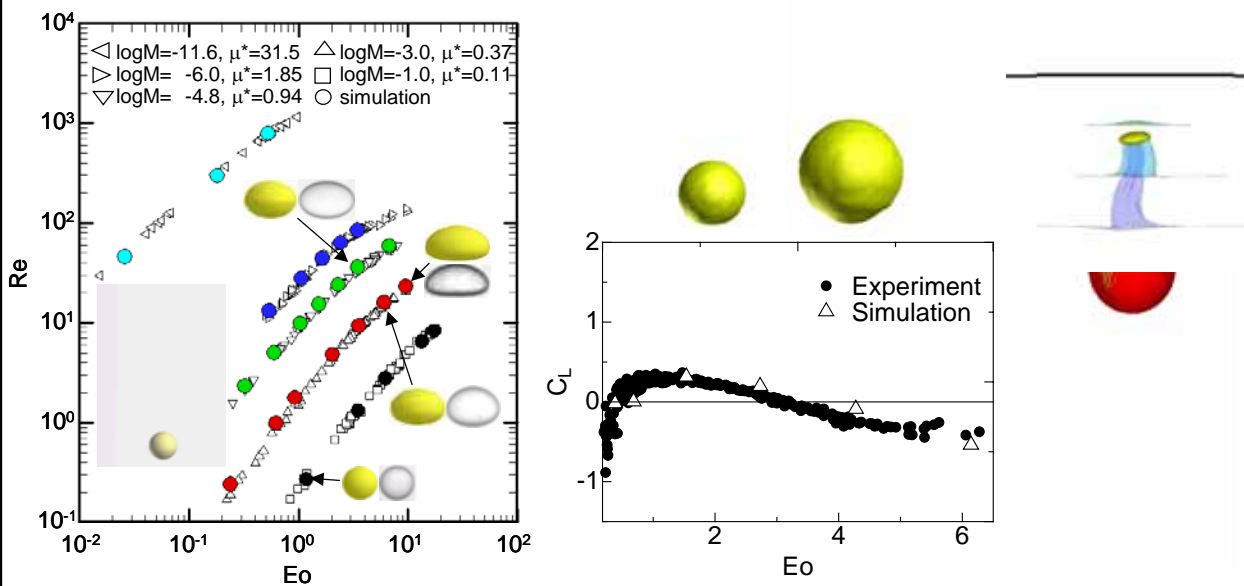
ITM

BTM

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Good predictions by using an appropriate combination!



When applying ITM to elementary phenomena in bubble columns, functions for simulating mass transfer, chemical reaction etc. are desired.

Interface Tracking Simulation of Mass Transfer through or onto Gas-Liquid Interface

Mass Transfer through bubbles at high Reynolds numbers

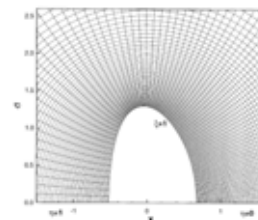
Adsorption and Desorption of Surfactant

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Interface Tracking Methods for Mass Transfer

Boundary-Fitted Coordinate Method

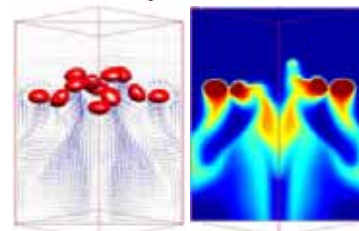
Ponoth & McLaughlin, 2000
Sugiyama et al., 2003



Ponoth & McLaughlin

Front Tracking Method

Koynov et al., 2005, 2006
Tryggvason et al., 2010
Darmana et al., 2007

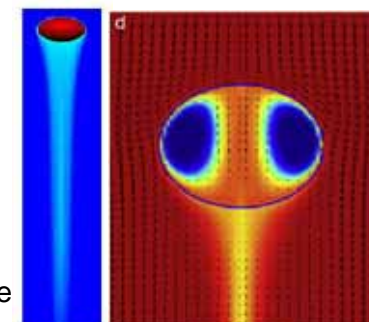


Darmana

Volume of Fluid and Level Set Methods

VOF: Davidson & Rudman, 2002
Bothe et al., 2004, 2010, 2011
Onea et al., 2006, 2009

LS: Yang & Mao, 2005
Wang et al., 2006



Bothe

Yang & Mao

Few methods can deal with volume change and bubbles at high Sc and high Re .

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Bubbles at High Re and Sc Numbers



CO₂ bubble in downward water flow

Schmidt number: 530

Reynolds number: 1500

Pipe diameter: 12.5mm

Initial bubble diameter: 29.0 mm

Initial Ingredients: CO₂ only

Final bubble diameter: 7.5 mm

Final ingredients: N₂ (80%), O₂ (20%)

Requirements for simulating bubbles at high Re and Sc numbers

- (1) Accurate conservation of species moles
- (2) Accurate evaluation of interfacial mass transfer
- (3) To capture a thin concentration boundary layer at high Sc

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Field Equations

Mass & Momentum Equations (One-Fluid Formulation)

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \mathbf{V} = 0$$

$$\frac{\partial \mathbf{V}}{\partial t} + \mathbf{V} \cdot \nabla \mathbf{V} = -\frac{1}{\rho} \nabla P + \frac{1}{\rho} \nabla \cdot \mu [\nabla \mathbf{V} + (\nabla \mathbf{V})^T] + \mathbf{g} + \frac{\sigma \kappa n \delta}{\rho}$$

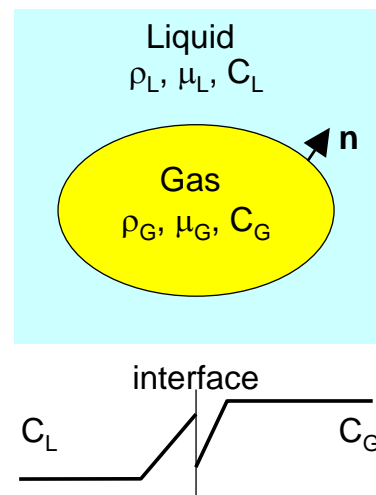
Conservation of Species Molar Concentration C_k

$$\frac{\partial C_k}{\partial t} + \mathbf{V} \cdot \nabla C_k = \nabla \cdot D_k \nabla C_k \quad (k = G, L)$$

Jump Conditions for C_k

$$C_G = m C_L \quad \text{m: distribution coefficient (Henry's law)}$$

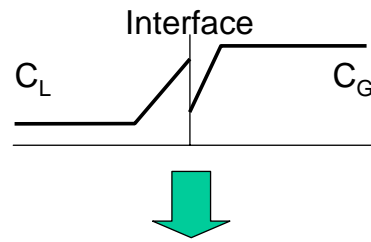
$$\mathbf{j} \cdot \mathbf{n} \Big|_{int} = -D_G \frac{\partial C_G}{\partial n} \Big|_{int} = -D_L \frac{\partial C_L}{\partial n} \Big|_{int}$$



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Variable Transformation

$$\Phi = \begin{cases} mC_L & x \in L \\ C_G & x \in G \end{cases}$$

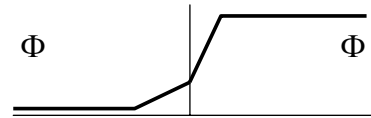


Single-Variable Formulation

$$\frac{\partial \Phi}{\partial t} + \mathbf{V} \cdot \nabla \Phi = \nabla \cdot \hat{D} \nabla \Phi$$

$$\Phi|_{Gint} = \Phi|_{Lint}$$

$$-D_G \frac{\partial \Phi}{\partial n} \Big|_{int} = -(D_L / m) \frac{\partial \Phi}{\partial n} \Big|_{int}$$



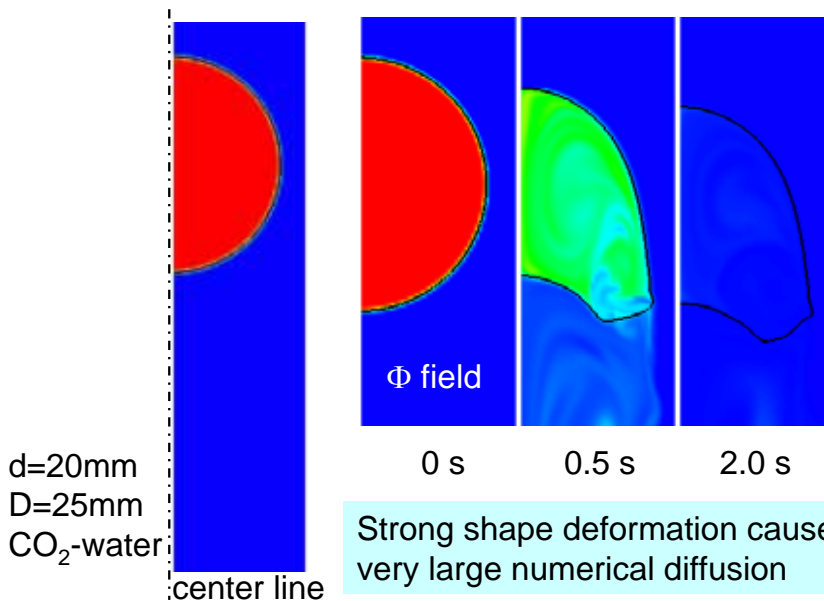
Solutions, C_L & C_G , are obtained by solving only the single equation for Φ .

This formulation has been often adopted.

(Bothe et al., 2004; Yang & Mao, 2005; Onea et al., 2006; 2009; Francois & Carlson, 2010)

Defect of Single-Variable Formulation

Example: 20mm CO₂ Bubble without Any Mass Transfer (No Volume Change)
 Yan & Mao's method (Level set, 5th order WENO, 3rd order TVD, RK)



CFD2011: Fleckenstein & Bothe stopped using this for simulating high Re bubbles

Two-Variable Formulation for Moles, M_G and M_L

$$\frac{\partial M_k}{\partial t} + \int_{S_k} C_k \mathbf{V} \cdot d\mathbf{S} = \int_{S_k} D_k \nabla C_k \cdot d\mathbf{S}$$

$$C_k = \frac{M_k}{\Theta_k}$$

M_k : mole in phase k [mol]

C_k : molar concentration [mol/m³]

Θ_k : volume of phase k [m³]

Accurate Advection of M_k

$$M_k^{n+1} = M_k^n - \Delta t \int_{S_k} C_k \mathbf{V} \cdot d\mathbf{S}$$

Transferred mole = Transferred fluid volume times C

Transferred volume: Volume Tracking Method based on

NSS (Non-uniform Subcell Scheme), Hayashi et al., 2006)

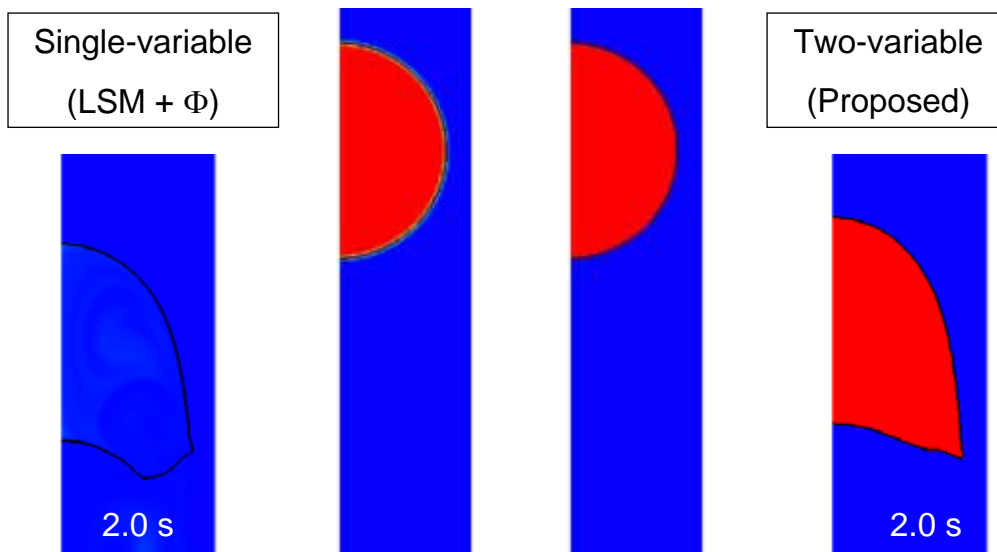
+ EI-LE scheme (Aulisa et al., 2003)



Accurate Conservation of Fluid Volume & Species Moles

Comparison between single- and two-variable formulations

20 mm CO₂ bubble in water (pipe diameter = 25 mm)



Two-variable formulation gives no numerical diffusion

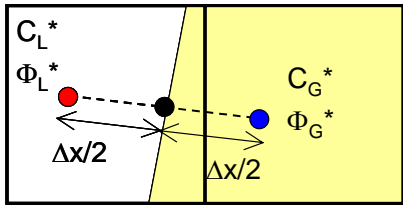
Trend: Use of a Solution based on Boundary Layer Approximation

Tryggvason et al., 2010, Bothe et al., 2010

$$C_L(\xi, \eta) = \left(\frac{C_{Gi+1j}}{m} \right) \left[1 - \operatorname{erf} \left(\frac{\xi}{\sqrt{4D_L\eta/V_B}} \right) \right] \rightarrow \mathbf{j \cdot n|_{int} = -D_L \left[\frac{\partial C_L}{\partial n} \right]_{int}}$$

This may not hold for actual boundary layers (wavy, flow separation, ..).

Evaluation of Interfacial Mass Flux using Single-Variable & Harmonic Mean



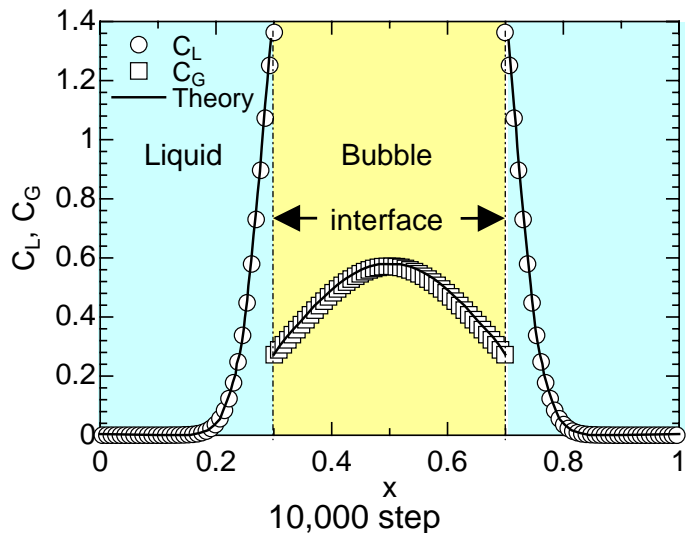
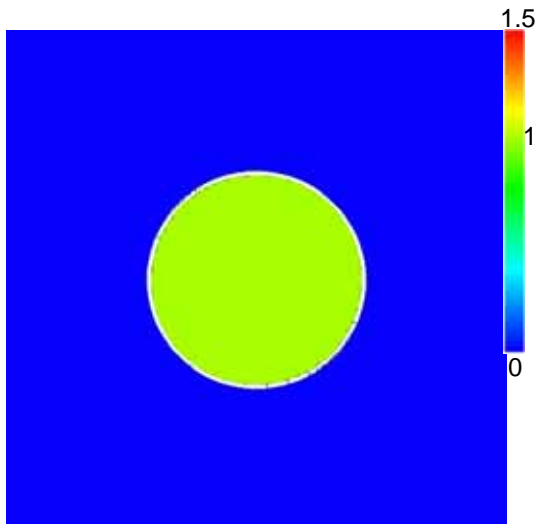
$$\mathbf{j \cdot n|_{int} = -D_G \frac{\Phi_{int} - \Phi_G^*}{\Delta x/2} = -(D_L/m) \frac{\Phi_L^* - \Phi_{int}}{\Delta x/2}}$$

Elimination of Φ_{int} \rightarrow $\mathbf{j \cdot n|_{int} = -\hat{D} \frac{\Phi_L^* - \Phi_G^*}{\Delta x}}$

$$\hat{D} = \frac{2\hat{D}_G\hat{D}_L}{\hat{D}_G + \hat{D}_L} \quad \hat{D}_k = \begin{cases} D_L/m & \text{for } k = L \\ D_G & \text{for } k = G \end{cases}$$

Validation: Mass Diffusion with Interface Jump

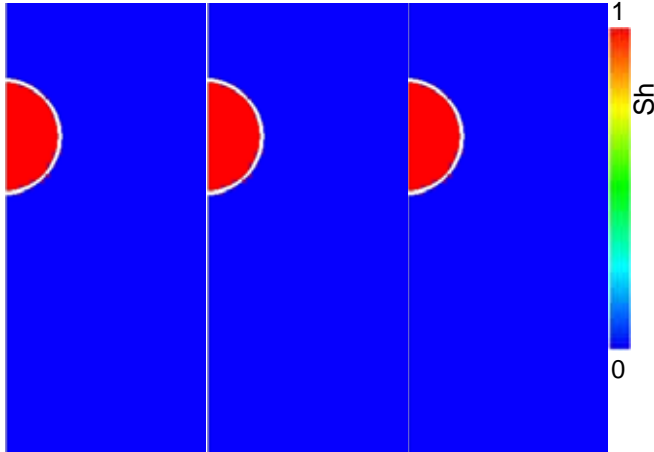
$C_G=1$ & $C_L=0$ at $t=0$
 $m = 0.2$ ($C_G=mC_L$, $C_L > C_G$ at interface)
 $D_G/D_L = 10$
 No volume change
 $(x, y) = (128, 128)$ cells



Interfacial mass transfer is accurately computed.

Case 1: $Eo=1, M=1 \times 10^{-4}$
 Case 2: $Eo=40, M=9.2 \times 10^{-3}$
 Case 3: $Eo=3.1, M=5 \times 10^{-7}$
Sc=1

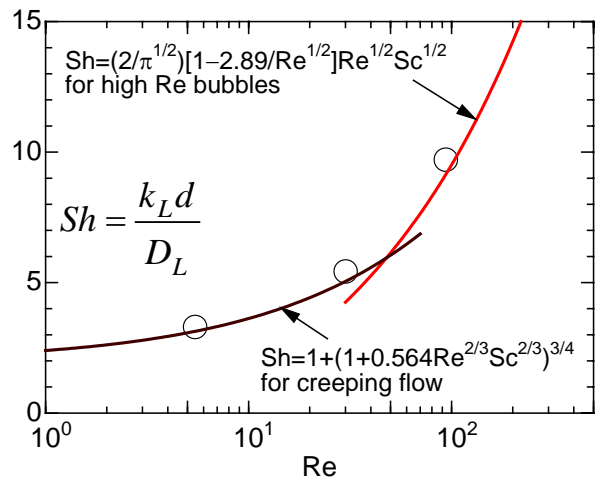
$m = 1$ (no jump in C), $C_G = \text{constant}$
 No volume change
 (R, H) = (40, 160 cells)



Case 1 Re=5.5 Case 2 Re=31 Case 3 Re=97

Mass transfer coefficient k_L

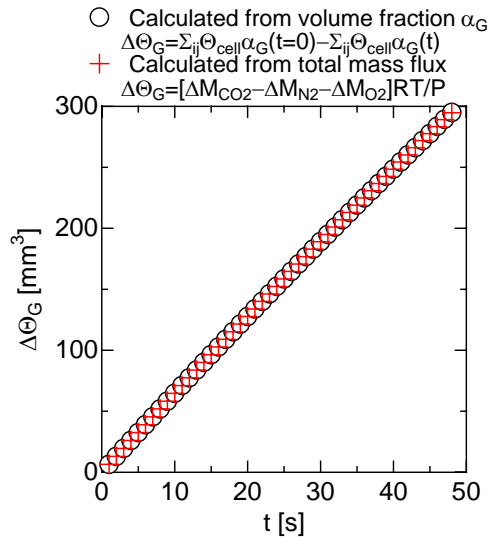
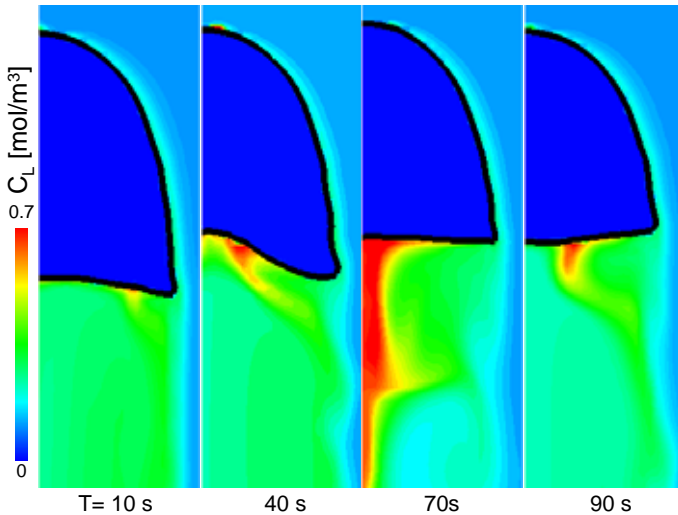
$$k_L = \frac{1}{A\Delta t} \int_{V \in L} (C_L^{n+1} - C_L^n) dV$$



Mass transfer from a bubble at low Re and low Sc is well predicted

Accuracy of Computation of Volume Change

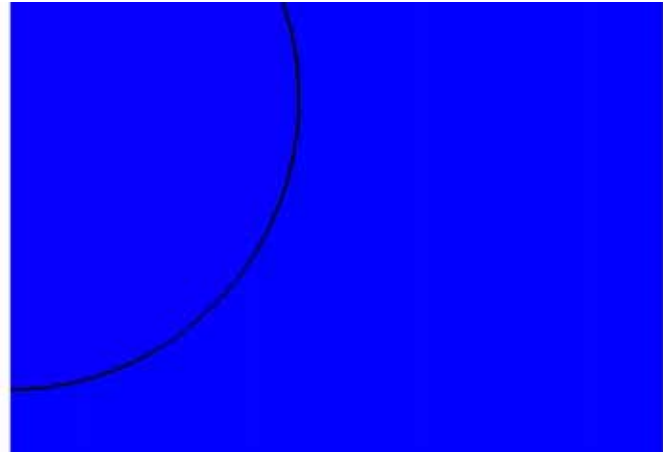
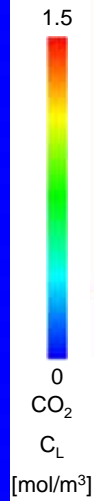
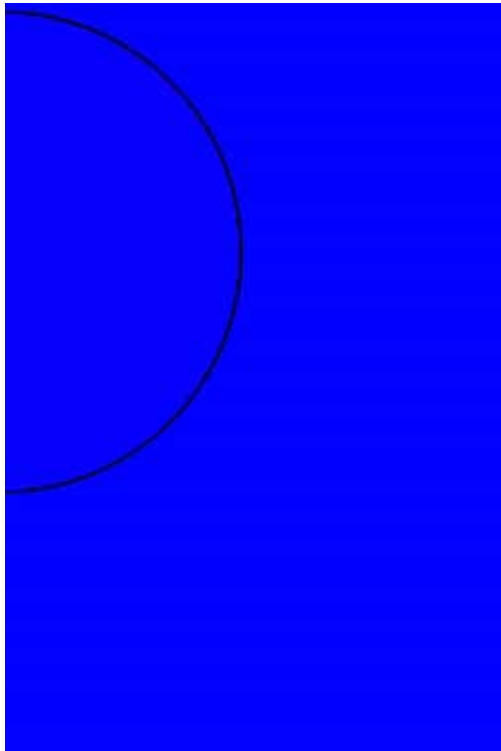
Bubble: $d=15 \text{ mm}$, CO_2 100% at $t = 0 \text{ sec}$
 Pipe: $D= 18.6 \text{ mm}$
 $Re = 2700, Sc = 530$
 Water: equilibrium concentration of N_2 & O_2
 Volume Change : Considered



The volume change due to mass transfer is accurately computed.

Simulation of 8 mm CO₂ Bubble in Stagnant Water

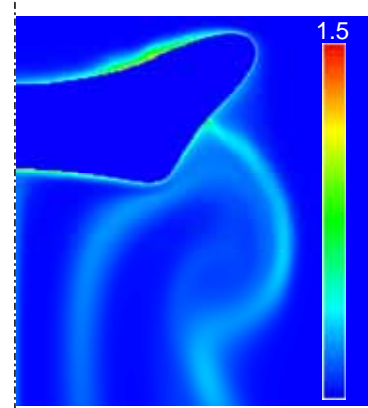
Re=1200, Sc=530, D = 18.2 mm



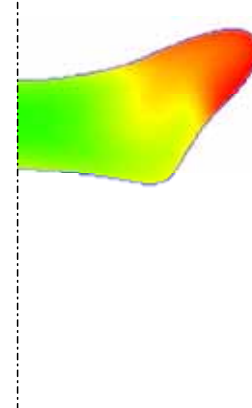
Interface oscillation plays an important role in mass transfer process.
Boundary layer approximation would not hold in separated regions

85 cells for the bubble radius

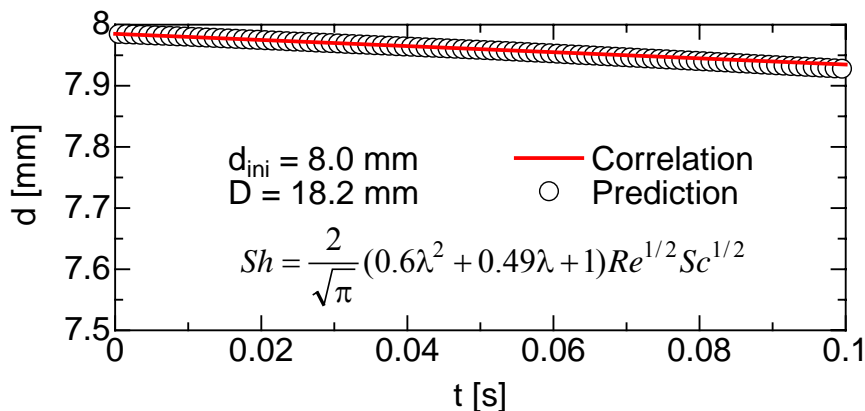
Comparison with Measured Bubble Dissolution Process



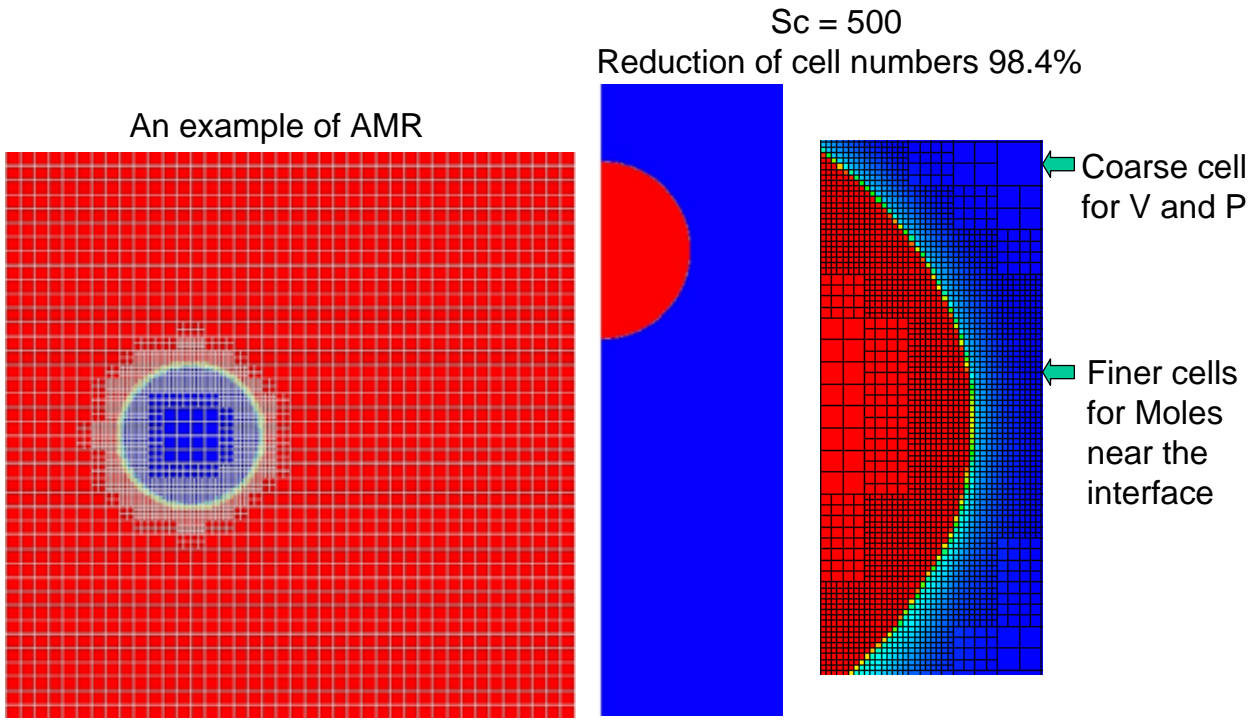
CL C_L [mol/m³] CO₂



CL C_G [mol/m³] N₂

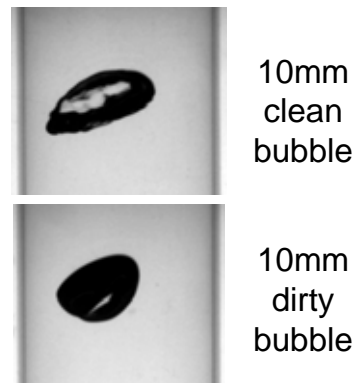
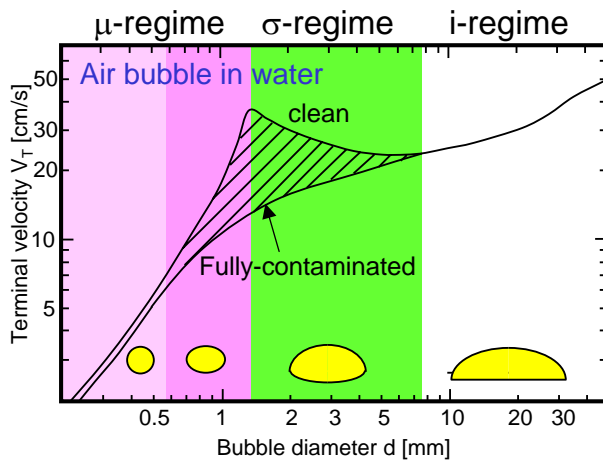


8mm CO₂ Bubble
Cell size = 47 μm
(170 cells for d)



High Resolution only for Concentration Boundary Layer

1. Reduction of surface tension σ
2. Marangoni effect (tangential force)
3. Increase in damping coefficient of capillary wave



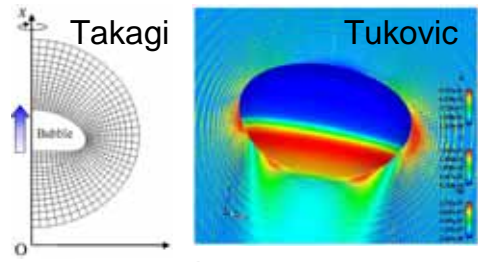
Decrease in V_T only in some bubble diameter range
Simulations have been performed only for μ -regime

No effects of surfactant on V_T at High Eo_D (Well-known fact)
Increase in V_T at low Eo_D (Almatroushi & Borhan, 2004)

$$Eo_D = \frac{\Delta\rho g D^2}{\sigma}$$

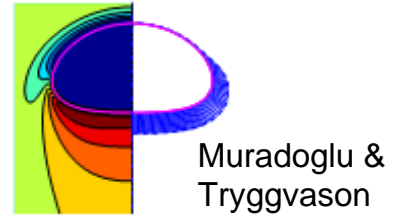
Boundary-Fitted Coordinate Method

- Cuenot et al., 1997 (spherical bubble)
- Takagi et al., 2003 (spherical bubble)
- Tukovic & Jasak, 2008 (deformed bubble)



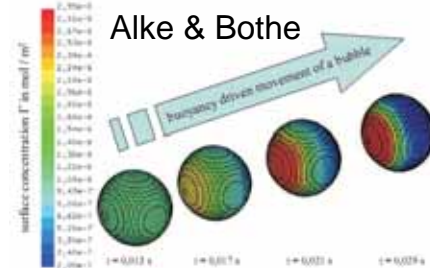
Front Tracking Method

- Muradoglu & Tryggvason, 2008 (low Re spheroidal bubble)



Volume of Fluid and Level Set Methods

- VOF: Renardy et al., 2002
James & Lowengrub, 2004
Alke & Bothe, 2009
- LS: Xu & Shao, 2003
Xu et al., 2006



Few methods can simulate highly distorted bubbles at high Re numbers

Field equations & Numerical Methods

Mass & Momentum Equations (Projection Method)

$$\nabla \cdot \mathbf{V} = 0$$

$$\frac{\partial \mathbf{V}}{\partial t} + \mathbf{V} \cdot \nabla \mathbf{V} = -\frac{1}{\rho} \nabla P + \frac{1}{\rho} \nabla \cdot \mu [\nabla \mathbf{V} + (\nabla \mathbf{V})^T] + \mathbf{g} + \frac{1}{\rho} [\sigma \kappa \mathbf{n} + \nabla_S \sigma] \delta$$

Ghost fluid method CSF model

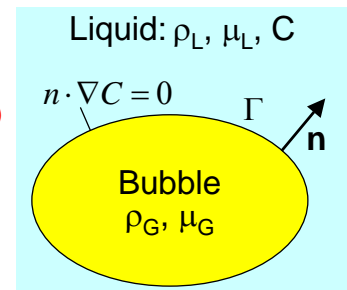
Bulk & Surface Concentrations of Surfactant

$$\frac{\partial C}{\partial t} + \mathbf{V} \cdot \nabla C = \nabla \cdot D_B \nabla C + \dot{S}_\Gamma \delta$$

Two step method (Muradoglu & Tryggvason)

$$\frac{\partial \Gamma}{\partial t} + \nabla_S \cdot \Gamma \mathbf{V}_S = \nabla_S \cdot D_S \nabla_S \Gamma - \dot{S}_\Gamma$$

Extrapolation method (Xu & Zhao)



Adsorption-Desorption Kinetics (Langmuir model)

$$\dot{S}_\Gamma = -\mathbf{n} \cdot (D_C \nabla C|_{int}) = k[C_S(\Gamma_{max} - \Gamma) - \beta \Gamma]$$

k: adsorption constant
β: desorption constant (Frumkin & Levich)

Advection of Interface

$$\frac{\partial \phi}{\partial t} + \mathbf{V} \cdot \nabla \phi = 0$$

Level set method (Susuman)

Surface Tension

$$\sigma = \sigma_0 + RT\Gamma_{max} \ln \left(1 - \frac{\Gamma}{\Gamma_{max}} \right)$$

Small Air Bubble in Contaminated Water (μ -Regime)

23/45

Fluids: air and water

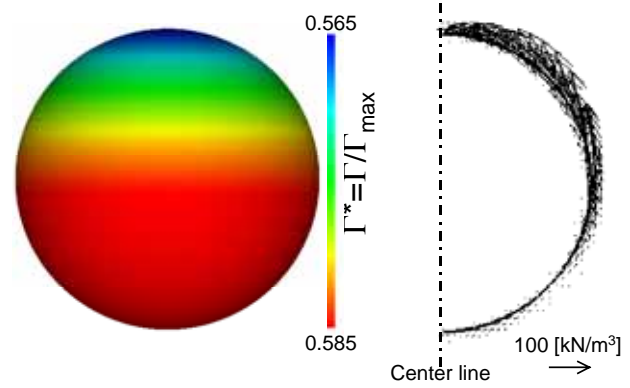
Surfactant: 1-pentanol

$k=5.08 \text{ m}^3/\text{mol}\cdot\text{s}$

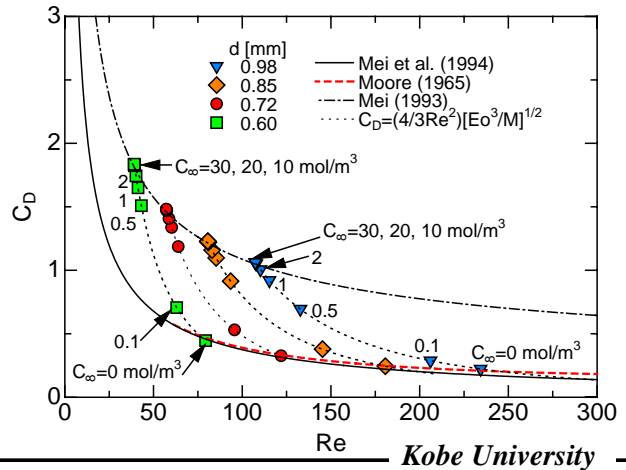
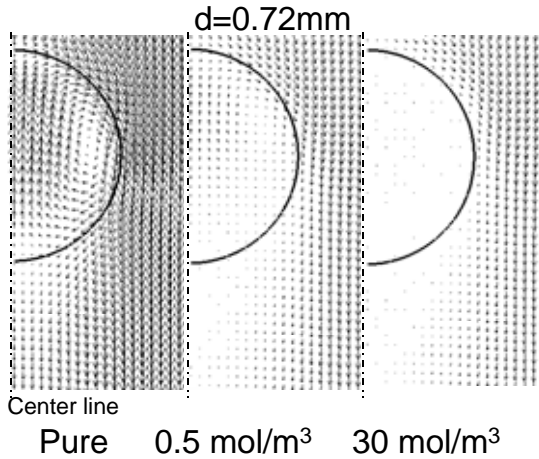
$\beta=21.7 \text{ mol/m}$

$\Gamma_{\text{max}}=5.9 \times 10^{-6} \text{ mol/m}^2$

Bubble diameter: 0.6-0.98 mm



$d=0.72\text{mm}$, $C=30\text{mol/m}^3$

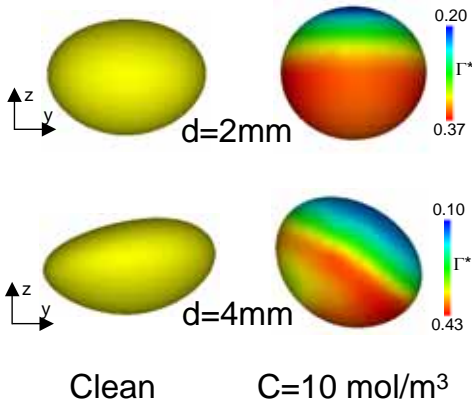


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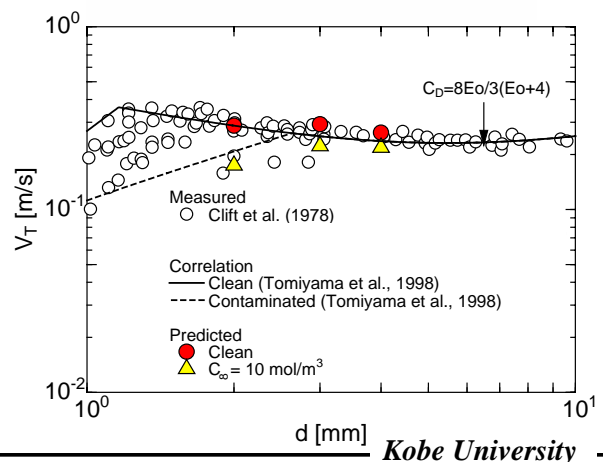
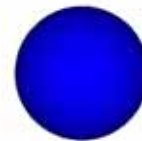
Distorted Air Bubble in Contaminated Water (σ -Regime)

24/45

4 mm air bubble in clean water



4 mm air bubble in contaminated water ($C=10\text{mol/m}^3$)

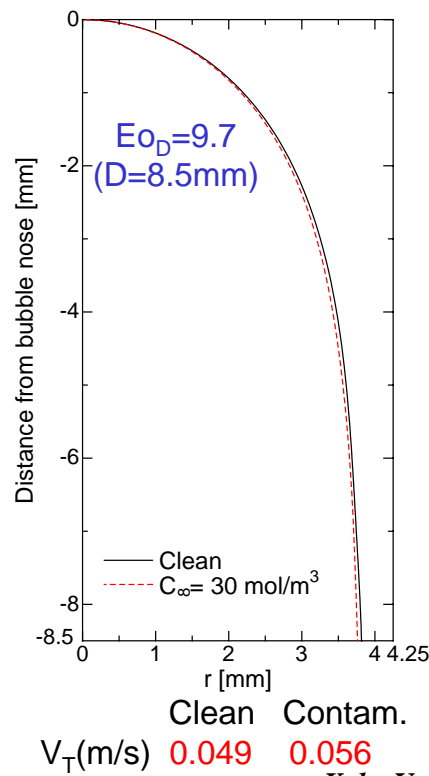
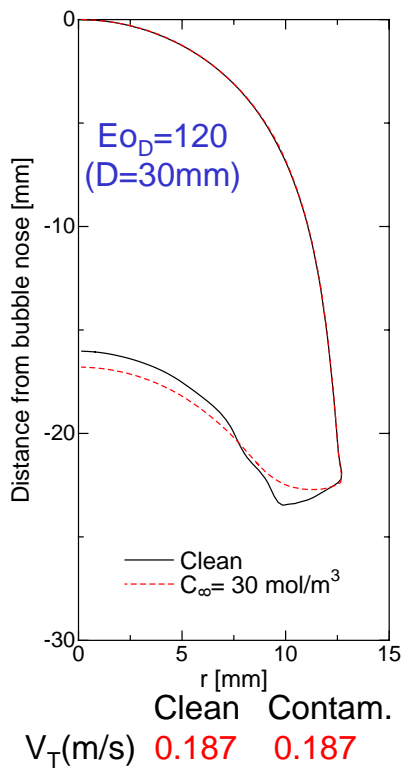


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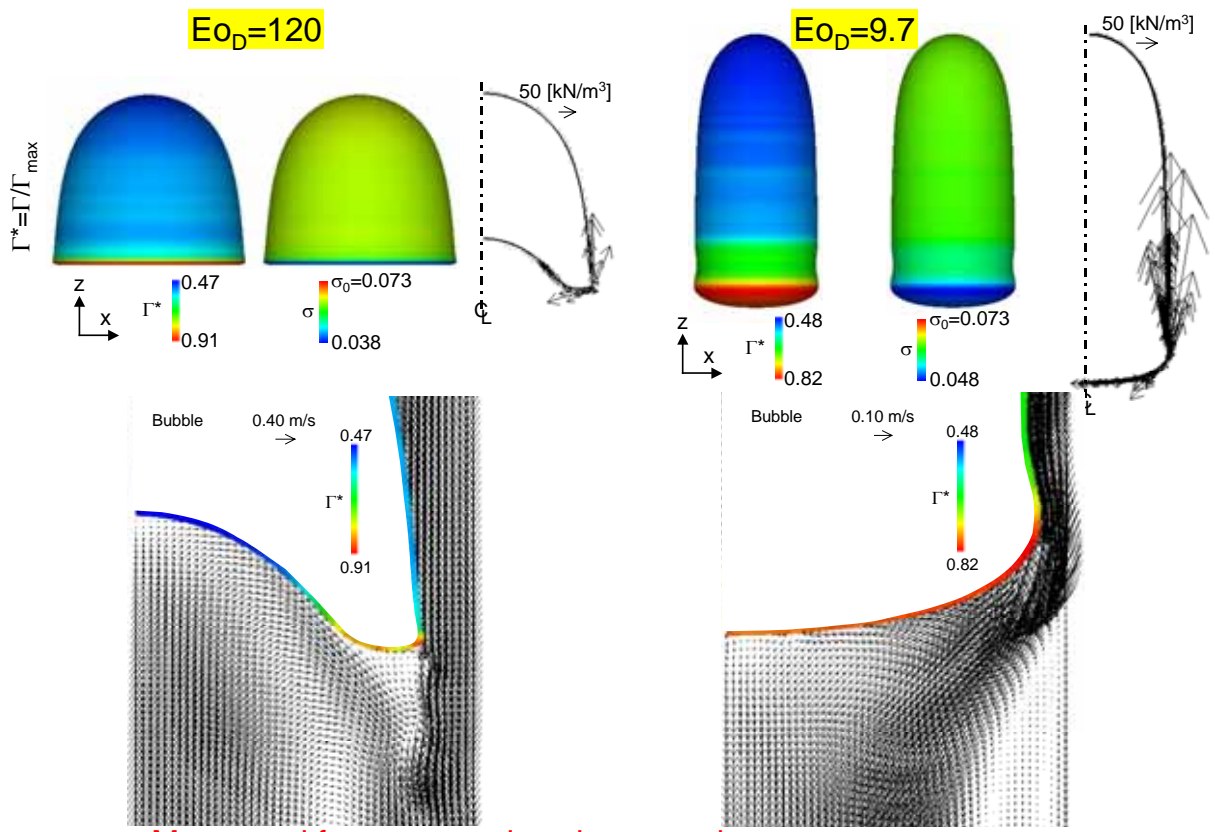
Taylor Bubbles in Air-Water System (σ and i-Regimes)

No effects of surfactant on V_T at High Eo_D (Well-known fact)
 Increase in V_T at low Eo_D (Almatroushi & Borhan, 2004)

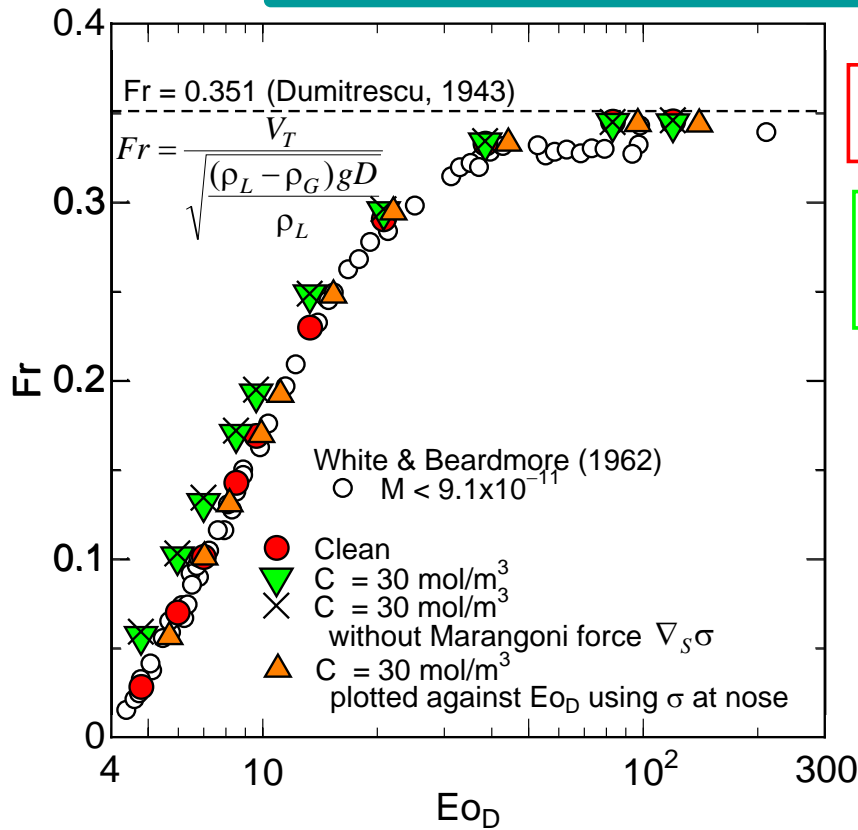
$$Eo_D = \frac{\Delta\rho g D^2}{\sigma}$$



Distributions of Surfactant, Marangoni Force & Velocity



Marangoni force acts only at bottom edge.
 Surfactant distribution strongly depends on velocity profile.

Cause of V_T Increase at Low Eo_D 

Good predictions for the clean system

Contaminated case high Fr at low Eo same Fr at high Eo

Negligible effect of Marangoni force

$$Eo_D = \frac{\Delta \rho g D^2}{\sigma_{nose}}$$

Surface tension at bubble nose is the key parameter

The nose surface tension determines the rise velocity of a low M Taylor Bubble

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Multi-Fluid Simulation of Bubbly Flow in a Bubble Column

Validation of Bubble Coalescence and Breakup Model

Effects of Fine Particles on Bubble Coalescence

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Dispersed Phases ($m = 0, \dots, N$)

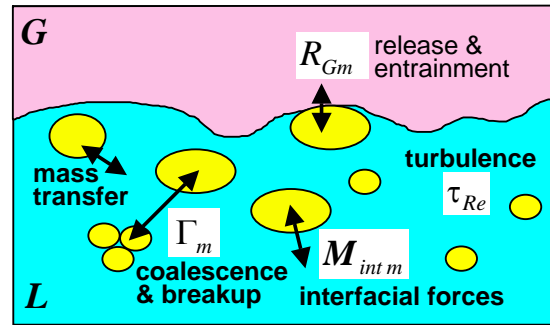
$$\frac{\partial n_m}{\partial t} + \nabla \cdot n_m \mathbf{V}_m = \sum_{m' \neq m} (\Gamma_{m' \rightarrow m} - \Gamma_{m \rightarrow m'}) - R_m$$

$$\alpha_m \rho_G \frac{D\mathbf{V}_m}{Dt} = -\nabla P + \mathbf{g} - (\mathbf{M}_{int m} + \mathbf{M}_{\Gamma_m} + \mathbf{M}_{R_m})$$

Continuous Gas & Liquid Phases

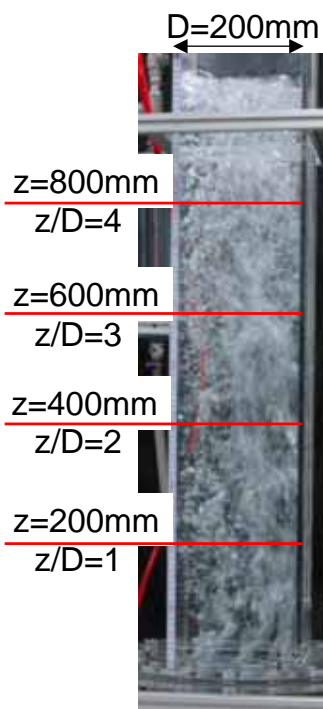
$$\frac{\partial \alpha_G}{\partial t} + \nabla \cdot (\alpha_G \mathbf{V}_C) = \sum_{m=0}^N R_m \quad \frac{\partial \alpha_L}{\partial t} + \nabla \cdot (\alpha_L \mathbf{V}_C) = 0$$

$$\alpha_C \rho_C \frac{D\mathbf{V}_C}{Dt} = -\nabla P + \mathbf{F}_S + \alpha_C \rho_C \mathbf{g} + \sum_{m=0}^N (\mathbf{M}_{int m} + \mathbf{M}_{R_m}) + \nabla \cdot \alpha_C (\boldsymbol{\tau} + \boldsymbol{\tau}_{Re})$$



- Experiments on single bubbles and drops in infinite liquids or in pipes
 - ➔ Drag and lift correlations for various bubbles and drops
- Experiments on mass transfer from single bubbles in vertical pipes
 - ➔ Sherwood number correlations
- Experiments on turbulent poly-dispersed bubbly pipe flows
 - ➔ Validation of two-phase turbulence model
- Experiments on poly-dispersed bubbly flow in a bubble column
 - ➔ Validation of bubble coalescence and breakup models

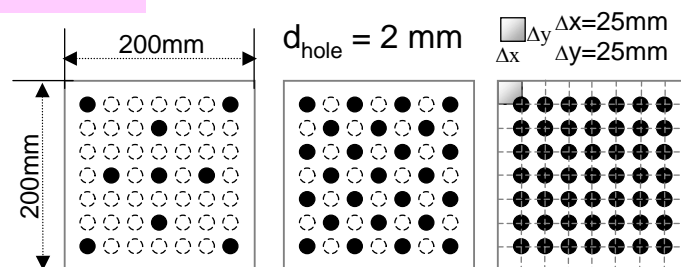
Experimental Setup & Conditions



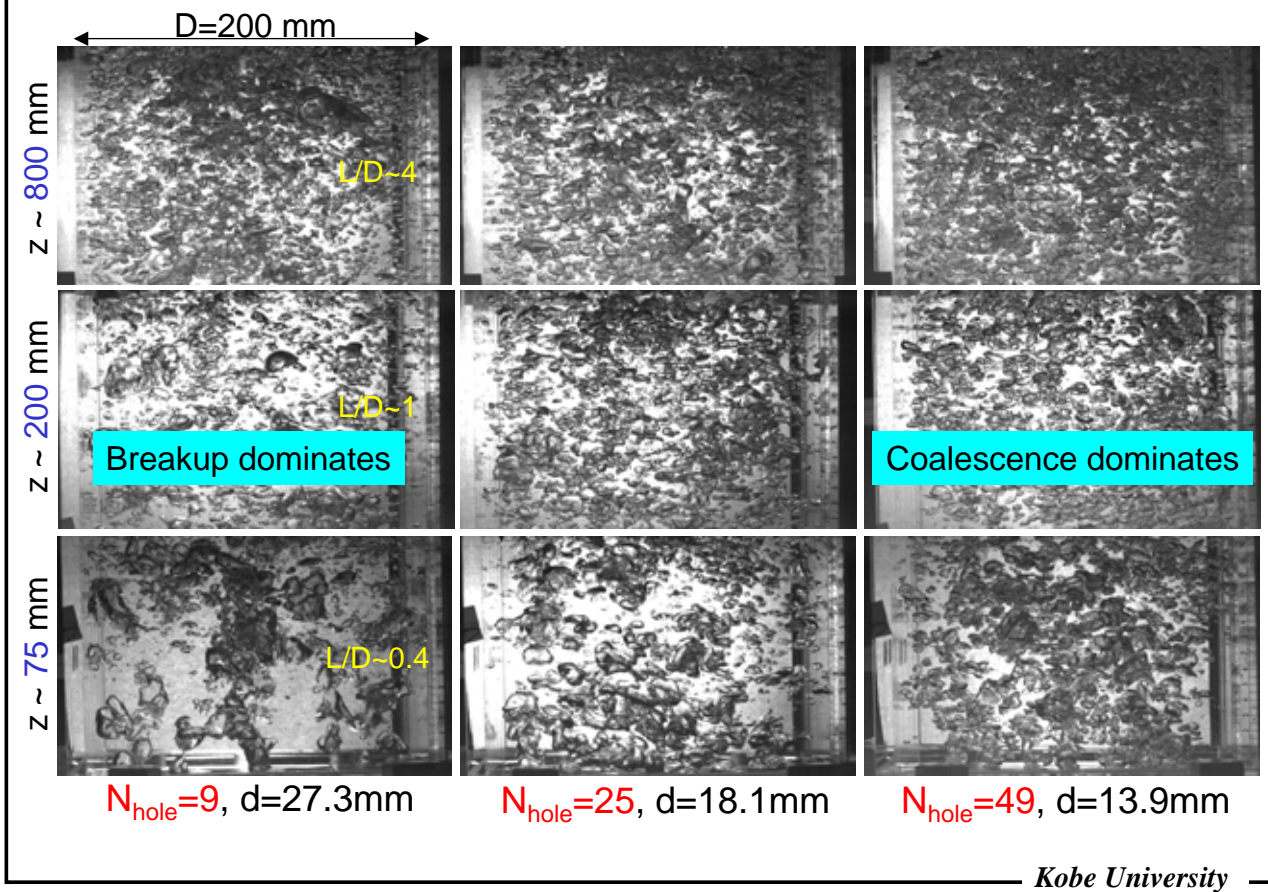
H x W x D : 1400 x 200 x 200 [mm]
 Fluids : Air-water at 25 °C and 1 atm
 Gas flow rate : $1.44 \times 10^{-3} \text{ m}^3/\text{s}$ ($J_G = 0.036 \text{ m/s}$)
 Initial liquid level : 1 m

Void distribution: Conductivity probe
 Bubble sizes: Images (2500 bubbles/elevation)

Gas diffusers



| | $N_{\text{hole}}=9$ | $N_{\text{hole}}=25$ | $N_{\text{hole}}=49$ |
|--------------------------|---------------------|----------------------|----------------------|
| Q_{hole} [cc/s] | 160 | 57.6 | 29.4 |
| d_{exp} [mm] | 27.3 | 18.1 | 13.9 |

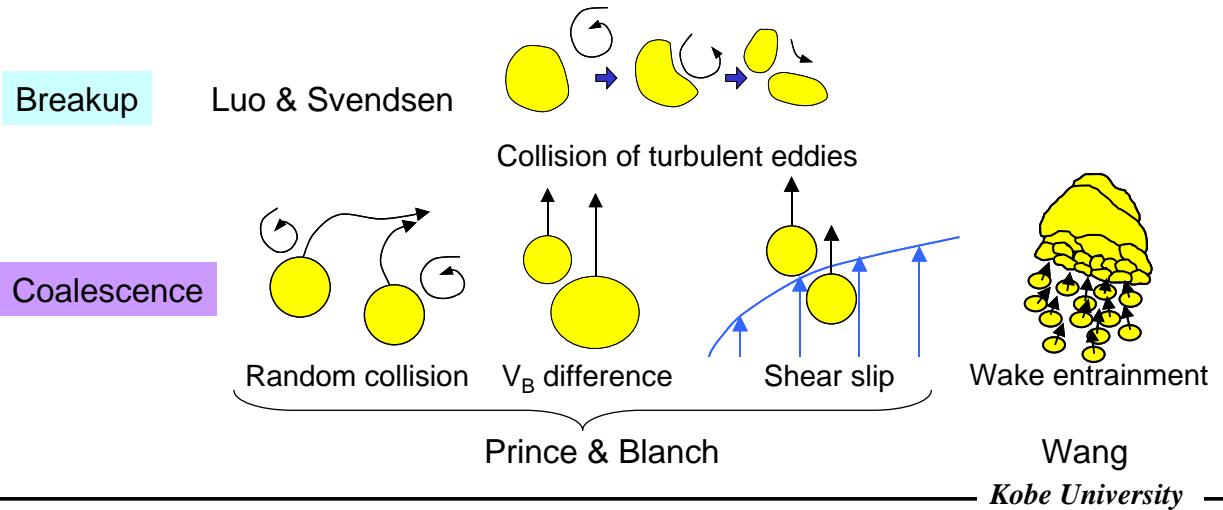


Bubble Coalescence & Breakup Models

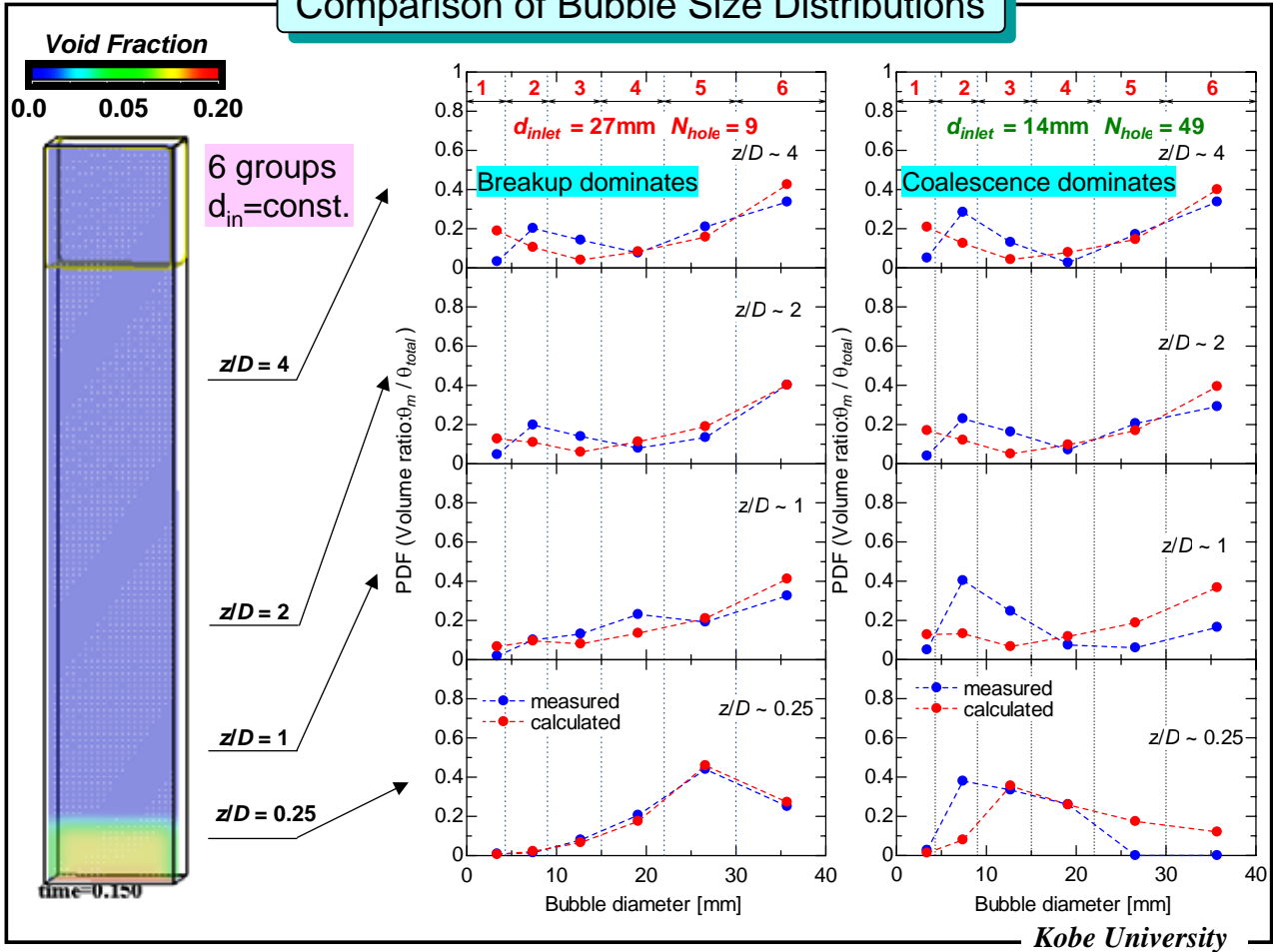
$$\frac{\partial n_m}{\partial t} + \nabla \cdot n_m \mathbf{V}_m = \sum_{m' \neq m} (\overset{\text{Source}}{\Gamma_{m' \rightarrow m}} - \overset{\text{Sink}}{\Gamma_{m \rightarrow m'}})$$

Source: $\sum_{m' \neq m} \Gamma_{m' \rightarrow m} = \int_{\theta_m}^{\infty} r_B(\theta_m, \theta_{m'}) f(\theta_{m'}) d\theta_{m'} + \frac{1}{2} \int_0^{\theta_m} r_C(\theta_{m'}, \theta_m) f(\theta_{m'}) f(\theta_m - \theta_{m'}) d\theta_{m'}$
 Breakup birth Coalescence birth

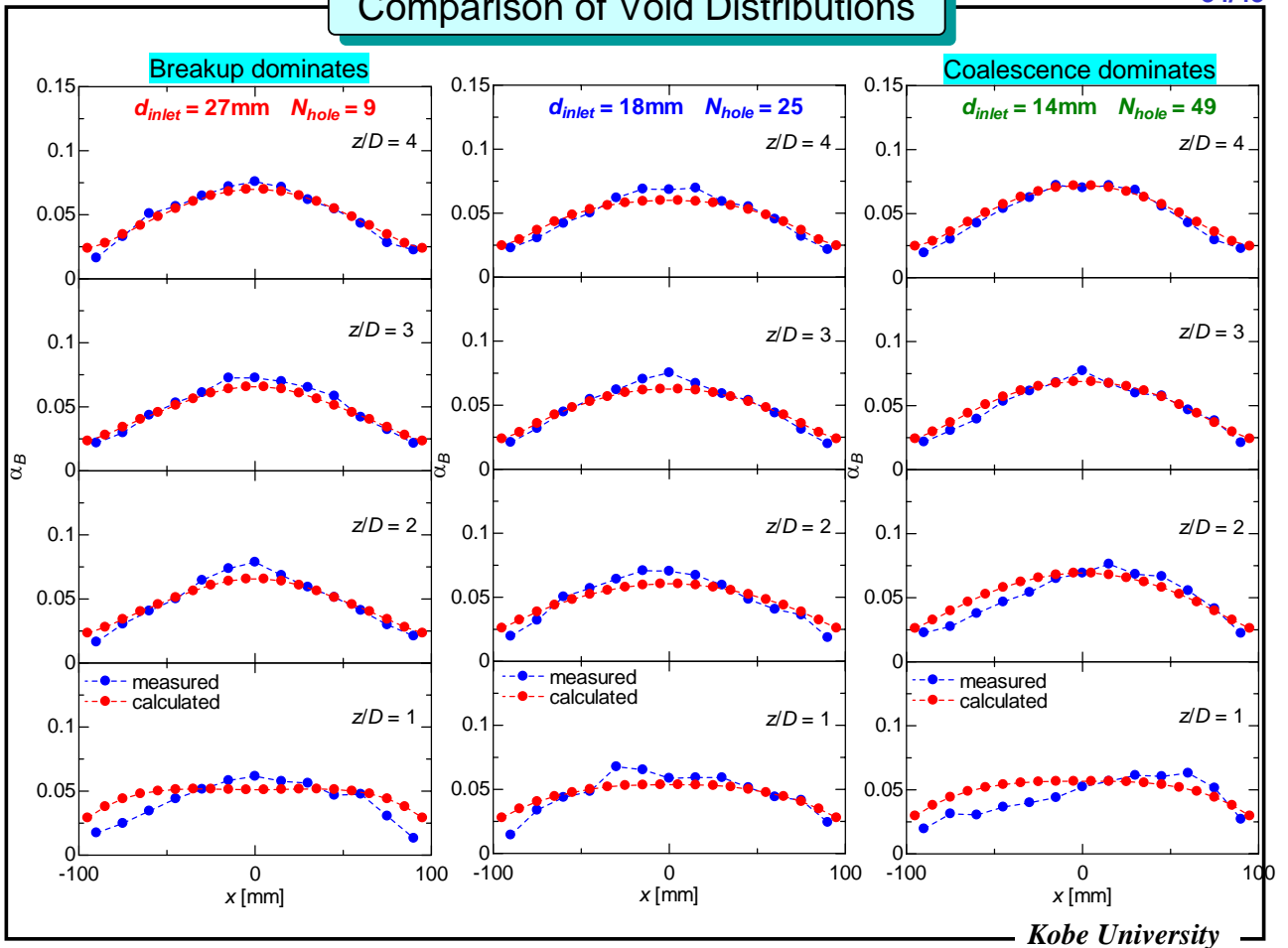
Sink: $\sum_{m' \neq m} \Gamma_{m \rightarrow m'} = \int_0^{\theta_m} \theta_{m'} r_B(\theta_{m'}, \theta_m) d\theta_{m'} \frac{f(\theta_m)}{\theta_m} + \int_0^{\infty} r_C(\theta_{m'}, \theta_m) f(\theta_{m'}) d\theta_{m'} f(\theta_m)$
 Breakup death Coalescence death



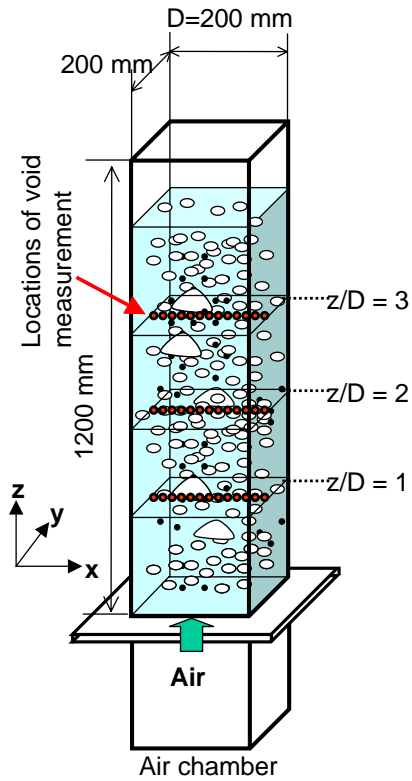
Comparison of Bubble Size Distributions



Comparison of Void Distributions



Experimental Setup and Condition



Silica particle: $d=60, 100, 150 \mu\text{m}$,
 $\rho=1320\text{kg/m}^3$
 hydrophilic, spherical

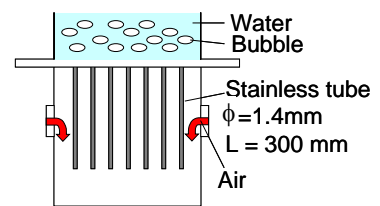
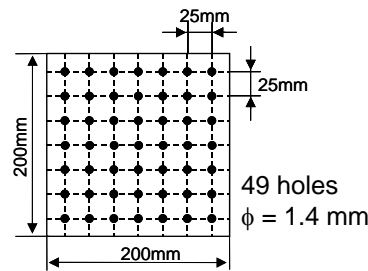
Particle concentration C_s : 0, 20, 40 vol. %

Gas volume flux J_G : 0.02, 0.034 m/s

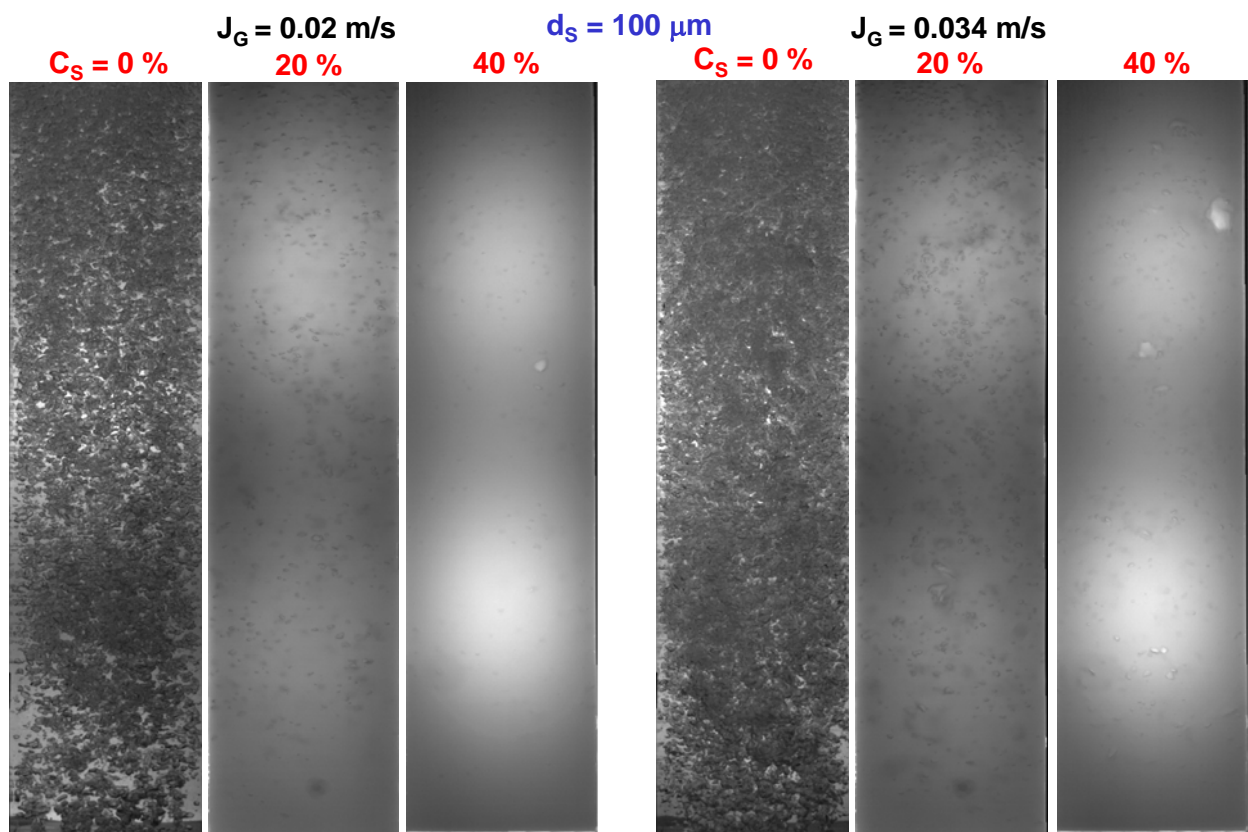
Inlet bubble diameter: 10.9, 12.6 mm

Initial water level: 1 m

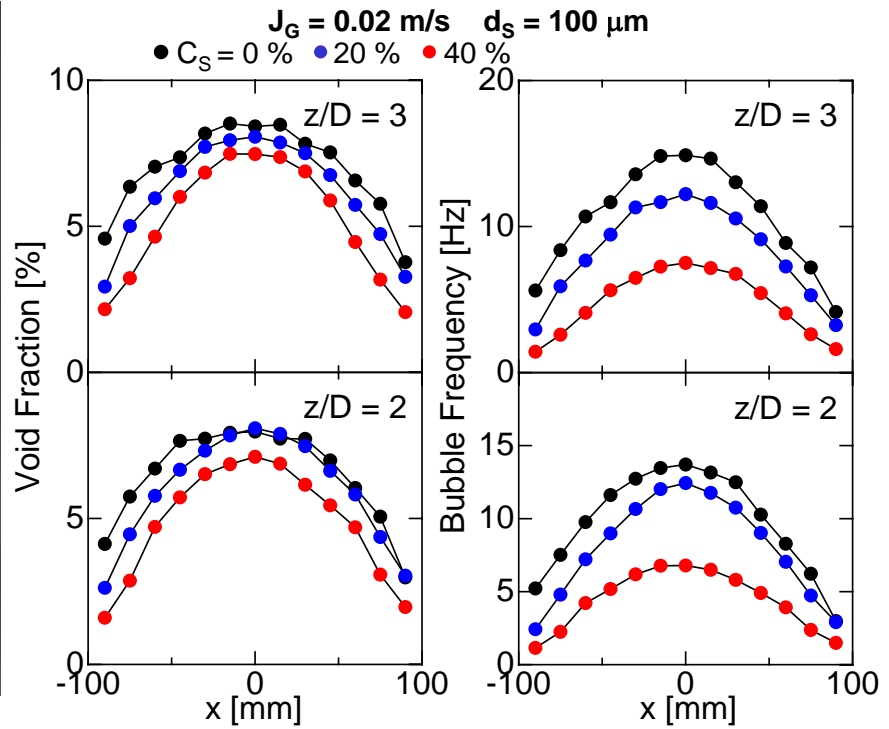
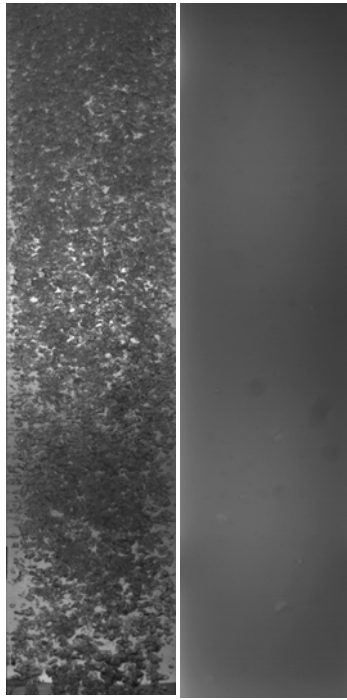
Void distribution: Conductivity probe



Effect of Particles on Flow Patterns

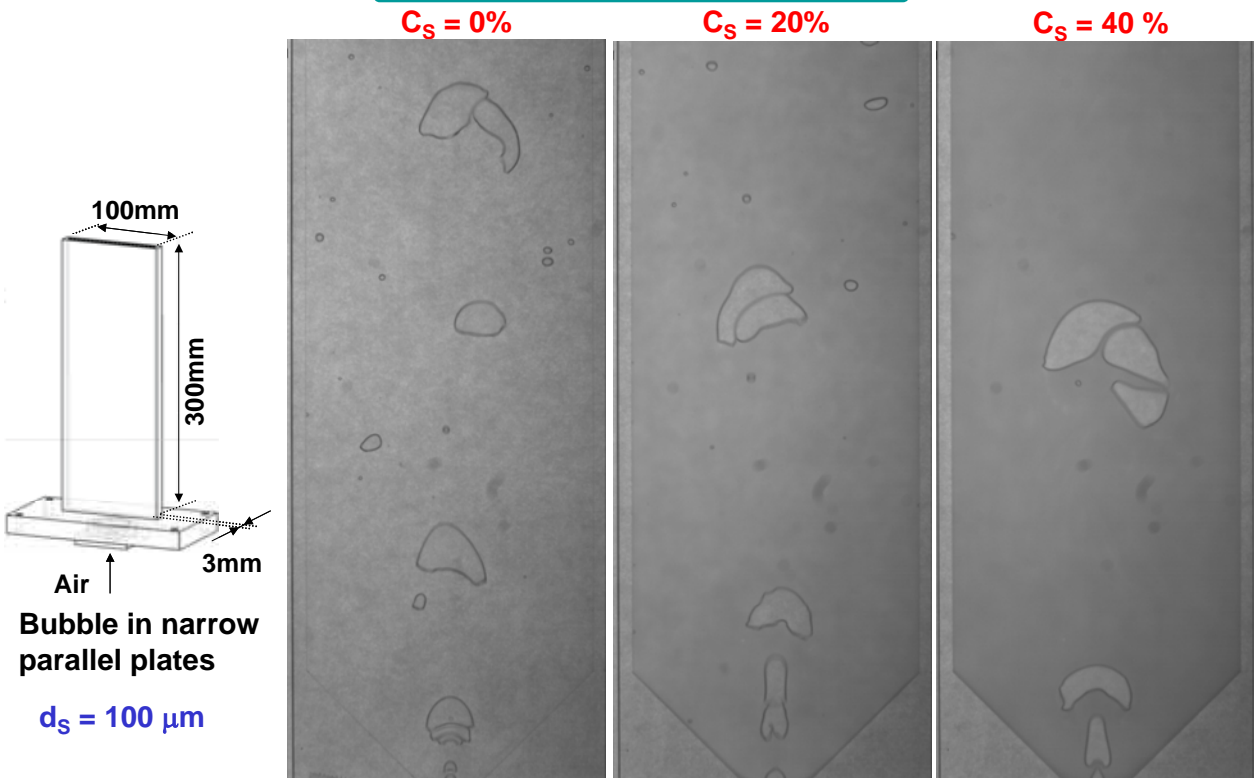


Particles seem to promote bubble coalescence.



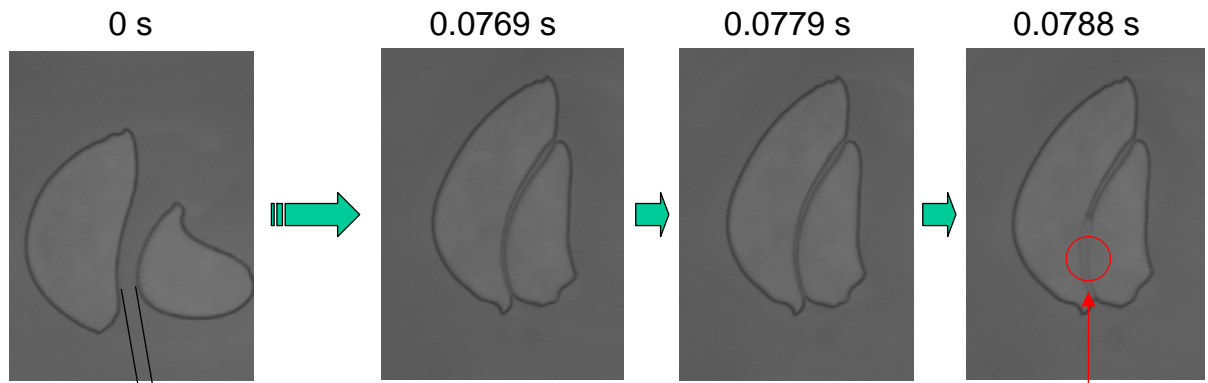
Void fraction decreases with increasing C_S .
 → Bubble velocity increases with C_S .

Number of bubbles decreases with increasing C_S .
 → Particles promote coalescence.



“Time required for coalescence” decreases with increasing C_S .
 → Increase in coalescence probability

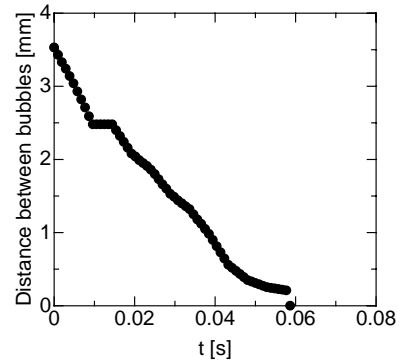
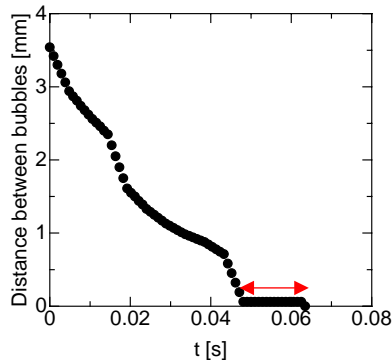
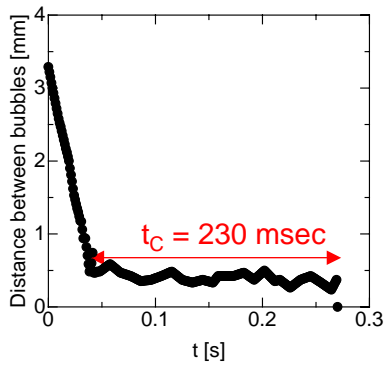
Measurement of Time required for Coalescence t_C



Distance between bubbles

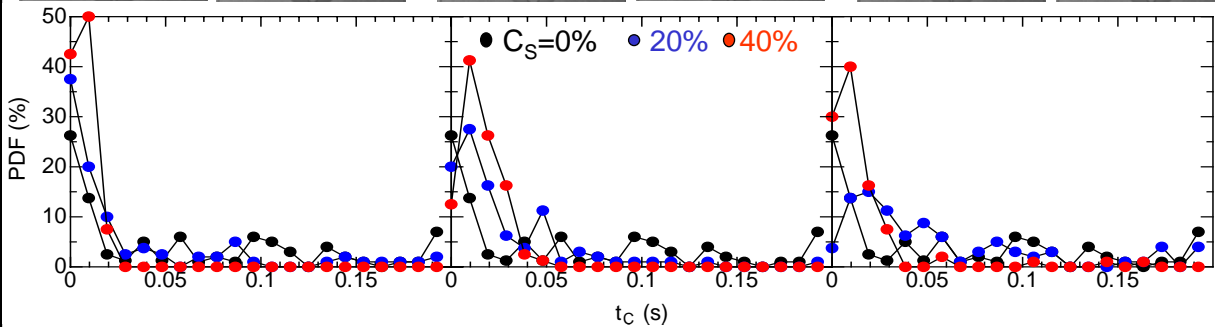
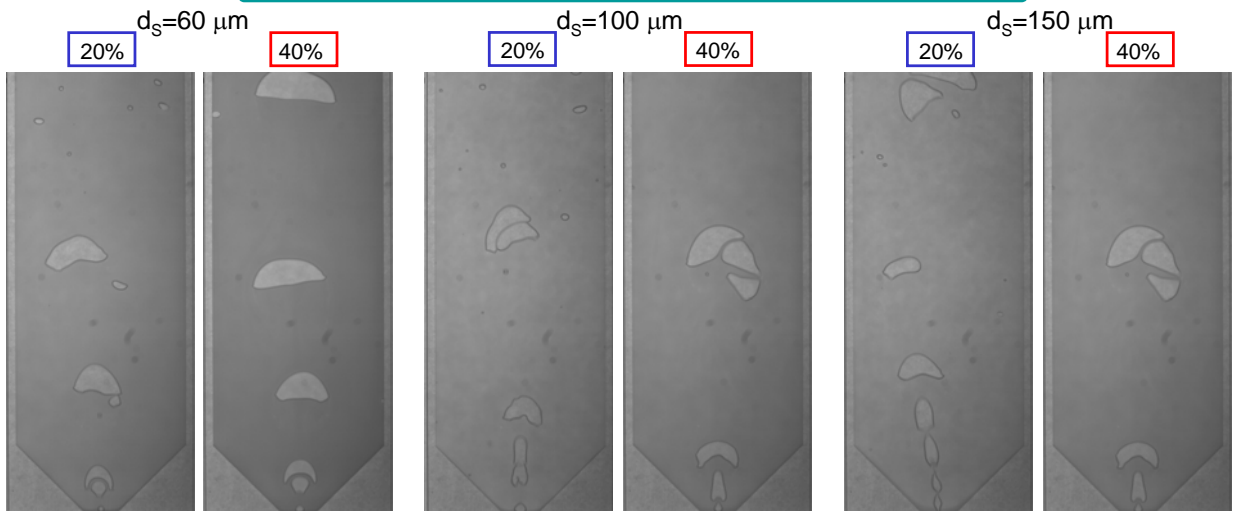
$t_C = 0.0788 - 0.0769 = 0.0019 \text{ s}$

Film breakup

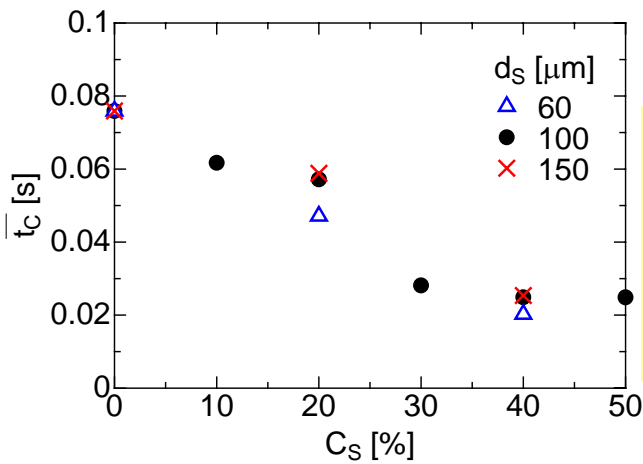


(a) Long contact: $t_C = 230 \text{ msec}$ (b) Short contact: $t_C = 15 \text{ msec}$ (c) Instantaneous coales. $t_C = 0$

Effect of Particle Size & Concentration on t_C



Coalescence probability increases with decreasing d_S and increasing C_S .

Introduction of a Particle Effect Multiplier to t_C 

t_C decreases with increasing C_S .

t_C decreases with d_S .

t_C reaches a minimum value at $C_S=40\%$.

Instantaneous coalescence

Coalescence efficiency: $P_C = \exp(-\beta t_C / \tau)$

β : particle effect multiplier

τ : contact time

$\beta=1$ at $C_S = 0\%$ (Gas-liquid)

$\beta=0$ for $C_S > 40\%$ (Instantaneous coalescence)

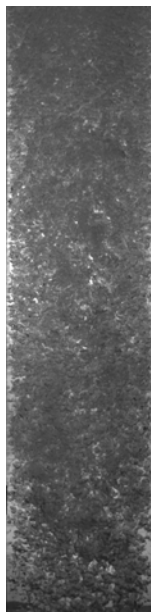
$\beta \sim 0.5$ at $C_S = 20\%$ (Experimental data)

Kobe University

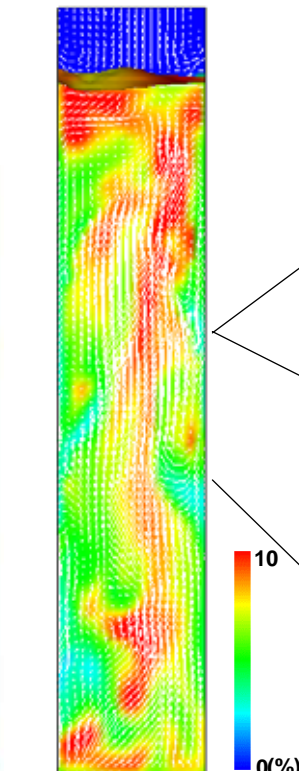
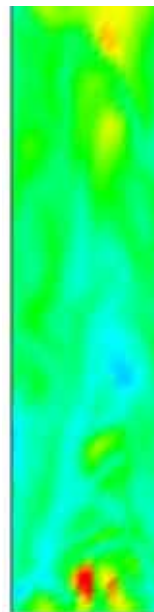
Comparison for $C_S = 0\%$

$J_G = 0.02$ m/s

$\beta = 1$

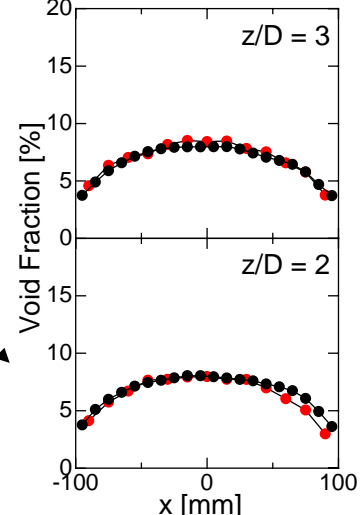
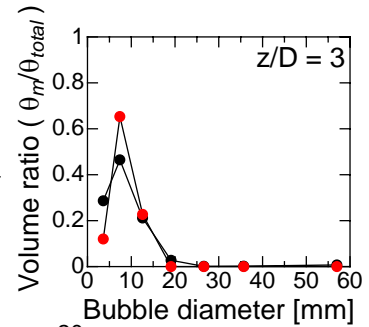


Void Fraction near the wall



Void Fraction on the center plane

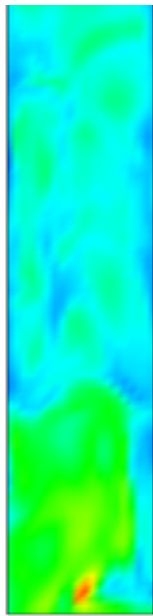
● Calculated ● Measured



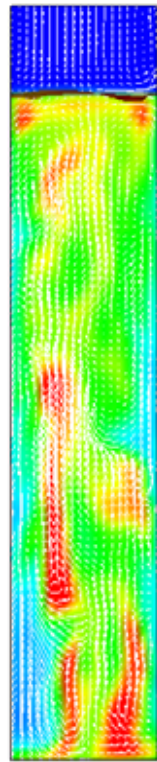
Kobe University

Comparison for $C_S = 20\%$, $d_S = 100\ \mu\text{m}$

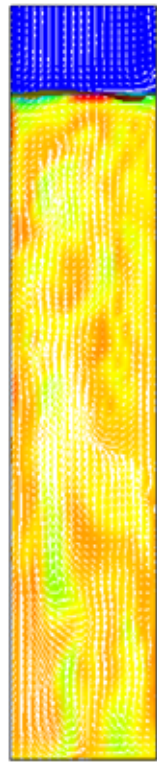
$J_G = 0.02\ \text{m/s}$
 $\beta = 0.5$



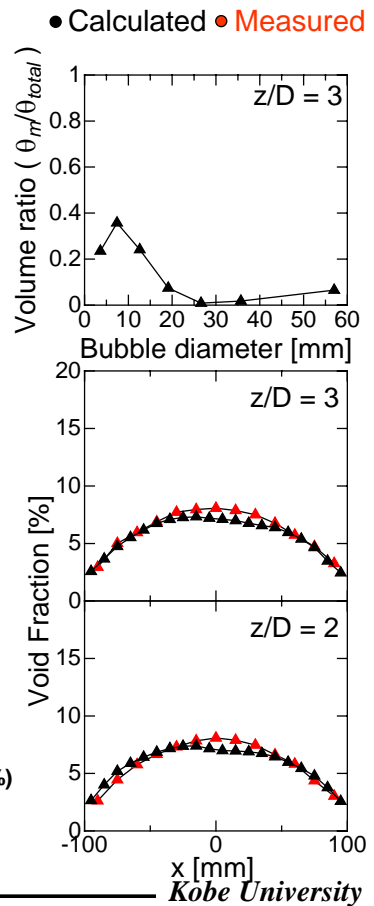
Void Fraction near the wall



Void Fraction on the center plane

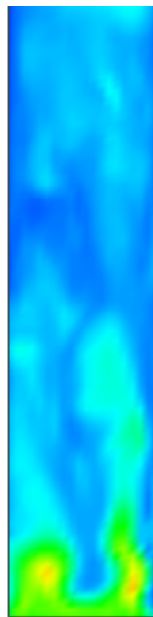


Particle Fraction on the center plane

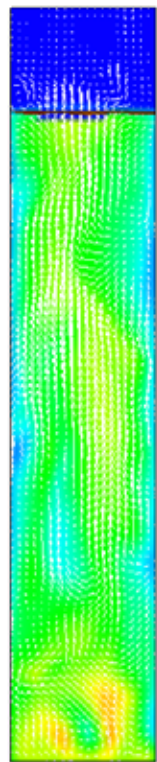


Comparison for $C_S = 40\%$, $d_S = 100\ \mu\text{m}$

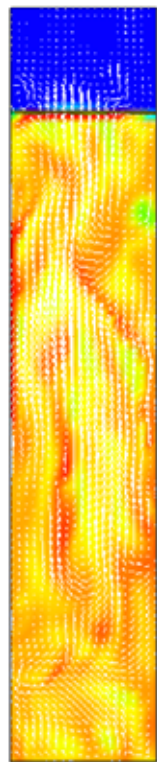
$J_G = 0.02\ \text{m/s}$
 $\beta = 0$



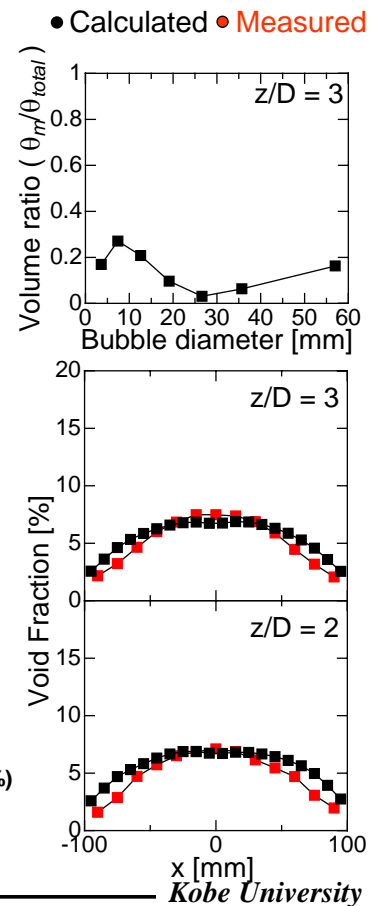
Void Fraction near the wall



Void Fraction on the center plane



Particle Fraction on the center plane



Summary

Accurate interface tracking simulation of contaminated bubbles and drops for a wide range of fluid properties and Reynolds number is possible, provided that physical properties for the adsorption and desorption kinetics are available.

Interface tracking simulation of mass transfer from a bubble for a wide range of Sc and Re numbers is also feasible, provided that HPCs are available.

Fine particles promote bubble coalescence and the effects of particles can be reasonably predicted by introducing a multiplier to the time required for coalescence.