

Parameter Identification in Partial Differential Equations

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Differentiation of data

Not strictly a parameter identification problem, but good motivation.
Appears often as a subproblem.

Given noisy observation of a function f ,

$$f^\delta(x) = f(x) + n^\delta(x), \quad x \in [0, 1]$$

Find its derivative




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Differentiation of data

Pointwise noise at measurement points identically normally distributed with mean 0 and variance δ .

Law of large numbers yields

$$\int_0^1 |n^\delta(x)|^2 dx \approx \delta^2$$



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Differentiation of data

Example: $n^\delta(x) = \sqrt{2\delta} \sin(2\pi kx)$

$$\frac{df^\delta}{dx}(x) = \frac{df}{dx}(x) + \sqrt{2}\pi\delta k \sin(2k\pi x)$$

Noise satisfies $\int_0^1 |n^\delta(x)|^2 dx = \delta^2$

But error in derivative is large




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Differentiation of data

Error:

$$\left(\int_0^1 \left(\frac{df^\delta}{dx}(x) - \frac{df}{dx}(x) \right)^2 dx \right)^{1/2} = 2\pi\delta k$$

$$\sup_{x \in [0,1]} \left| \frac{df^\delta}{dx}(x) - \frac{df}{dx}(x) \right| = \sqrt{22}\pi\delta k$$

Arbitrarily large since k can be arbitrarily large



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Conclusion 1

Without regularization and without further information, the error between the exact and noisy solution can be arbitrarily large, even if the noise is arbitrarily small.



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Differentiation of data

Additional information:

Solution is smooth (e.g. twice differentiable)

Regularization: e.g. smoothing by elliptic PDE

$$-\alpha \frac{d^2 f_\alpha}{dx^2}(x) + f_\alpha(x) = f^\delta(x), \quad f_\alpha(0) = f_\alpha(1) = 0$$

Variational principle for elliptic PDEs: minimize

$$H_\alpha(f_\alpha) = \int_0^1 (f_\alpha(x) - f^\delta(x))^2 dx + \alpha \int_0^1 \left(\frac{df_\alpha}{dx}(x) \right)^2 dx$$



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Differentiation of data

Detailed estimate:

Lecture notes, p. 7



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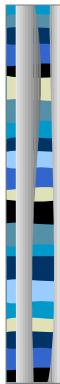
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Computerized Tomography

(A nice link between Austria & Norway)

Problem: reconstruct a spatial density f in a domain D from measurements of X-rays traveling through the domain.

X-rays travel along rays, parametrized by distance s from origin and its unit normal vector ω




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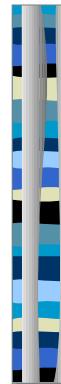
Computerized Tomography

X-ray beam has intensity I

Model: decay of intensity in a small distance Δt proportional to I , f , and distance

$$\frac{\Delta I(sw + tw^\perp)}{\Delta t} = -I(sw + tw^\perp)f(sw + tw^\perp)$$

Limit Δt to zero

$$\frac{dI(sw + tw^\perp)}{dt} = -I(sw + tw^\perp)f(sw + tw^\perp)$$



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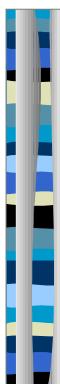
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Measurement: intensity at emitter and detector

$$I_L(s, w) := I(sw + Lw^\perp) \text{ and } I_0(s, w) := I(sw)$$

Parameter identification problem for system of ordinary differential equations (first-order)

Overdetermined for given f since initial and final value are known




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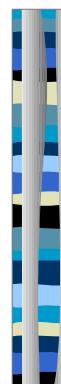


Computerized Tomography

Integrate ODE

$$\ln I(sw + Lw^\perp) - \ln I(sw) = - \int_0^L f(sw + tw^\perp) dt$$

Leads to inversion of the Radon transform

$$\ln I_0(s, w) - \ln I_L(s, w) = \mathcal{R}f(s, w)$$



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Computerized Tomography

Radon Transform

$$\mathcal{R}f(s, w) = \int_{\mathbb{R}} f(sw + tw^\perp) dt$$

Exact inversion formula by Johann Radon 1917.
SVD computed by McCormick et. al. in 1960s,
Nobel Prize for Medicine in the 1970

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Computerized Tomography

Radial symmetry, $r = \sqrt{s^2 + t^2}$

$$\mathcal{R}f(s, w_0) = 2 \int_s^\rho \frac{rF(r)}{\sqrt{r^2 - s^2}} dr$$

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Computerized Tomography

Radial symmetry, $r = \sqrt{s^2 + t^2}$

$$\mathcal{R}f(s, w_0) = 2 \int_s^\rho \frac{rF(r)}{\sqrt{r^2 - s^2}} dr$$

Abel Integral equation !

SVs decay with half speed as for differentiation

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Groundwater Filtration

Identification of diffusivity of sediments from an observation of the piezometric head (describing the flow)

Knowledge of diffusivity allows to draw conclusions about the structure of the groundwater

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Groundwater Filtration

Mathematical model

$$-\operatorname{div} (a \nabla u) = f$$

Diffusivity a , piezometric head u (measured),
density of water sources and sinks f

+ Appropriate boundary conditions (e.g. $u=0$)



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Groundwater Filtration

Instead of exact data, we only know noisy
measurement

$$u^\delta(x) = u(x) + n^\delta(x), \quad x \in \Omega$$

Typically the noise $n(x)$ comes from identically
normally distributed random variables

From a law of large numbers this gives an
estimate

$$\int_0^1 |n^\delta(x)|^2 dx \approx \delta^2$$



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Groundwater Filtration

Consider 1D version of inverse problem, for
simplicity with $\frac{du}{dx}(0) = 0$

We can integrate the equation to obtain

$$a(x) \frac{du}{dx}(x) = \int_0^x f(y) dy$$

Hence, if the derivative of u does not vanish,
 a is determined uniquely (Identifiability)



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Groundwater Filtration

Detailed estimate:

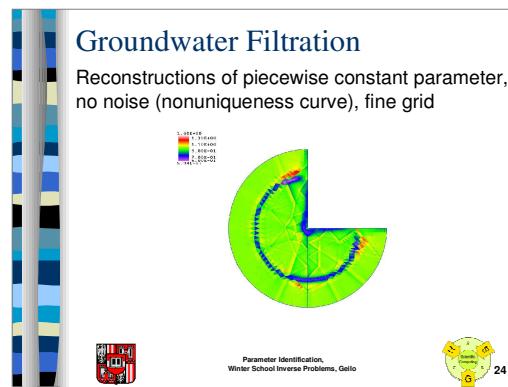
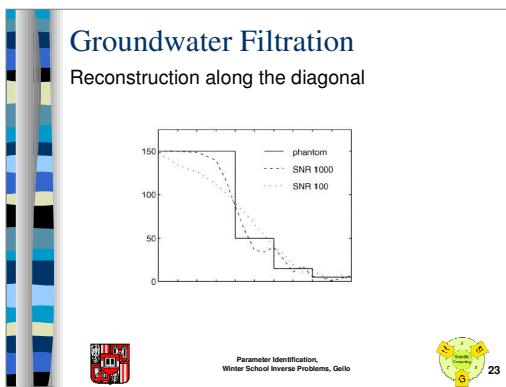
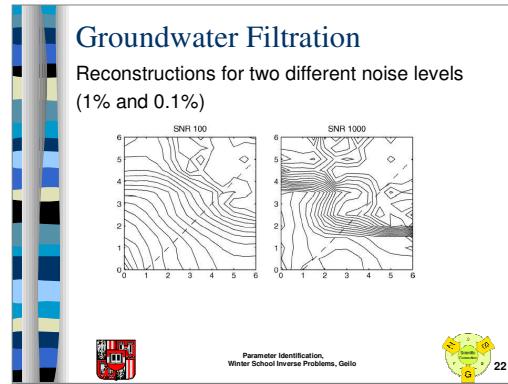
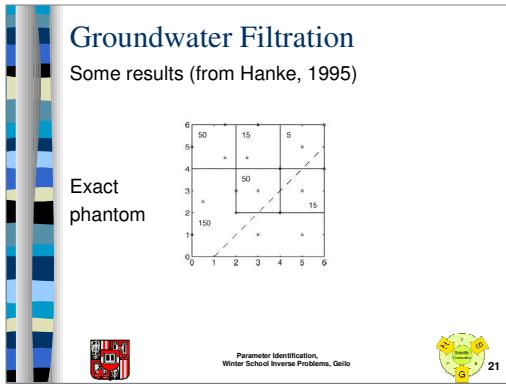
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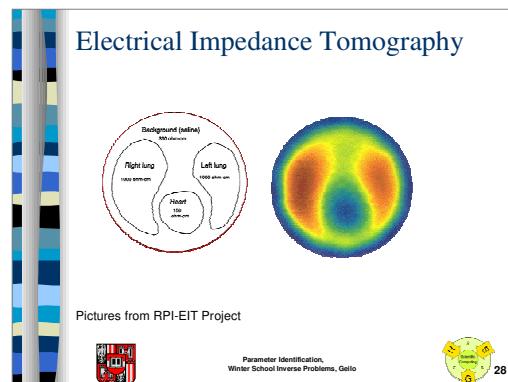
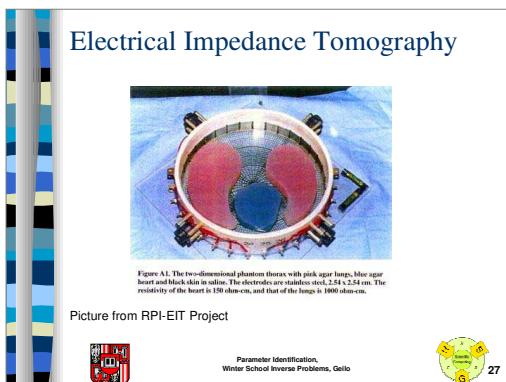
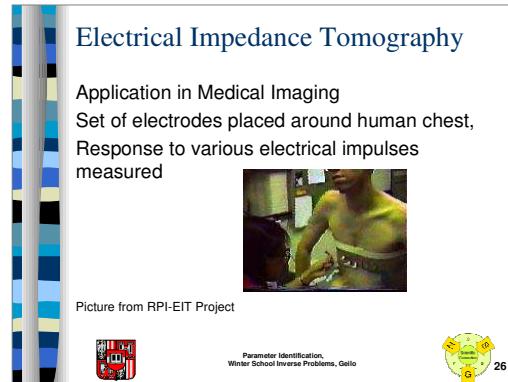
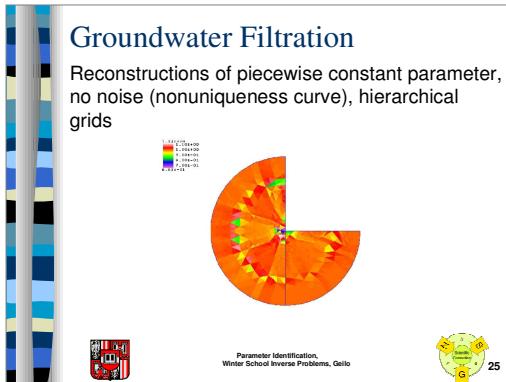


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Electrical Impedance Tomography

Mathematical model: Maxwell equations, reduced to potential equation

$$\operatorname{div} (a \nabla u) = 0 \quad \text{in } D$$

$$u = f \quad \text{on } \partial D$$

u is electrical potential, f applied voltage pattern
 a denotes the (unknown) conductivity



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Electrical Impedance Tomography

Measurement: current density on the boundary for different voltage patterns

$$g_f = a \frac{\partial u}{\partial n} \quad \text{on } \partial D$$

Idealized mathematical model: Dirichlet-Neumann map

$$\Lambda_a : f \mapsto g_f$$


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Electrical Impedance Tomography

INVERSE CONDUCTIVITY PROBLEM

In practice finite number of measurements

$$\Lambda_a(f_j) \text{ for } j = 1, \dots, N$$

N typically very large



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Electrical Impedance Tomography

To compute output, we have to solve solutions of N partial differential equations

$$-\operatorname{div} (a \nabla u^j) = 0$$

Problem of extremely large scale



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Electrical Impedance Tomography

Interesting special case: piecewise constant conductivities

$$a(x) = \begin{cases} a_1 & \text{if } x \in \Omega \subset D \\ a_2 & \text{if } x \in D \setminus \Omega. \end{cases}$$

Real interest is the subset where a takes a value different from a_1 (e.g. a tumour)

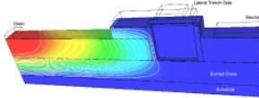
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Semiconductor Devices

Related problem to impedance tomography appears for semiconductor devices: **inverse dopant profiling**



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Inverse Dopant Profiling

Identify the **device doping profile** from measurements Current-Voltage map

Analogous definitions of voltage and current, but more complicated mathematical model

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Mathematical Model

Stationary Drift Diffusion Model:
PDE system for potential V , electron density n and hole density p

$$\begin{aligned} \operatorname{div}(\epsilon_s \nabla V) &= q(n - p - C) \\ \operatorname{div}(D_n \nabla n - \mu_n n \nabla V) &= 0 \\ \operatorname{div}(D_p \nabla p + \mu_p p \nabla V) &= 0 \end{aligned}$$

in Ω (subset of R^2)

Doping Profile $C(x)$ enters as source term

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Boundary Conditions

Boundary of Ω : homogeneous Neumann boundary conditions on Γ_+ and

$$V = V_D = U + V_{bi} = \textcolor{red}{U} + U_T \ln \left(\frac{n_D}{n_i} \right)$$

$$\textcolor{blue}{n} = n_D = \frac{1}{2} \left(\textcolor{red}{C} + \sqrt{\textcolor{red}{C}^2 + 4n_i^2} \right)$$

$$\textcolor{blue}{p} = p_D = \frac{1}{2} \left(-\textcolor{red}{C} + \sqrt{\textcolor{red}{C}^2 + 4n_i^2} \right)$$

on Dirichlet boundary Γ_D (Ohmic Contacts)

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Device Characteristics

Measured on a contact Γ_0 on Γ_D :
Outflow current density

$$\textcolor{blue}{I} = (D_n \nabla n - \mu_n n \nabla V - D_p \nabla p - \mu_p p \nabla V) \cdot \nu$$

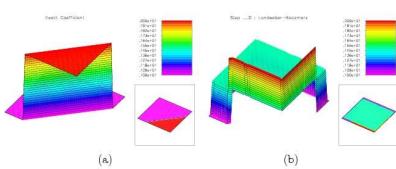
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Numerical Tests

Test for a P-N Diode



Real Doping Profile



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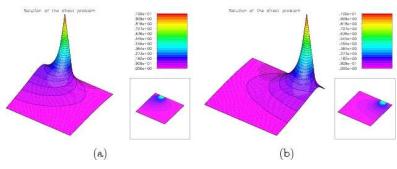
Initial Guess



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Numerical Tests

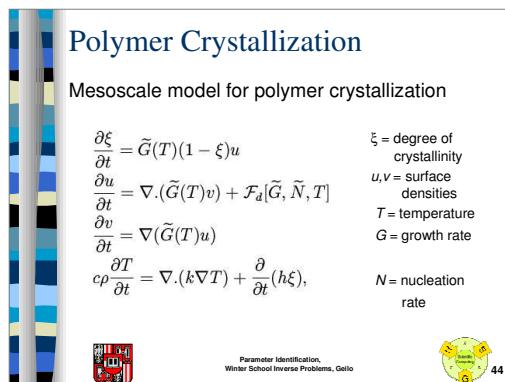
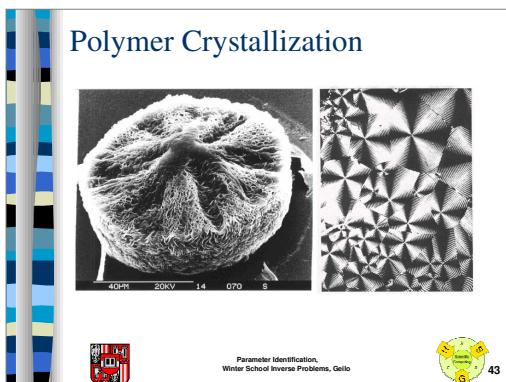
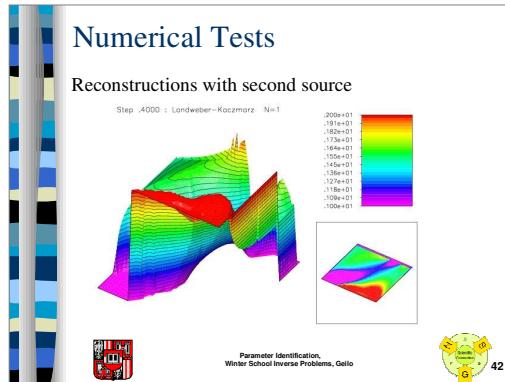
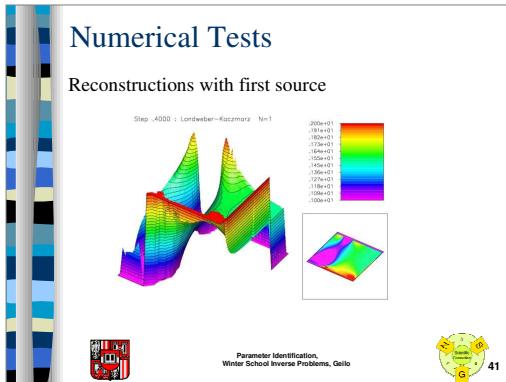
Different Voltage Sources



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Polymer Crystallization

Source term

$$\mathcal{F}_1[\tilde{G}, \tilde{N}, T](x, t) := 2\tilde{N}(T(x, t))_t$$

$$\mathcal{F}_2[\tilde{G}, \tilde{N}, T](x, t) := 2\pi\tilde{G}(T(x, t))(\tilde{N}(T(x, t)) - \tilde{N}(T(x, 0)))$$

$$\mathcal{F}_3[\tilde{G}, \tilde{N}, T](x, t) := 4\pi\tilde{G}(T(x, t)) \int_0^t \tilde{G}(T(x, s))(\tilde{N}(T(x, s)) - \tilde{N}(T(x, 0))) ds$$


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Polymer Crystallization

Traditional way of determining nucleation rate as function of temperature:

- Make separate experiment for each value of T
- Count (by eyes !!!) the final number of crystals (typically $\gg 10^6$).
- Divide by volume

Extremely expensive, extremely time-consuming !



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Polymer Crystallization

Idea: determine nucleation rate as function of temperature by single nonisothermal experiment

Measured data: temperature T at the boundary of the sample

Degree of crystallinity at final time



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Polymer Crystallization

Results

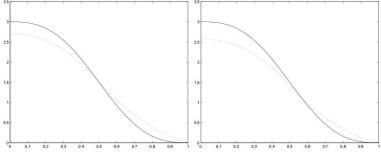


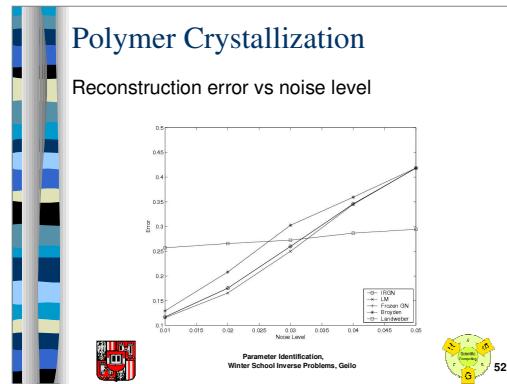
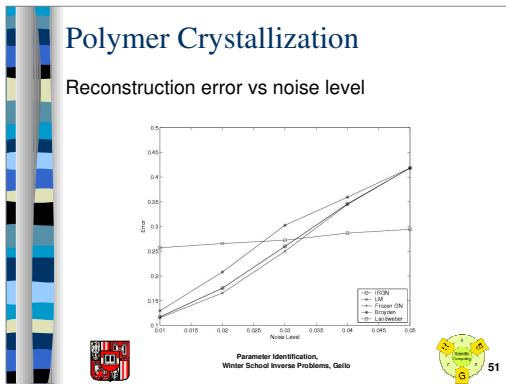
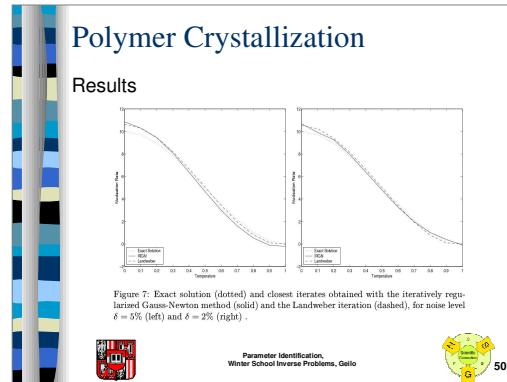
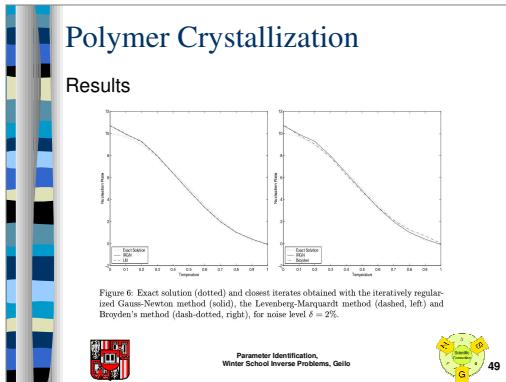
Figure 2: Exact solution (solid) and iterate at $k = k_s$ (dotted) vs. u , $\delta = 0.025$, $f = 50$. The left figure corresponds to an experiment with temperature values in the whole range

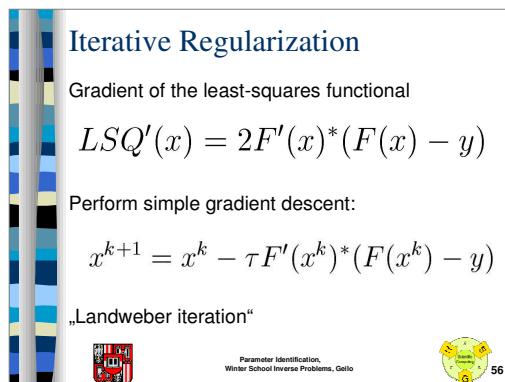
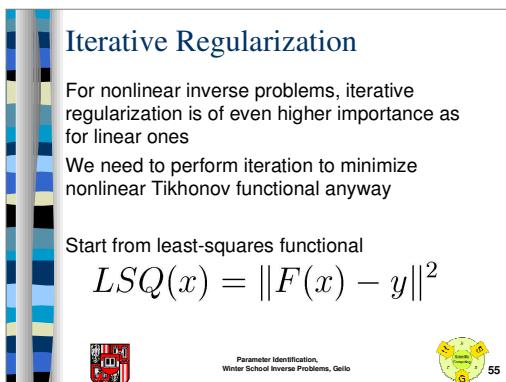
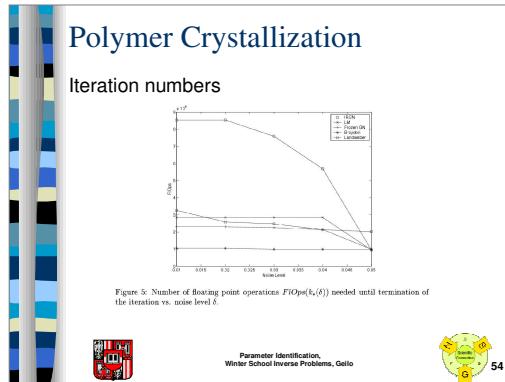
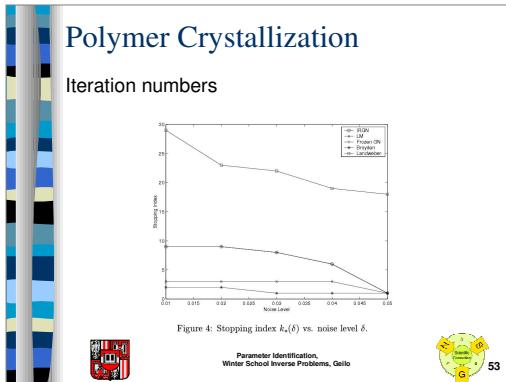


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Iterative Regularization

Where is the regularization parameter ??!!

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Iterative Regularization

Regularization by appropriate early stopping,
e.g. discrepancy principle

$$k_* = \inf\{k \in \mathbb{N} \mid \|F(x^k) - y^\delta\| \leq \delta\}$$

Each step of Landweber iteration is well-posed,
so we obtain iterative regularization method

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Iterative Regularization

Analysis of Landweber iteration: *Hanke-Neubauer-Scherzer 1994, Scherzer 1995*

Usual semi-convergence properties as for linear problems, but restriction on the nonlinearity of the problem (replacing regularity of $F'(x^*)$)

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Iterative Regularization

Landweber iteration is forward Euler time
discretization (with time step τ of the flow)

$$x'(t) = -F'(x(t))^*(F(x(t)) - y)$$

„Asymptotical regularization“ (*Tautenhahn 1994*)
Iteration

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Iterative Regularization

We can consider other time discretizations:

Backward Euler

$$x^{k+1} = x^k - \tau F'(x^{k+1})^*(F(x^{k+1}) - y)$$

„Iterated Tikhonov Regularization“ (Hanke-Groetsch 1999, Groetsch-Scherzer 1999)



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Iterative Regularization

We can consider other time discretizations:

Semi-Implicit Euler

$$x^{k+1} = x^k - \tau F'(x^k)^*(F(x^k) - y + F'(x^k)(x^{k+1} - x^k))$$

„Levenberg-Marquardt“ (Hanke 1995)

Arbitrary Runge-Kutta methods possibly, even inconsistent ones (Rieder, 2005)



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Iterative Regularization

Levenberg-Marquardt can be rewritten to

$$(F'(x^k)^* F'(x^k) + \alpha^k)(x^{k+1} - x^k) = -F'(x^k)^*(F(x^k) - y)$$

Tikhonov regularization of the linear Newton-equation

$$F'(x^k)(x^{k+1} - x^k) = -(F(x^k) - y)$$



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Iterative Regularization

General approach for Newton-type methods:
Apply linear regularization to

$$F'(x^k)(x^{k+1} - x^k) = -(F(x^k) - y)$$

„Iteratively regularized Gauss-Newton“
(Kaltenbacher-Neubauer-Scherzer 1996, Kaltenbacher 1997)
„Newton-CG“ (Hanke, 1998)



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Iterative Regularization

„Newton-Landweber“ (Kaltenbacher 1999)
 „Broyden“ (Kaltenbacher, 1997)
 „Multigrid / Discretization“ (Kaltenbacher, 2000-2003)

In particular for parameter identification:
 „Levenberg-Marquardt SQP“ (B-Mühlhuber 2002)
 „Newton-Kaczmarz“ (B-Kaltenbacher 2004)

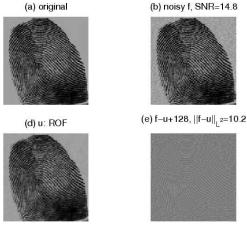


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Total Variation Denoising

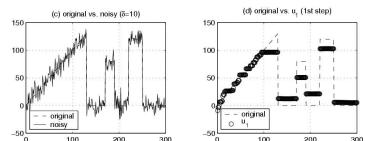


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Total Variation Denoising



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TV-Images

Basic problem in denoising:

given noisy version g of image f , find
 „optimal“ approximation u of f

Basic problem in deblurring:

given noisy version g of Kf , find
 „optimal“ approximation u of f



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TV-Images

What is optimal ?

Satisfy other criterion, e.g., minimization

$$J(u) \rightarrow \min_u$$

over a class of approximations



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ROF-Model

What is a suitable class of approximations ?

Discrepancy principle: allow images with

$$\int_{\Omega} (u - g)^2 dx \leq \sigma^2$$

with σ being the noise level:

$$\int_{\Omega} (f - g)^2 dx \leq \sigma^2$$



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ROF-Model

Rudin-Osher-Fatemi 1989 (ROF):

Minimize total variation

$$J(u) = \int_{\Omega} |\nabla u| dx$$

subject to

$$\int_{\Omega} (u - g)^2 dx \leq \sigma^2$$



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ROF-Model

Deblurring:

Minimize total variation

$$J(u) = \int_{\Omega} |\nabla u| dx$$

subject to

$$\int_{\Omega} (Ku - g)^2 dx \leq \sigma^2$$



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ROF-Model

Acar-Vogel 1994:

ROF is regularization method (well-posedness + convergence as σ to zero)

Chambolle-Lions 1997:

Solution unique, there exists Lagrange multiplier λ , problem equivalent to

$$\frac{\lambda}{2} \int_{\Omega} (u - g)^2 dx + J(u) \rightarrow \min_u$$

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ROF-Model

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Total Variation

Rigorous definition of total variation

$$|u|_{TV} := \sup_{\mathbf{g} \in C_0^\infty(\Omega)^d} \int_{\Omega} u \operatorname{div} \mathbf{g} dt$$

$$BV(\Omega) := \{ u \in L^1(\Omega) \mid |u|_{TV} < \infty \}.$$

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Total Variation

BV includes discontinuous functions:

1D example

$$u^R(x) = \begin{cases} 1 & \text{if } |x| \leq R \\ 0 & \text{else.} \end{cases}$$

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Total Variation

1D example

$$\int_{\Omega} u \operatorname{div} g \, dt = \int_{-R}^R \frac{dg}{dt} \, dt = g(R) - g(-R)$$

$$|u|_{TV} := \sup_{g \in C_0^\infty([-1,1])} [g(R) - g(-R)] = 2$$



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Total Variation

Formal optimality for TV-Denoising

$$u - y^\delta = \alpha \operatorname{div} \left(\frac{\nabla u}{|\nabla u|} \right)$$

Term on the right-hand side corresponds to
mean curvature of level sets of u



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Duality

Consider TV-regularization

$$\inf_u \left[\int_{\Omega} (u - y^\delta)^2 \, dt + \alpha |u|_{TV} \right] =$$

$$\inf_u \sup_{\mathbf{g}} \left[\int_{\Omega} (u - y^\delta)^2 \, dt + 2\alpha \int_{\Omega} u \operatorname{div} \mathbf{g} \, dt \right]$$



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Duality

Exchange inf and sup

$$\sup_{\mathbf{g}} \inf_u \left[\int_{\Omega} (u - y^\delta)^2 \, dt + 2\alpha \int_{\Omega} u \operatorname{div} \mathbf{g} \, dt \right]$$

Solve inner inf problem for u

$$u = y^\delta - \alpha \operatorname{div} \mathbf{g}$$



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Duality

Remains sup problem for \mathbf{g}

$$\sup_{\mathbf{g}} \left[\alpha^2 \int_{\Omega} (\operatorname{div} \mathbf{g})^2 dt + 2\alpha \int_{\Omega} (y^\delta - \alpha \operatorname{div} \mathbf{g}) \operatorname{div} \mathbf{g} dt \right]$$

Rewrite equivalently

$$-\int_{\Omega} (\alpha \operatorname{div} \mathbf{g} - y^\delta)^2 dx \rightarrow \max_{|\mathbf{g}|_\infty \leq 1}$$



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Duality

Introduce $p := \alpha g$

$$\int_{\Omega} (\operatorname{div} p - y^\delta)^2 dt \rightarrow \min_{|p|_\infty \leq \alpha}$$

Dual TV problem (Chambolle 2003)

Dual space of BV

$$BV^* := \{ q = \operatorname{div} p \mid p \in L^\infty(\Omega) \}$$

$$\|q\|_{BV^*} := \inf \{ \|p\|_\infty \mid q = \operatorname{div} p \}$$



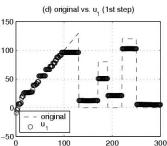
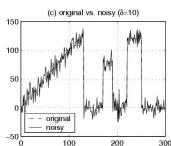
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Stair-Casing

Lecture notes p.20



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ROF-Model

Meyer 2001 (and others before):

ROF has a systematic error, since

$$J(\hat{u}) < J(f)$$

This means that jumps in the reconstructed image are smaller than jumps in the true image



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Noise Decomposition

How to resolve the systematic error ?

Take some λ , and minimize TV-functional

$$\frac{\lambda}{2} \int_{\Omega} (u - g)^2 dx + J(u) \rightarrow \min_u$$

to obtain u_1 . Then take residual („noise“)

$$v_1 = g - u_1$$



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Noise Decomposition

Minimize

$$\frac{\lambda}{2} \int_{\Omega} (u - g - v_1)^2 dx + J(u) \rightarrow \min_u$$

to obtain next iterate u_2

The second step may **increase total variation !**



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Iterated Total Variation

Obtain new iterate u_k by minimizing

$$\frac{\lambda}{2} \int_{\Omega} (u - g - v_{k-1})^2 dx + J(u) \rightarrow \min_u$$

then decompose noise

$$v_k = v_{k-1} + f - u_k$$

Need not change the code, just the data !

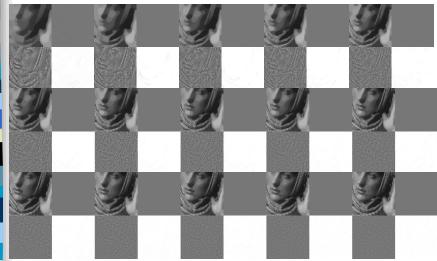


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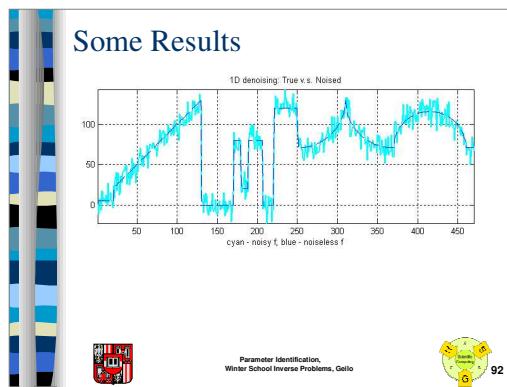
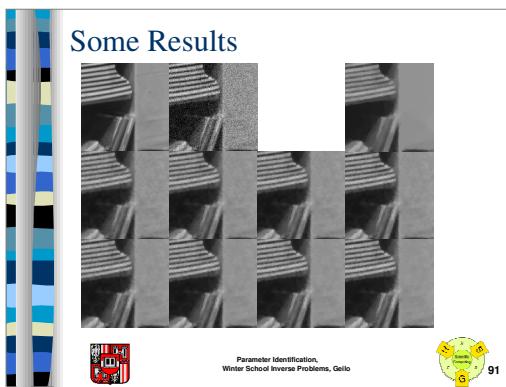
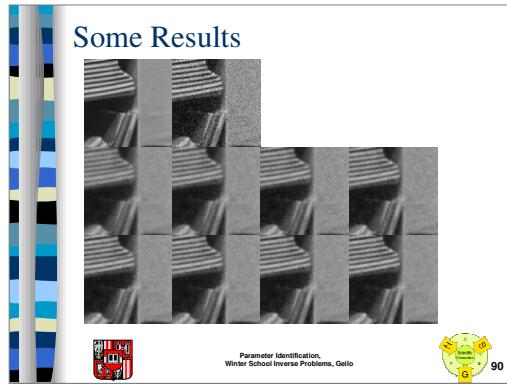
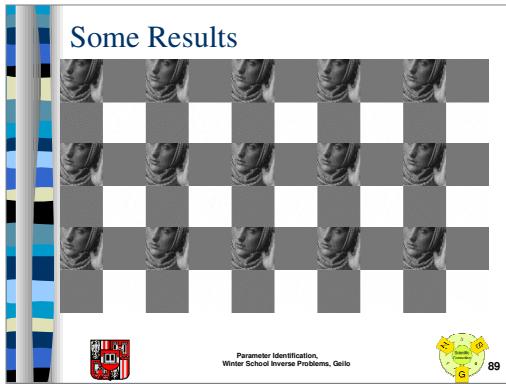
Some Results



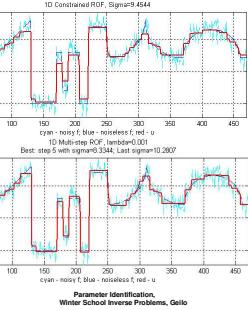
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Some Results



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Convergence Analysis

So why does that work ???

Expand fitting term, throw away constant terms:

$$\frac{\lambda}{2} \int_{\Omega} (u - g)^2 dx + J(u) - \int_{\Omega} v_{k-1}(u - g) dx \rightarrow \min_u$$

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Convergence Analysis

Optimality condition implies

$$v_k = v_{k-1} + f - u_k \in \partial J(u_k)$$

Write equivalent optimization problem

$$\begin{aligned} & \frac{\lambda}{2} \int_{\Omega} (u - g)^2 dx + J(u) - J(u_{k-1}) \\ & - \int_{\Omega} v_{k-1}(u - u_{k-1}) dx \rightarrow \min_u \end{aligned}$$

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Convergence Analysis

Hence, iterated TV is „proximal point algorithm“

$$\frac{\lambda}{2} \int_{\Omega} (u - g)^2 dx + D(u, u_{k-1}) \rightarrow \min_u$$

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Convergence Analysis

The second term is called **Bregman distance**

$$D(u, u_{k-1}) = J(u) - J(u_{k-1}) - \langle v_{k-1}, u - u_{k-1} \rangle$$

with $v_{k-1} \in \partial J(u_{k-1})$



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Convergence Analysis

Example: quadratic case

$$J(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 \, dx$$

yields

$$D(u, u_{k-1}) = \frac{1}{2} \int_{\Omega} |\nabla u - \nabla u_{k-1}|^2 \, dx$$

Iterated Tikhonov / Levenberg-Marquardt



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Inverse TV Flow

Does the reconstruction depend on λ ?

To understand the dependence on the parameter, consider limit $\lambda \rightarrow 0$.

Let $\lambda_N = T/N$, set $t_k = k \lambda_N$. We compute the solution of the iterated TV flow with this specific parameter and set

$$u^N(t_k) = u_k$$

Linear interpolation for other times.



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Inverse TV Flow

We obtain a limiting flow,

$$u^N(t) \rightarrow u(t)$$

This flow is uniquely defined

Formally, $\frac{\partial}{\partial t} J'(u(t)) = f - u(t)$



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Inverse TV Flow

The „inverse TV flow“ can be considered as an „inverse scale space method“
(Scherzer-Weickert 99)

Start with zero image, and gradually insert smaller and smaller scales. Noise should return in the end, this happened in experiments



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Inverse TV Flow

The parameter λ can be interpreted as a time step for an implicit discretization of the inverse TV flow

Should work well and be independent of the parameter, as long as λ is small enough. This can be observed in experiments, too.



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Nonlinear Inverse Problems

The algorithm can immediately be generalized to arbitrary linear and even nonlinear inverse problems:

Find next iterate as minimizer of

$$Q_k(u) = \frac{\lambda}{2} \|F(u) - g\|^2 + D(u, u_{k-1})$$



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Nonlinear Inverse Problems

Example: estimate diffusion coefficient q in

$$-\operatorname{div}(q \nabla u) = f$$

from measurement of u

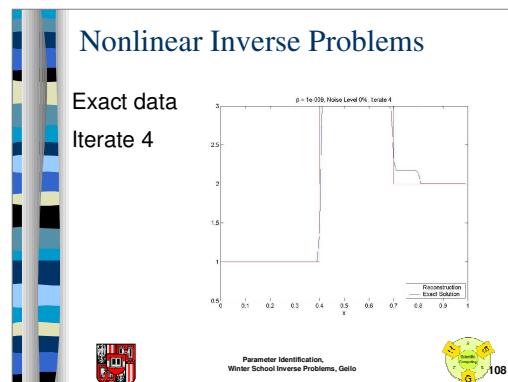
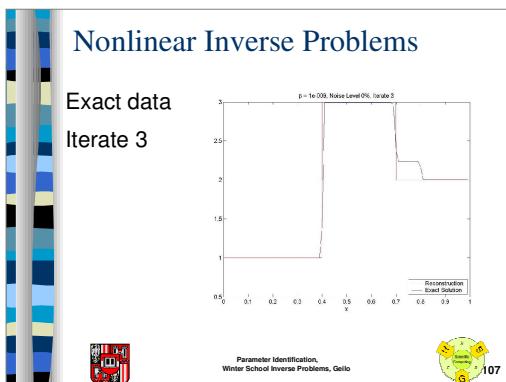
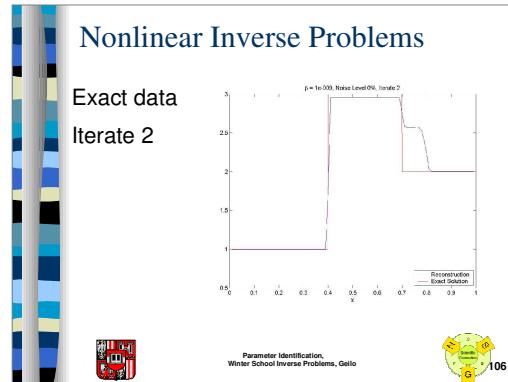
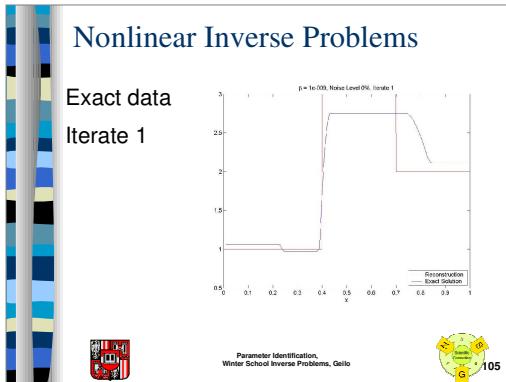
Operator $F: q \rightarrow u(q)$

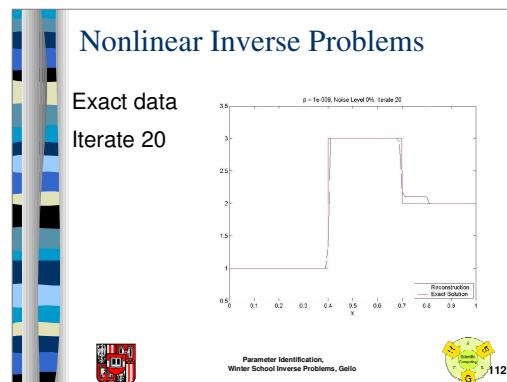
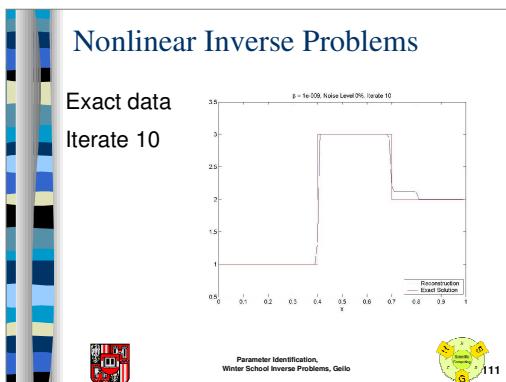
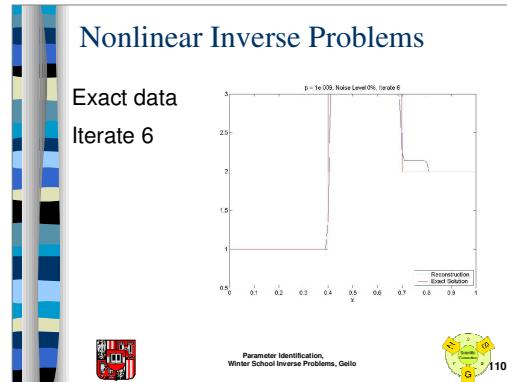
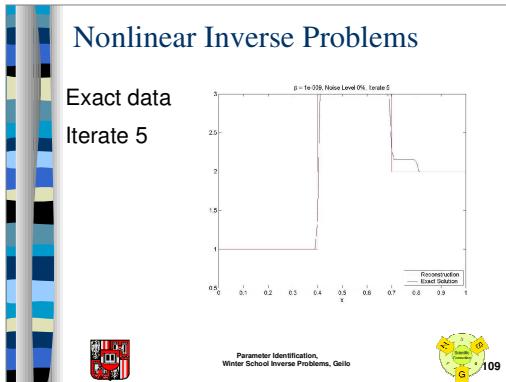
Application: find material layers in groundwater modeling (piecewise constant)

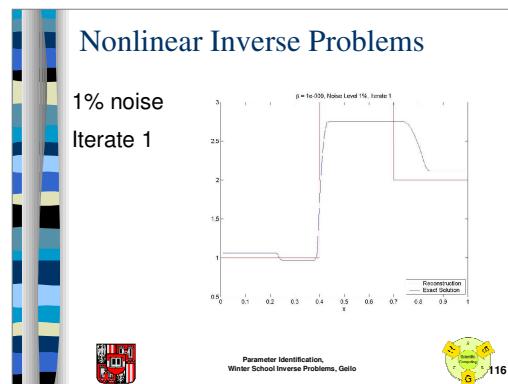
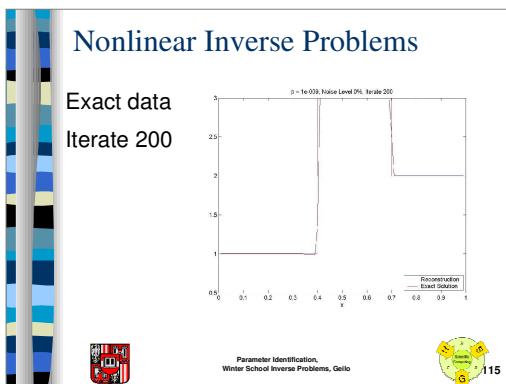
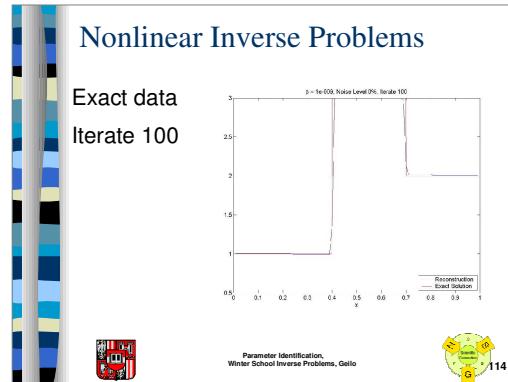
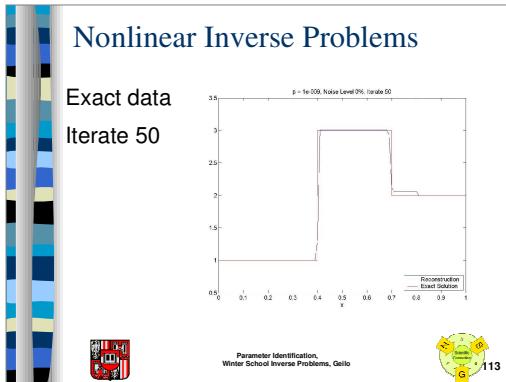


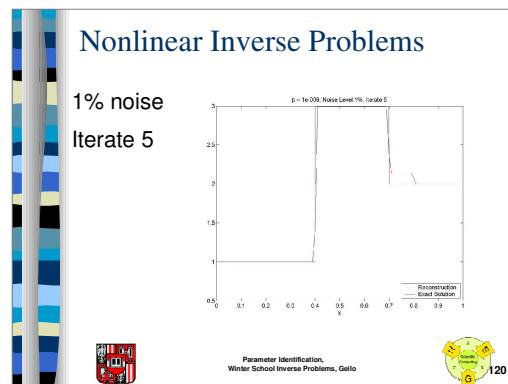
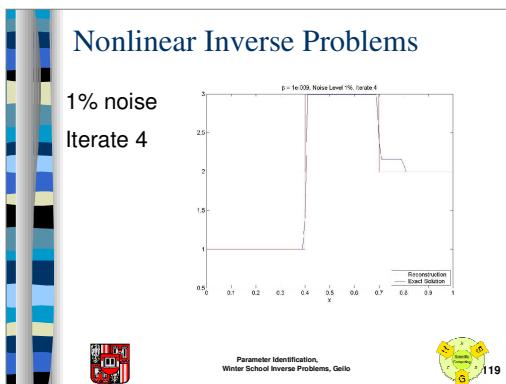
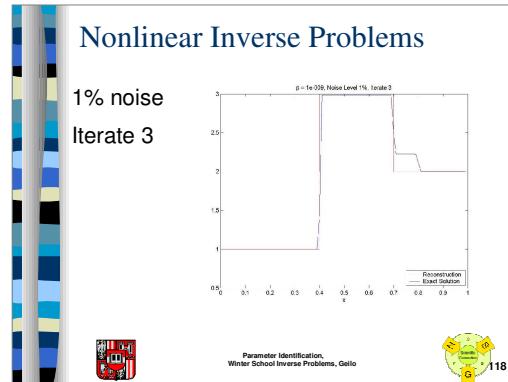
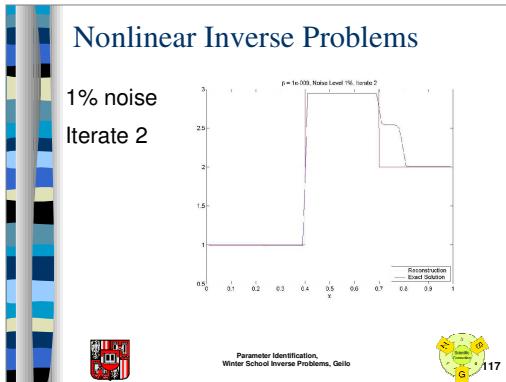
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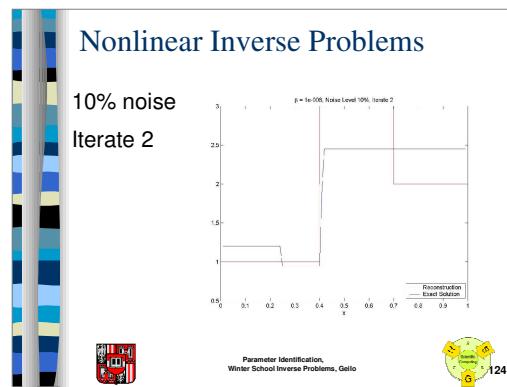
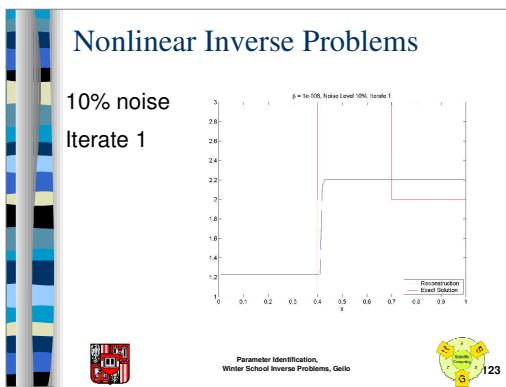
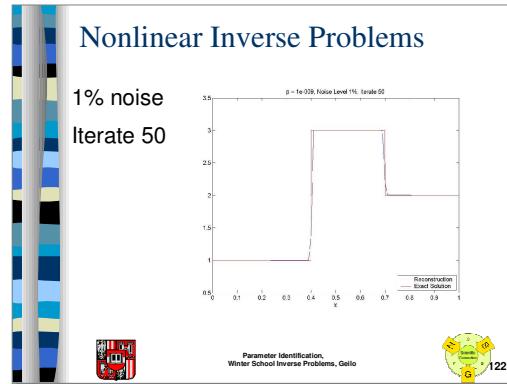
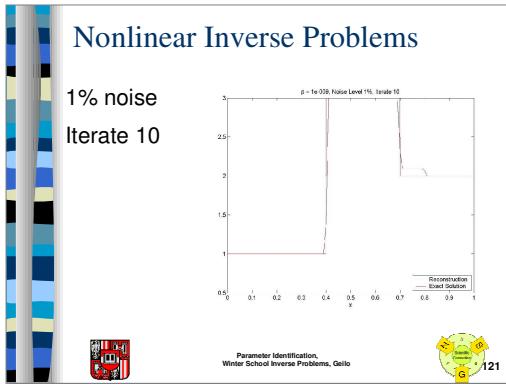


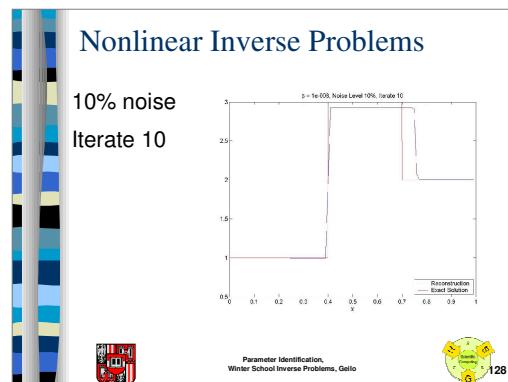
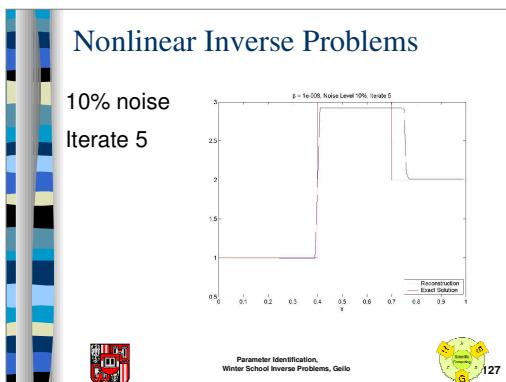
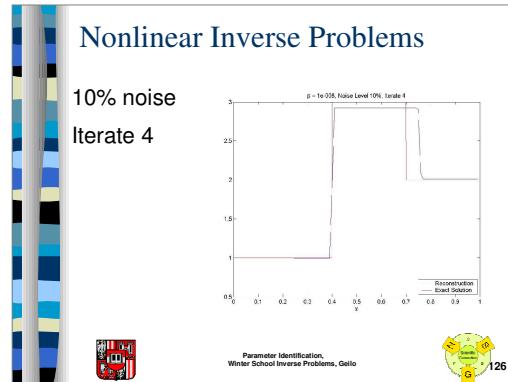
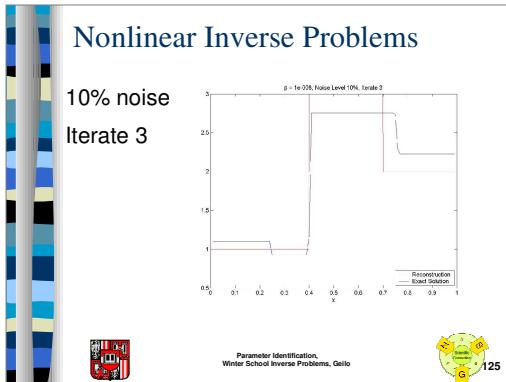












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Implementation Aspects

Landweber iteration:

- ▷ Discretize parameter
- ▷ Discretize state + state equation
- ▷ Compute gradient by adjoint method
- ▷ Try to use same kind of discretization for adjoint as for state equation („Discretized Adjoint = Adjoint of Discretized Problem“)
- ▷ For multiple state equations compute gradients corresponding to each state equation immediately

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Implementation Aspects

Landweber iteration with Black Box Solver:

- ▷ Discretize parameter
- ▷ Discretize state + state equation (determined by solver)
- ▷ If adjoint method not possible, try to compute gradients by finite differencing
- ▷ For multiple state equations compute gradients corresponding to each state equation immediately

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Implementation Aspects

Quasi-Newton Methods (BFGS):

- ▷ Discretize parameter
- ▷ Discretize state + state equation (determined by solver)
- ▷ Compute gradients by adjoint method, use adjoint to update quasi-Newton matrix
- ▷ Solve linear system for update

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Implementation Aspects

Newton Methods (LM / IRGN / NCG):

Discretize parameter

Discretize state + state equation

Never compute Newton-Matrix (NOT SPARSE)!!

$$F'(a) = -B \circ \frac{\partial e}{\partial u}(\Phi(a); a)^{-1} \circ \frac{\partial e}{\partial a}(\Phi(a); a)$$

Try to use iterative method for computing the Newton step, only implement application of $F'(a)$ and its adjoint (two PDEs)



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Implementation Aspects

One shot methods:

Discretize parameter + state + state equation + adjoint + adjoint equation

Solve the full KKT system simultaneously, e.g.

$$\begin{aligned} 0 &= 2\alpha^k(a^{k+1} - a^k) + \frac{\partial e}{\partial a}(u^k; a^k)^* w^k \\ 0 &= 2B^*(Bu^k - y^\delta + Bv^k) + \frac{\partial e}{\partial a}(u^k; a^k)^* w^k \\ 0 &= \frac{\partial e}{\partial u}(u^k; a^k)v^k + \frac{\partial e}{\partial a}(u^k; a^k)(a^{k+1} - a^k) \end{aligned}$$



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Implementation Aspects

One shot methods:

KKT system is **sparse and symmetric**

KKT system is not positive definite (no CG !)

Use BiCGStab, MINRES, QMR or GMRES

Apply appropriate preconditioning, note that there is a small parameter (α), acting like singular perturbation

Preconditioning strategies: Battermann-Sachs 2000,
Battermann-Heinckenschloss 2001, Ascher-Haber 2001-2003,
B-Müllhuber 2002, Griewank 2005



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Kaczmarz Methods

For problems like impedance tomography, the operators are of the form

$$e(u; a) = (e_1(u^1; a), \dots, e_N(u^N; a)) = 0$$

$$Bu = (B_1 u^1, \dots, B_N u^N)$$

Idea: perform multiplicative splitting (like in Gauss-Seidel), cyclic iteration over the N subproblems



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Kaczmarz Methods

Subproblem j :

$$e_j(u^j; a) = 0, \quad B_j u^j = y_j^\delta$$

Landweber-Kaczmarz (Natterer 1996, Kowar-Scherzer 2003)

$$e(u^{k,j}; a^{k,j-1}) = 0, \quad \frac{\partial e_j}{\partial u}(u^{k,j}; a^{k,j-1})^* w^{k,j} + B_j^*(B_j u^{k,j} - y_j^\delta) = 0$$

$$a^{k,j} = a^{k,j-1} + \tau^j \frac{\partial e_j}{\partial a}(u^j; a^{k,j-1})^* w^j, \quad j = 1, \dots, N$$


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Kaczmarz Methods

Levenberg-Marquardt Kaczmarz (B-Kaltenbacher 2004)

State equation: $e(u^{k,j}; a^{k,j-1}) = 0$

Linear KKT:

$$B^*(Bu^{k,j} + Bv^{k,j} - y^\delta) + \frac{\partial e}{\partial u}(u^{k,j}; a^{k,j-1})^* w^{k,j} = 0$$

$$\alpha_{kj}(a^{k,j} - a^{k,j-1}) + \frac{\partial e}{\partial a}(u^{k,j}; a^{k,j-1})^* w^{k,j} = 0$$

$$\frac{\partial e}{\partial u}(u^{k,j}; a^{k,j-1}) v^{k,j} + \frac{\partial e}{\partial a}(u^{k,j}; a^{k,j-1})(a^{k,j} - a^{k,j-1}) = 0$$


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Kaczmarz Methods

Simple model problem: identify q in

$$-\Delta u_j + qu_j = 0$$

from measurements $f_j = u_j$

for different sources $\frac{\partial u_j}{\partial \nu} = g_j$



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Numerical Solution

Newton-Kaczmarz approach is equivalent to minimizing

$$\frac{1}{2} \left\| \frac{\partial v_n}{\partial \nu} - g_n \right\|_{H^{-1/2}(\partial \Omega)}^2 + \frac{\alpha_n}{2} \|s + q_n - q_0\|_{H^1(\Omega)}^2$$

for $s = q_{n+1} - q_n$ in each step, subject to

$$-\Delta v_n + q_n v_n + s u_n = 0$$


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Numerical Solution

Since we need another function to realize the $H^{1/2}$ -norm, corresponding KKT system is 4×4

$$-\Delta v_n + qv_n + su_n = 0$$

$$-\Delta \lambda + q\lambda = 0$$

$$-\Delta \mu + \mu + (1 - q)v_n - su_n = 0$$

$$-\Delta s + s + \frac{1}{\alpha} u_n \lambda = 0$$

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Numerical Solution

With boundary conditions

$$v_n = 0$$

$$\lambda - \mu = \phi_n - \phi_{g_n}$$

$$\frac{\partial \mu}{\partial \nu} = 0$$

$$s = 0$$

and

$$\int_{\Omega} (\nabla \phi_g \cdot \nabla \psi + \phi_g \psi) dx = \int_{\partial\Omega} g \psi d\sigma \quad \forall \psi \in H^1(\Omega)$$

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Numerical Results

Numerical experiment with $N=20$ localized sources g_j (5 on each side)

Exact solution

$$\hat{q} = 3 + 5 \sin(\pi x) \sin(\pi y)$$

Constant initial value $q_0 \equiv 3$

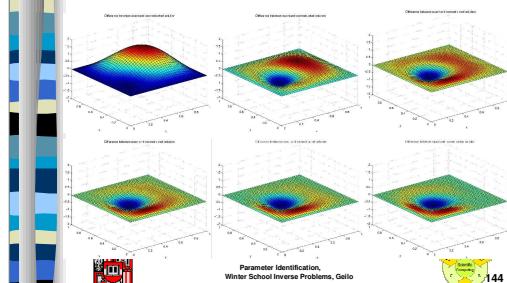
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Numerical Results

Iterates 1, 2, 3, 4, 5, 6



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